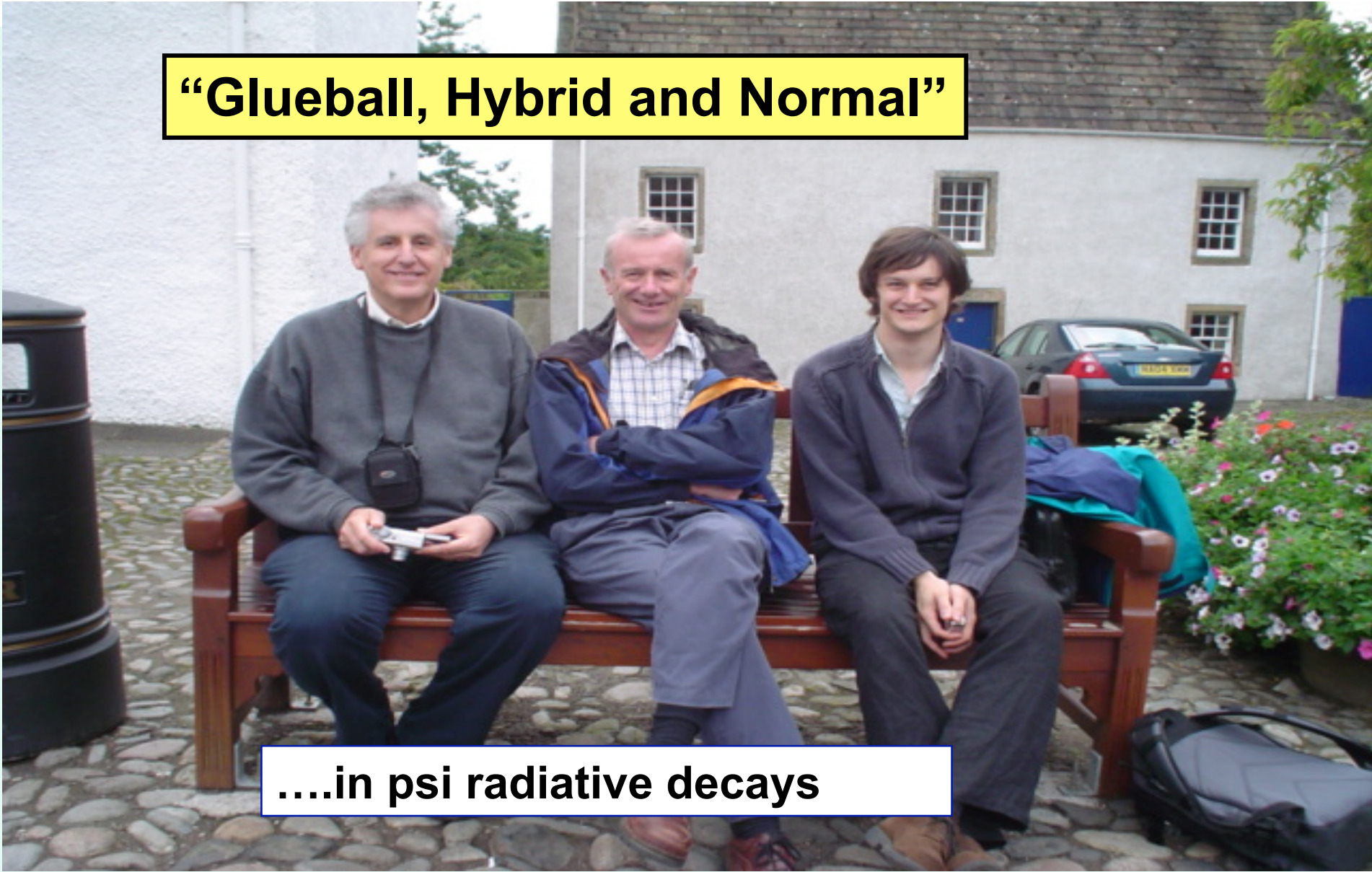


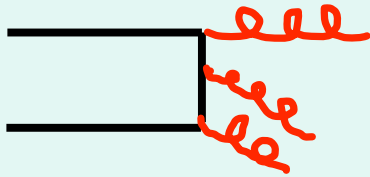
**“Glueball, Hybrid and Normal”**



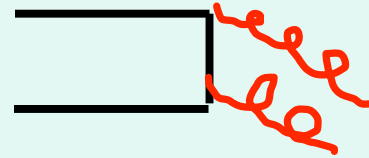
**....in psi radiative decays**

**Frank Close  
SURA05**

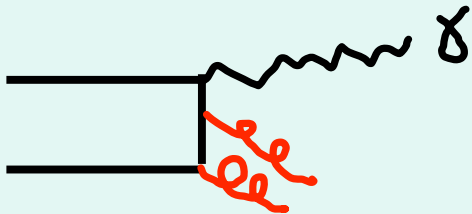
# pQCD Folklore



Psi hadronic and chi\_1 widths.....

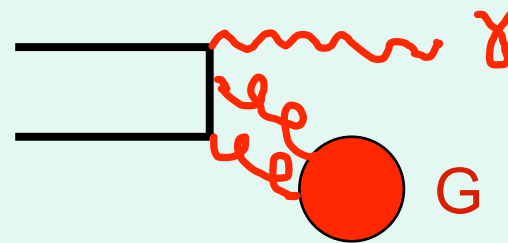


....Smaller than chi\_0 and chi\_2 widths



$$\frac{\Gamma_B(\psi(3.1) \rightarrow \gamma + 2 \text{ gluons} \rightarrow \text{hadrons})}{\Gamma_{\text{tot}}(\psi(3.1) \rightarrow 3 \text{ gluons} \rightarrow \text{hadrons})} = \frac{16}{5} \frac{\alpha}{\alpha_s} \left( 1 - 2.6 \frac{\alpha_s}{\pi} \right)$$

Expt ~ 8% .....



..... dominated by glueballs



$$\Gamma(\psi) = 91 \text{ keV}. \quad (12 \pm 2 \text{ via } \gamma^* = \text{e.m.})$$

$$\Gamma(\eta_c) = 17.3 \pm 2.7 \text{ MeV}$$

Consistent with 3g versus 2g



$$\Gamma(\chi_0) = 10 \text{ MeV}; \quad \chi_1 = 0.9 \text{ MeV}; \quad \chi_2 = 2.1 \text{ MeV}$$

Consistent with axial decoupled from 2g



$$\frac{\Gamma(\chi_0 \rightarrow \gamma\gamma)}{\Gamma(\chi_2 \rightarrow \gamma\gamma)} \sim 6 > \frac{15}{4}$$

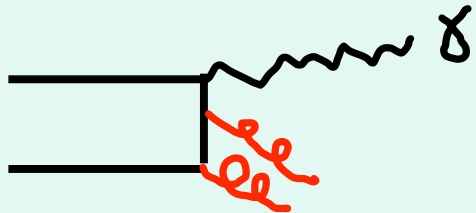
Approx consistent with 2g

?



(data vary widely; need precision)

# pQCD Folklore



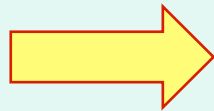
$$\frac{\Gamma_B(\psi(3.1) \rightarrow \gamma + 2 \text{ gluons} \rightarrow \text{hadrons})}{\Gamma_{\text{tot}}(\psi(3.1) \rightarrow 3 \text{ gluons} \rightarrow \text{hadrons})} = \frac{16}{5} \frac{\alpha}{\alpha_s} \left( 1 - 2.6 \frac{\alpha_s}{\pi} \right)$$

Expt ~ 8% .....

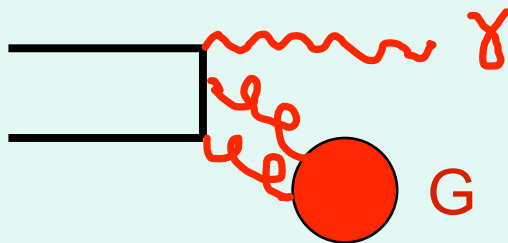


Qualitatively expect this order of magnitude anyway?

(What is actual precision b.r.?)



Need x (M\_hadron) dependence pQCD; + expt precision not just integrated rate.

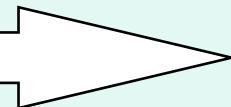


..... dominated by glueballs

?

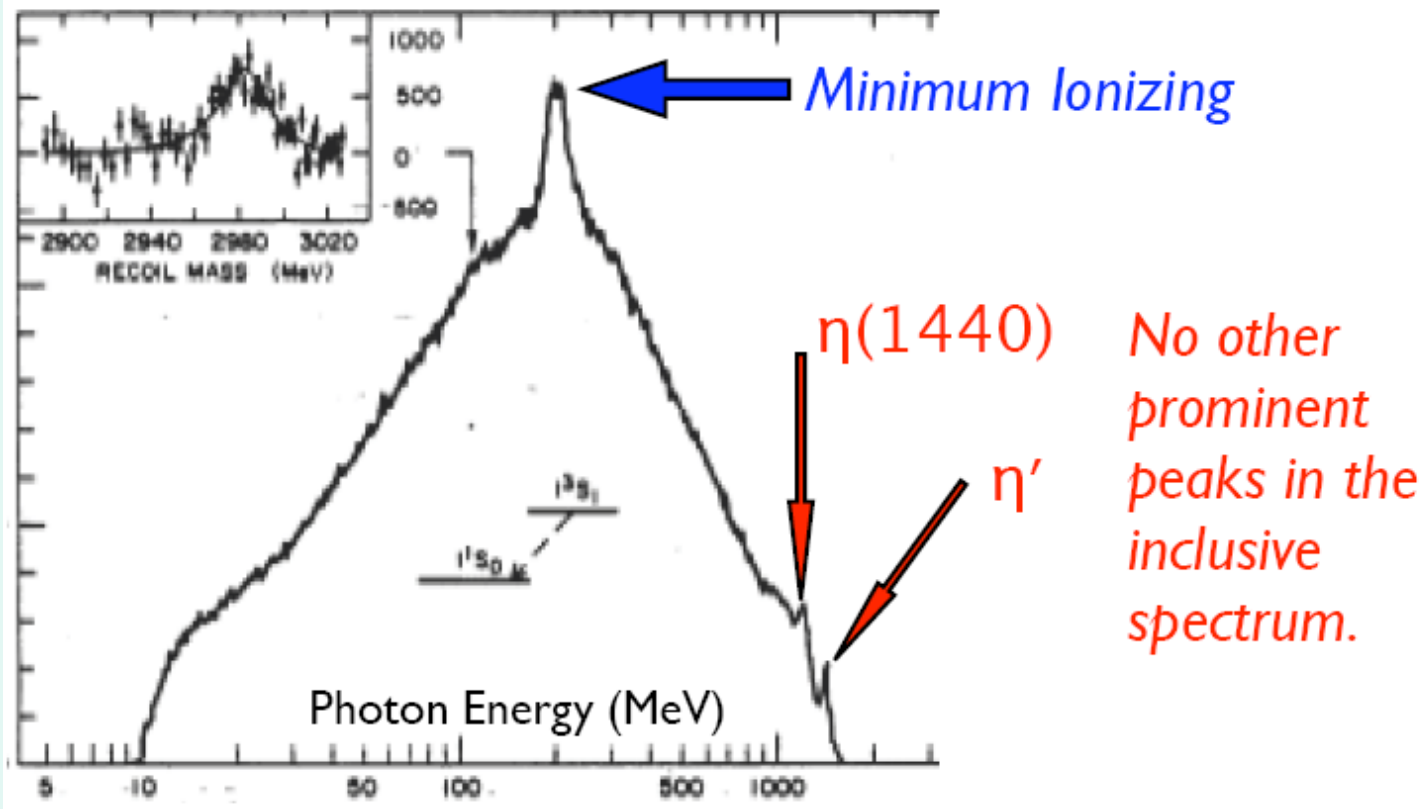


next



# psi to gamma X (see Jim Napolitano talk)

Bloom & Peck (Crystal Ball) *Ann.Rev.Nuc.Part.Sci.* 33(1983)143



# Some prominent radiative b.r.'s ( $10^{-4}$ )

total  $\psi \rightarrow \gamma X \sim 800 (10^{-4})$

**0-+:**  $\eta$  8;  $\eta'$  43;  $\eta(1405/1475) > 30$

**2++:**  $f_2(1270)$  14;  $f_2(1525)$  5;  $f_2(1950) > 7$

**1++:**  $f_1(1285)$  6;  $f_1(1420)$  8;  $f_1(1510)$  4?

----> Prominent  $qq^*$  2++ which are not glueballs  
and prominent 1++ which pQCD naively thought to be zero

# Some prominent radiative b.r.'s ( $10^{-4}$ )

total psi  $\rightarrow$  gamma X  $\sim 800 (10^{-4})$

**0-+:** eta 8; eta' 43; eta(1405/1475)  $> 30$

**2++:** f2(1270) 14; f2(1525) 5; f2(1950)  $> 7$

**1++:** f1(1285) 6; f1(1420) 8; f1(1510) 4?

**0++:** f0(1370) 4pi  $> 80$ ; f0(1500)  $> 6$ ; f0(1700)  $> 9$

**0++:** f0(980) = ??

# Some prominent radiative b.r.'s ( $10^{-4}$ )

total  $\psi \rightarrow \gamma X \sim 800 (10^{-4})$

**0-+:**  $\eta$  8;  $\eta'$  43;  $\eta(1405/1475) > 30$

$\eta \pi \pi > 60$

VV in  $0^- > 40?$

**0++:**  $f_0(1370) 4\pi > 80$ ;  $f_0(1500) > 6$ ;  $f_0(1700) > 9$

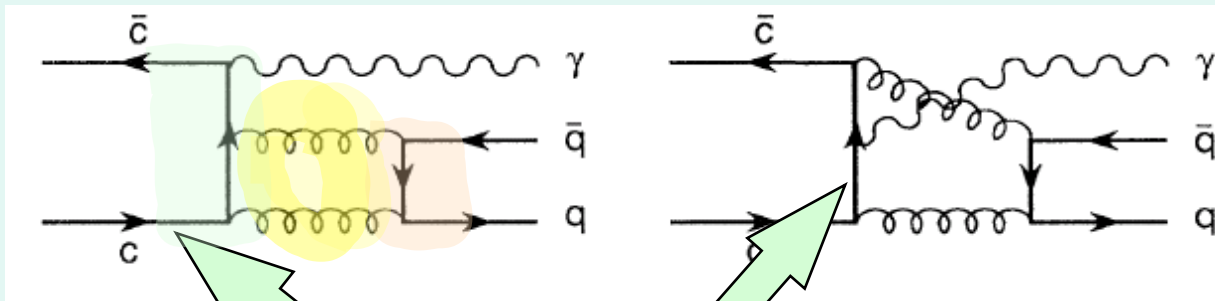
$4\pi > 120$

JPC need measuring

Isolating  $0^{++}$  a challenge anyway. Also look for exotic  $1^-+$



# The JPC dependence of psi radiative in pQCD



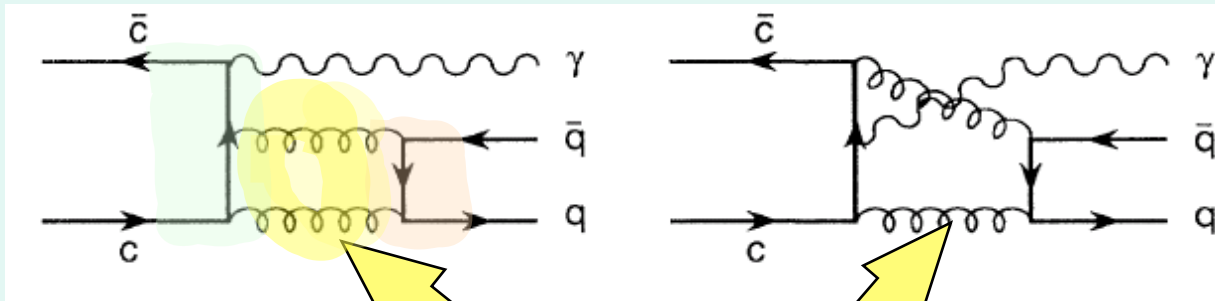
Psi decay piece

$$b_{\text{rad}}(\overline{Q}Q_V \rightarrow \gamma + R_J) = \frac{c_{R^X} |H_J(x)|^2}{8\pi(\pi^2 - 9)} \frac{m_R}{M_V^2} \Gamma(R_J \rightarrow gg)$$

qq\* to gamma gamma  
+ pQCD =  
qq\* to gg

Close Farrar Li

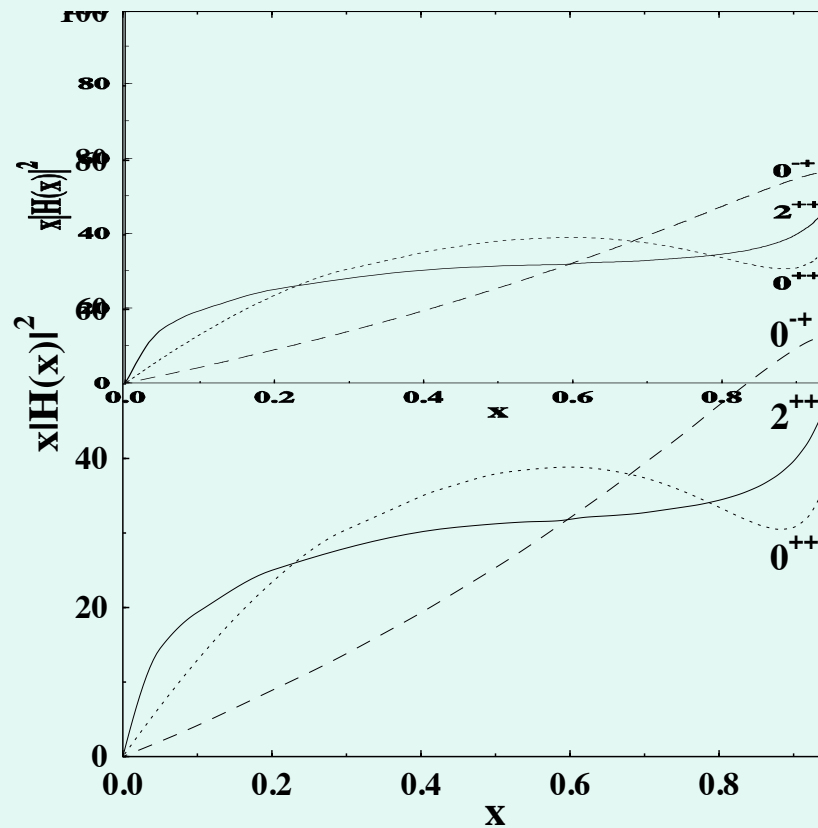
# The JPC dependence of psi radiative in pQCD



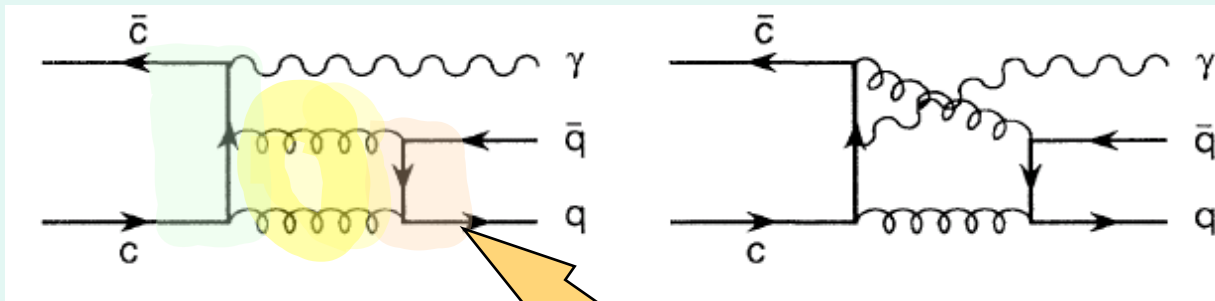
$H_J(x)$  Loop integral

$$b_{\text{rad}}(\overline{Q}Q_V \rightarrow \gamma + R_J) = \frac{c_{R^X} |H_J(x)|^2}{8\pi(\pi^2 - 9)} \frac{m_R}{M_V^2} \Gamma(R_J \rightarrow gg)$$

What  $HJ(x)$  look like for  $0-, 0+, 2+$



# The JPC dependence of psi radiative in pQCD

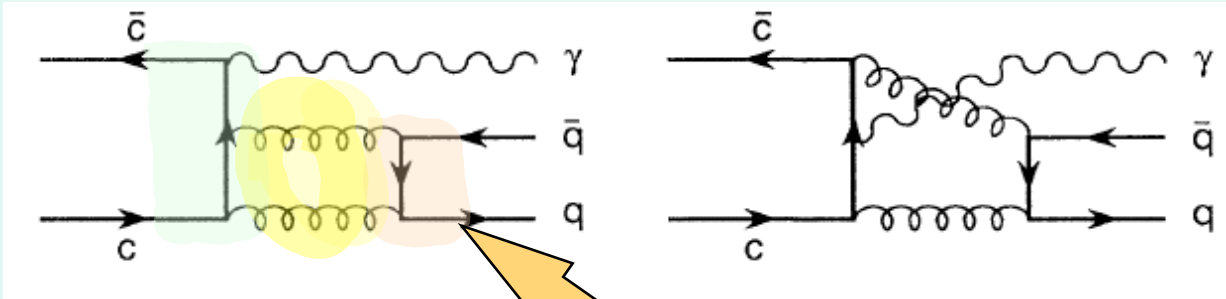


using data as input.....

...quantify this gg measure

$$b_{\text{rad}}(\overline{Q}Q_V \rightarrow \gamma + R_J) = \frac{c_{R^X} |H_J(x)|^2}{8\pi(\pi^2 - 9)} \frac{m_R}{M_V^2} \Gamma(R_J \rightarrow gg)$$

# The JPC dependence of psi radiative in pQCD



using data as input.....

...quantify this gg measure

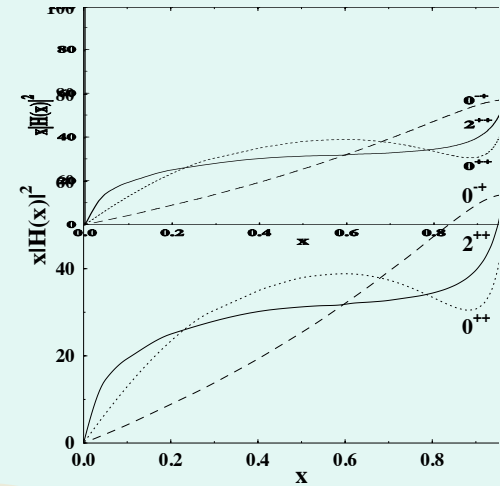
$$b_{\text{rad}}(\overline{Q}Q_V \rightarrow \gamma + R_J) = \frac{c_{R^X} |H_J(x)|^2}{8\pi(\pi^2 - 9)} \frac{m_R}{M_V^2} \Gamma(R_J \rightarrow gg)$$

b.r(R → gg)



qq\* in gamma gamma convert to gg  
Close Farrar Li

Rewrite the b.r. in a more user friendly way.....



$$10^3 br(J/\psi \rightarrow \gamma 2^{++}) = \left(\frac{m}{1.5 \text{ GeV}}\right) \left(\frac{\Gamma_{R \rightarrow gg}}{26 \text{ MeV}}\right) \frac{x |H_S(x)|^2}{35}$$

$$10^3 br(J/\psi \rightarrow \gamma 0^{++}) = \left(\frac{m}{1.5 \text{ GeV}}\right) \left(\frac{\Gamma_{R \rightarrow gg}}{96 \text{ MeV}}\right) \frac{x |H_S(x)|^2}{35}$$

$$10^3 br(J/\psi \rightarrow \gamma 0^{-+}) = \left(\frac{m}{1.5 \text{ GeV}}\right) \left(\frac{\Gamma_{R \rightarrow gg}}{50 \text{ MeV}}\right) \frac{x |H_S(x)|^2}{45}$$

O(1)

O(1)

## Tensor 2++

$$10^3 b(J/\psi \rightarrow \gamma 2^{++}) = \left( \frac{m}{1.5 \text{ GeV}} \right) \left( \frac{\Gamma_{R \rightarrow gg}}{26 \text{ MeV}} \right) \frac{x |H_T(x)|^2}{34}$$

f2(1270)

1.4 ± 0.14

Total width = 185 ± 20 MeV

41 ± 7 MeV

b.r = 0.22 ~

$O(\alpha_s^2) \sim 0.1-0.2$

## Tensor 2++

$$10^3 b(J/\psi \rightarrow \gamma 2^{++}) = \left( \frac{m}{1.5 \text{ GeV}} \right) \left( \frac{\Gamma_{R \rightarrow gg}}{26 \text{ MeV}} \right) \frac{x |H_T(x)|^2}{34}$$

f2(1270)

1.4 ± 0.14

41 ± 7 MeV

Total width = 185 ± 20 MeV

b.r = 0.22 ~

$O(\alpha_s^2) \sim 0.1-0.2$

f2(1520)

0.45 ± 0.07

12 ± 2 MeV

Total width = 76 ± 10 MeV

b.r = 0.15 ~

$O(\alpha_s^2) \sim 0.1-0.2$



# Axial 1++

$$10^3 b(J/\psi \rightarrow \gamma 1^{++}) = \left( \frac{m}{1.45 \text{ GeV}} \right) \left( \frac{\Gamma_{R \rightarrow gg}}{12 \text{ MeV}} \right) \frac{x |H_1(x)|^2}{30}$$

f1(1285)

0.65 +/- 0.1

Total width = 24 +/- 3 MeV

~ 8 MeV

b.r ~ 0.3

$O(\alpha_s^2) \sim 0.1-0.2$

f1(1420)

also

b.r ~ 0.3

## Scalar 0<sup>++</sup>

$$10^3 b(J/\psi \rightarrow \gamma 0^{++}) = \left( \frac{m}{1.5 \text{ GeV}} \right) \left( \frac{\Gamma_{R \rightarrow gg}}{96 \text{ MeV}} \right) \frac{x |H_S(x)|^2}{35}$$

f<sub>0</sub>(1710)

0.85 ± 0.1

Total width ~ 140 MeV

$$\Gamma(f_0(1710) \rightarrow gg) = \frac{(78 \pm 10) \text{ MeV}}{b(f_0(1710) \rightarrow KK)}$$

$$b(f_0(1710) \rightarrow gg) \geq 0.52 \pm 0.07$$

## Scalar $0^{++}$

$$10^3 b(J/\psi \rightarrow \gamma 0^{++}) = \left( \frac{m}{1.5 \text{ GeV}} \right) \left( \frac{\Gamma_{R \rightarrow gg}}{96 \text{ MeV}} \right) \frac{x |H_S(x)|^2}{35}$$

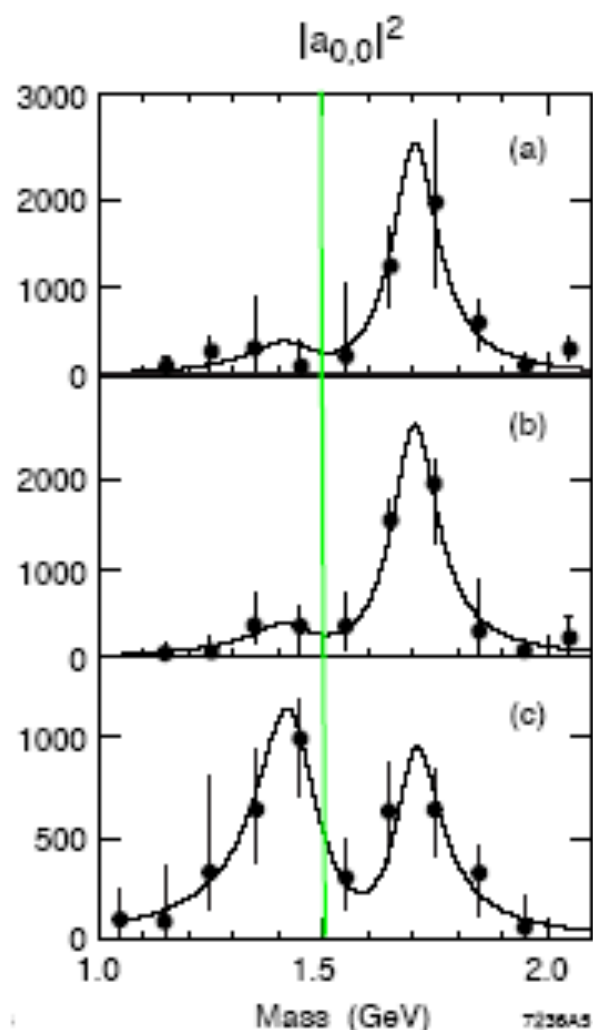
$f_0(1500)$

$$(0.57 \pm 0.08) \times 10^{-3} \leq b(J/\psi \rightarrow \gamma f_0(1500)) \leq (1.15 \pm 0.15) \times 10^{-3}$$

$$0.5 \pm 0.1 \leq b(f_0(1500) \rightarrow gg) \leq 0.9 \pm 0.2$$

# “Partial Wave Analysis”

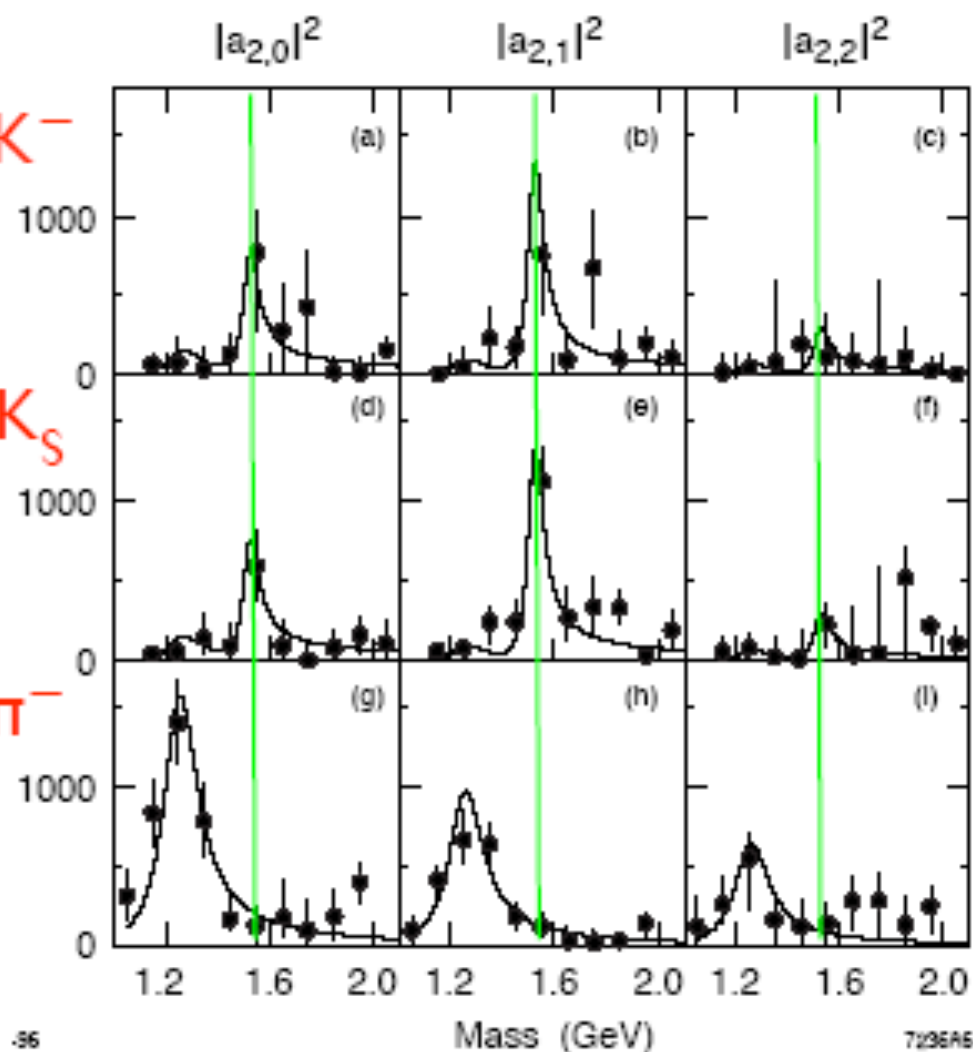
SLAC Mark-III (Dunwoodie, et al.)



$K^+K^-$

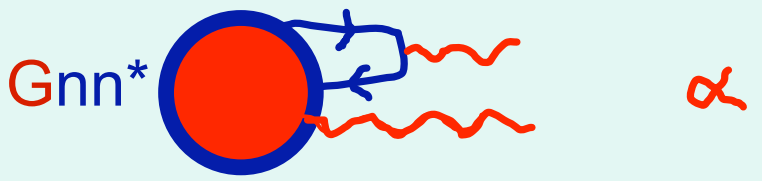
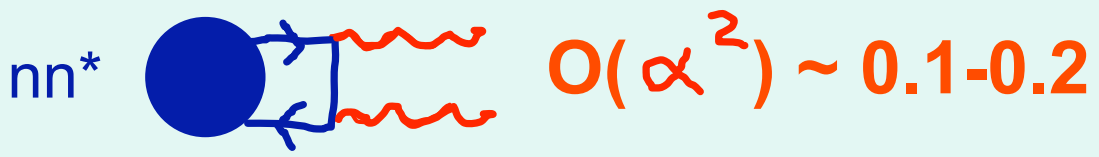
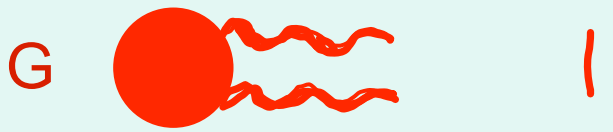
$K_S K_S$

$\pi^+\pi^-$



b.r(R → gg)

# HYBRID MESONS TOO?



should be significant

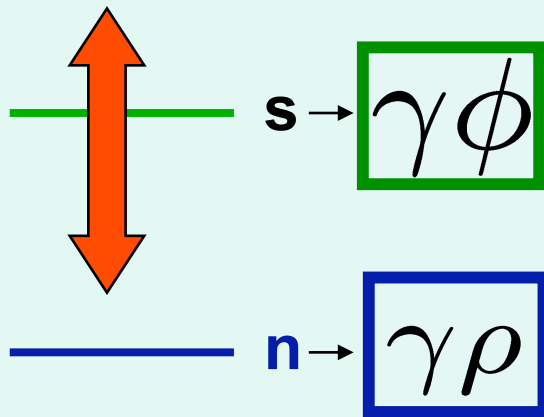
Needs calculation of HJ(x) for 1-+  
 apply to l=0 partners of pi1(1600)  
 Or 0-+, 2-+ hybrid components...  
 ....eta(1440)???

# More on Glue in $C=+$ Mesons

- How psi radiatives can sort out the PDG

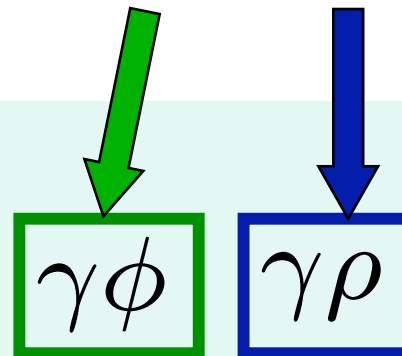
# Scalar Glueball and Mixing

a simple example for expt to rule out



Meson	G	$ss^*$	$nn^*$
1710	0.39	0.91	0.15
1500	-0.65	0.33	-0.70
1370	-0.69	0.15	0.70

0- 0-  
meson  
decays  
LEAR/  
WA102  
  
FC Kirk



$$\psi \rightarrow \gamma [\gamma V]$$

>1 billion

Coming soon  
from BES and  
CLEO-c

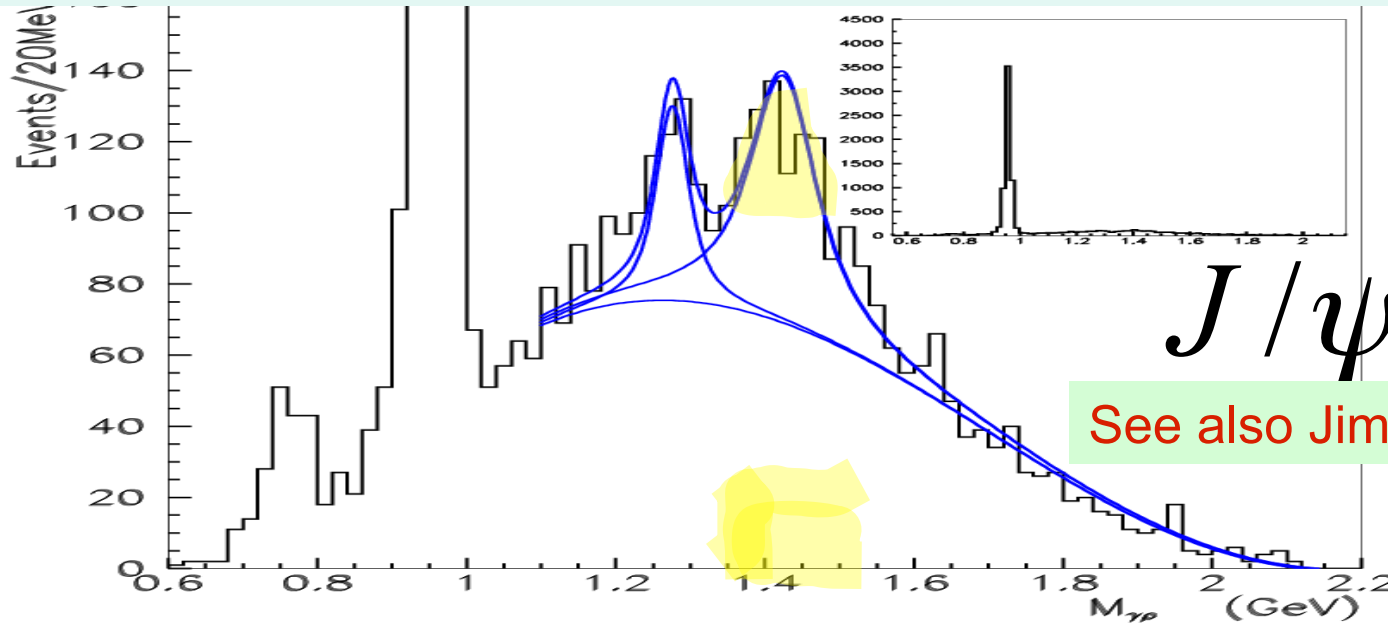
1000 per meson

**A flavour filter for  
0++ 0-+ 2++ (and 1++)  
mesons glueballs et al**

Challenge:

Turn Lattice QCD Glueball spectrum into physics

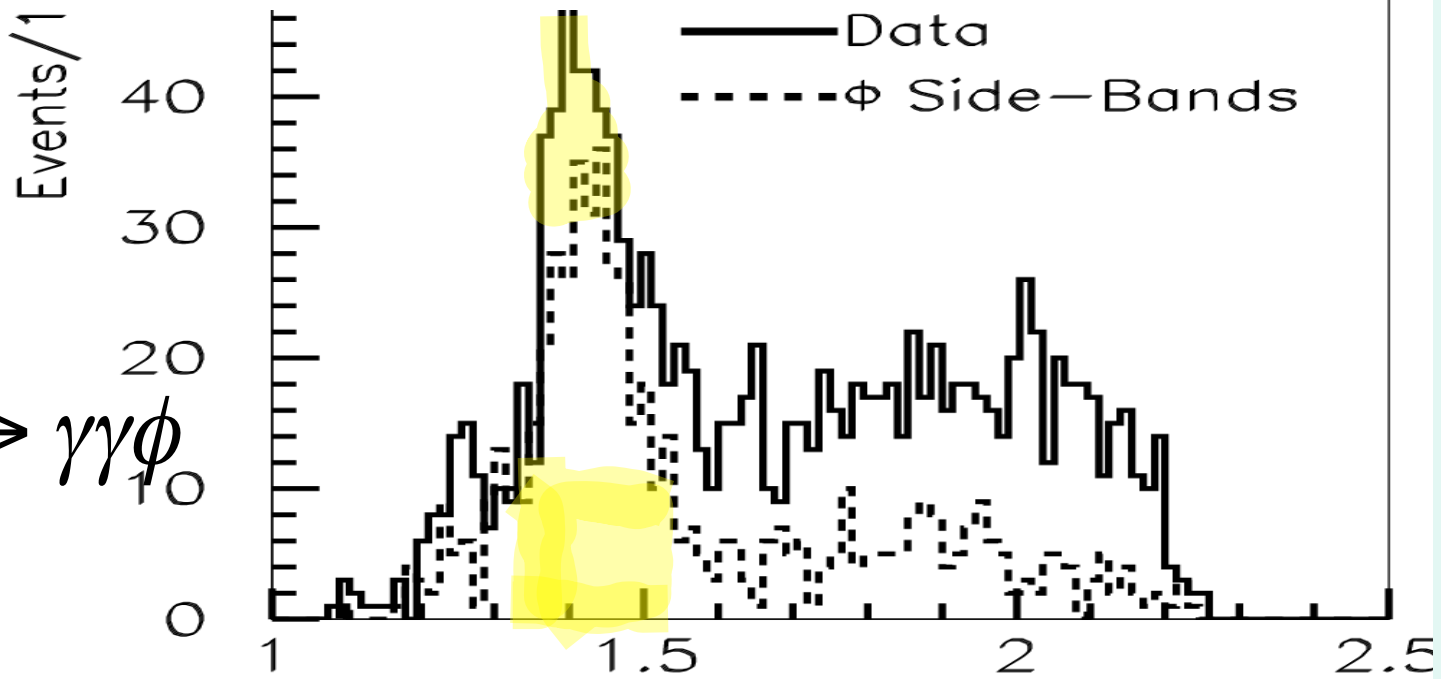




$J/\psi \rightarrow \gamma\gamma\rho$

See also Jim Napolitano talk

$J/\psi \rightarrow \gamma\gamma\phi$



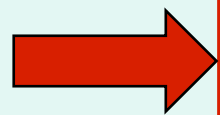
# Unusual properties of $f_0(1370)$ , $f_0(1500)$ $f_0(1710)$

## Scalar Puzzle

$$\psi \rightarrow \gamma [\gamma V]$$

$$\psi \rightarrow 0^{++} V$$

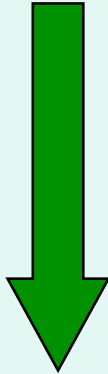
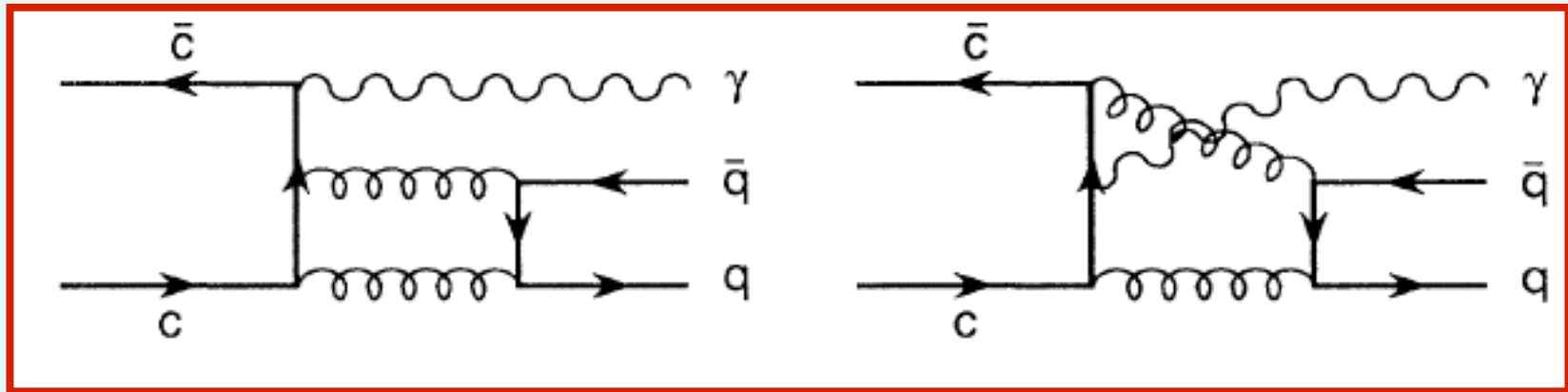
A consistent pattern in these two experiments can establish role of **Scalar Glueball**



**Challenge: quantify the predictions**

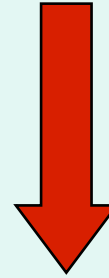
# A further puzzle with Tensors in $\psi$ radiative

- $2^{++}$  can have helicity 0, 1, 2
- Ratio of helicity 2:0 =  $y$ ; 1:0 =  $x$
- What do pQCD and data say?



$$y=0, \quad x = \frac{\sqrt{3}}{2} \frac{M_T}{M_\psi},$$

**Expt:  $y=0; x=0.9$**



$$y = \sqrt{6} \frac{1}{(1 - 2k^2/M_T^2)},$$

$$x = \sqrt{3} \frac{M_T}{M_\psi} \frac{M_T E_\psi - k^2}{M_T^2 - 2k^2}$$

$$\psi \rightarrow \gamma [\gamma V]$$

>1 billion

Coming soon  
from BES and  
CLEO-c

1000 per meson

**A flavour filter for  
0++ 0-+ 2++ (and 1++)  
mesons glueballs et al**

Challenge:

Turn Lattice QCD Glueball spectrum into physics

Lose factor 2 and do psiprime for chi?

$qq^*$  in psi rad can quantify if have gamma gamma

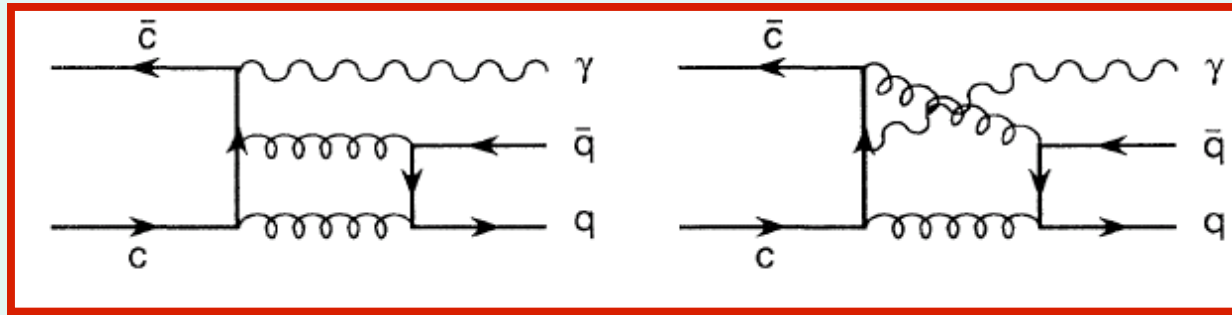
pp to pRp compare signals in psi to gamma R

Psi via gamma\* to "forbidden" states

If no big  $2^{++} \sim 2^{-2.5\text{gev}}$  in psi radiative: why not?

Lattice/phenom/principle challenge:

Why is iota(1440) not glueball?

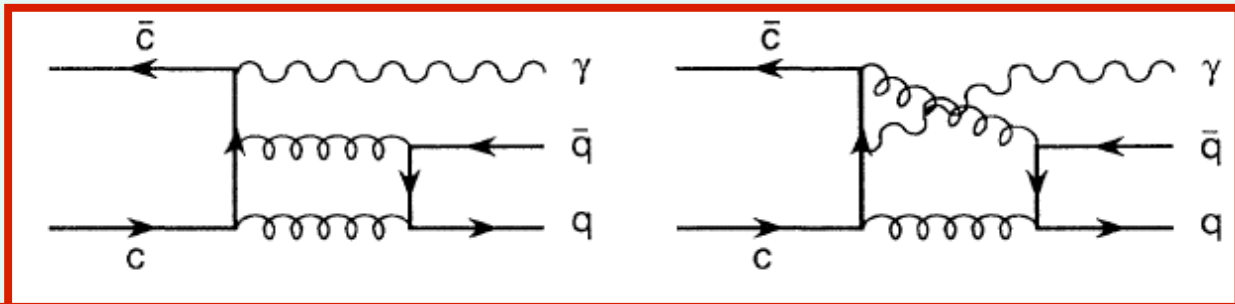


$$A_\lambda \propto \int d\Omega \{ A(\psi \rightarrow \gamma g + g -)(\theta) d_{2\lambda}^J(\theta) A(g + g - \rightarrow M) + A(\psi \rightarrow \gamma g + g +)(\theta) d_{0\lambda}^J(\theta) A(g + g + \rightarrow M) \}$$

$$Z \equiv (M_\psi^2 + M_T^2) / (M_\psi^2 - M_T^2)$$

$$I(Z, \lambda) \equiv \int \frac{d\Omega (\sin\theta)^{4-2\lambda} (1 - \cos\theta)^{2\lambda}}{Z^2 - \cos^2\theta}$$

	$A^C$	$A^H$
$A_2$	0	$kM_T D(Z, 2)$
$A_1$	$\frac{1}{\sqrt{2}} \frac{E_\psi}{M_\psi} k^2 D(Z, 1)$	$\frac{1}{\sqrt{2}} \left[ kM_\psi - \frac{k^2 E_\psi}{M_\psi} \right] D(Z, 1)$
$A_0$	$\frac{2}{\sqrt{6}} \frac{E_\psi}{M_T} k^2 D(Z, 0)$	$\frac{1}{\sqrt{6}} \frac{M_\psi}{M_T} \left[ kM_\psi - \frac{2k^2 E_\psi}{M_\psi} \right] D(Z, 0)$



	$A^C$	$A^H$
y:	$A_2$	$0$
x:	$A_1$	$\frac{1}{\sqrt{2}} \left[ kM_\psi - \frac{k^2 E_\psi}{M_\psi} \right] D(Z,1)$
1:	$A_0$	$\frac{1}{\sqrt{6}} \frac{M_\psi}{M_T} \left[ kM_\psi - \frac{2k^2 E_\psi}{M_\psi} \right] D(Z,0)$

$$y=0, \quad x = \frac{\sqrt{3}}{2} \frac{M_T}{M_\psi},$$

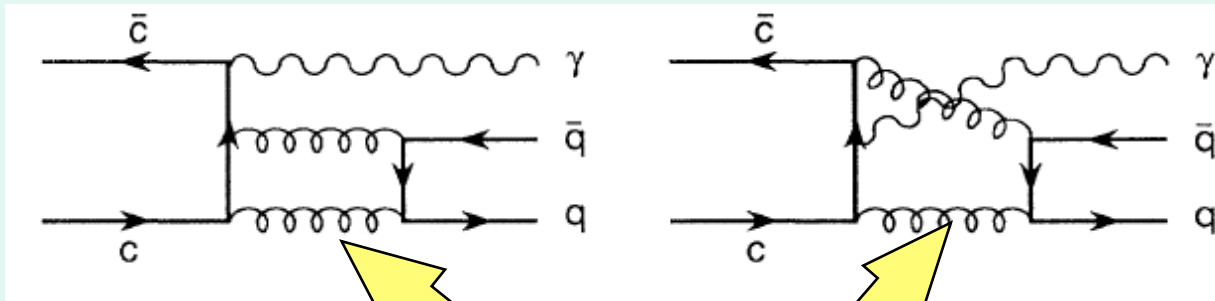
Expt:  $y=0; x=0.9$

$$y = \sqrt{6} \frac{1}{(1 - 2k^2/M_T^2)},$$

$$x = \sqrt{3} \frac{M_T}{M_\psi} \frac{M_T E_\psi - k^2}{M_T^2 - 2k^2}$$



# The JPC dependence of psi radiative in pQCD



$H_J(x)$  Loop integral

$$b_{\text{rad}}(\overline{Q}Q_V \rightarrow \gamma + R_J) = \frac{c_{R^X} |H_J(x)|^2}{8\pi(\pi^2 - 9)} \frac{m_R}{M_V^2} \Gamma(R_J \rightarrow gg)$$

Where just to show off:

$$H_{0^-+}(x) = \frac{4}{x} \left[ L(1-2x) - L(1) - \frac{1-x}{2-x} \left( 2L(1-x) - \frac{\pi^2}{3} + \frac{1}{2} \ln^2(1-x) \right) - \frac{x}{1-2x} \ln(2x) \right] + i4\pi \frac{1-x}{(2-x)x} \ln(1-x)$$

where  $L(x)$  is a Spence function, defined as

$$L(x) = - \int_0^x \frac{dx}{x} \ln(1-x).$$