

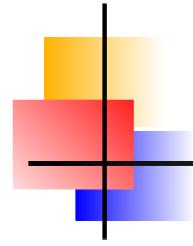
Centre-sector tunneling, confinement and the quark Fermi surface

ECT, Trento, Italy, 20 July 2010

Kurt Langfeld

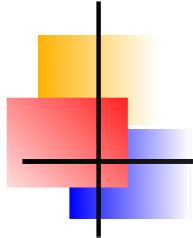
*School of Comp. and Mathematics and The HPCC,
Univ. of Plymouth, UK*

*Andreas Wipf, Björn Wellegehausen (PhD)
TPI, University of Jena, Germany*



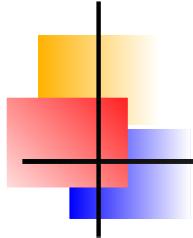
Introduction:

- Yang-Mills moduli and confinement



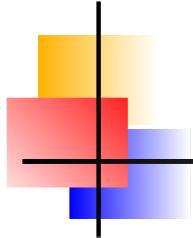
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- Phenomenological impact of centre-sector-tunneling
→ Fermi-Einstein condensation in $SU(2N)$ QCD-like theories



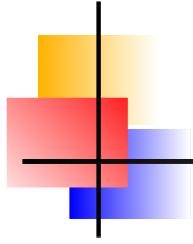
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- Does centre-sector-tunneling take place in the hadronic phase?
 - lattice gauge simulations: YM + qHiggs
 - centre-sector-tunneling and the 't Hooft loop
 - tunneling coefficient (new!)



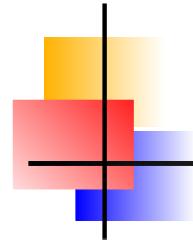
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- Does Fermi-Einstein condensation take place
in $SU(3)$ with matter?



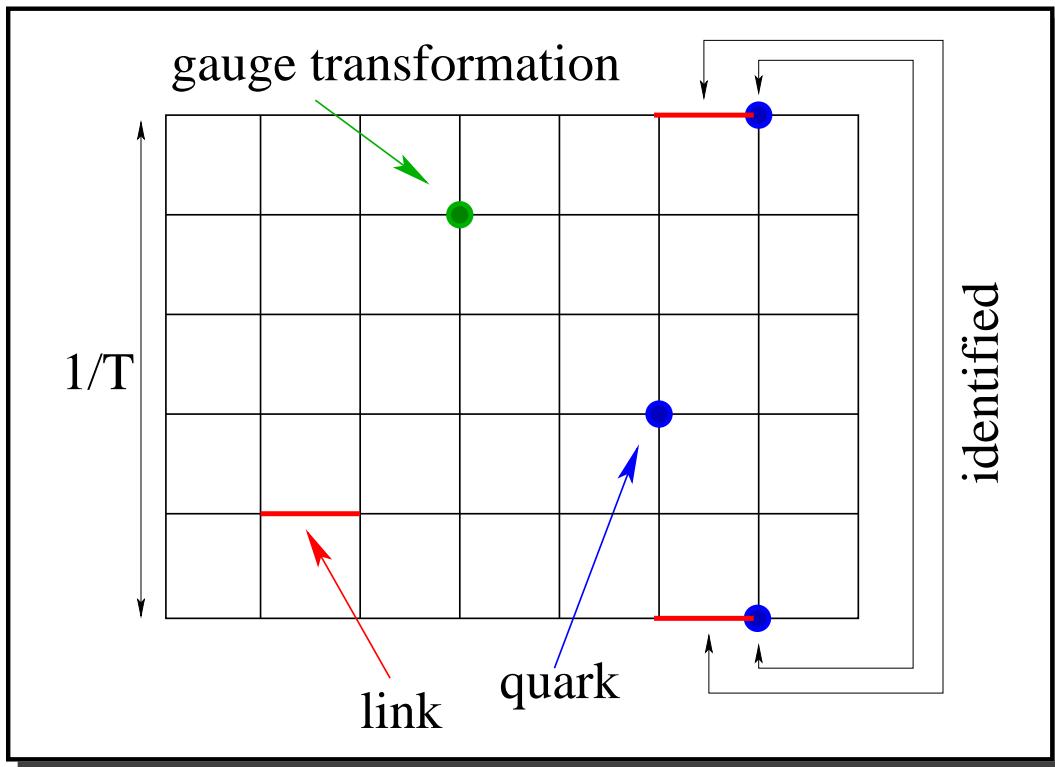
Yang-Mills moduli

- use lattice gauge theory throughout

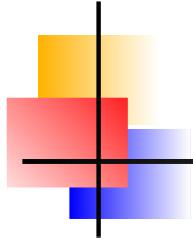


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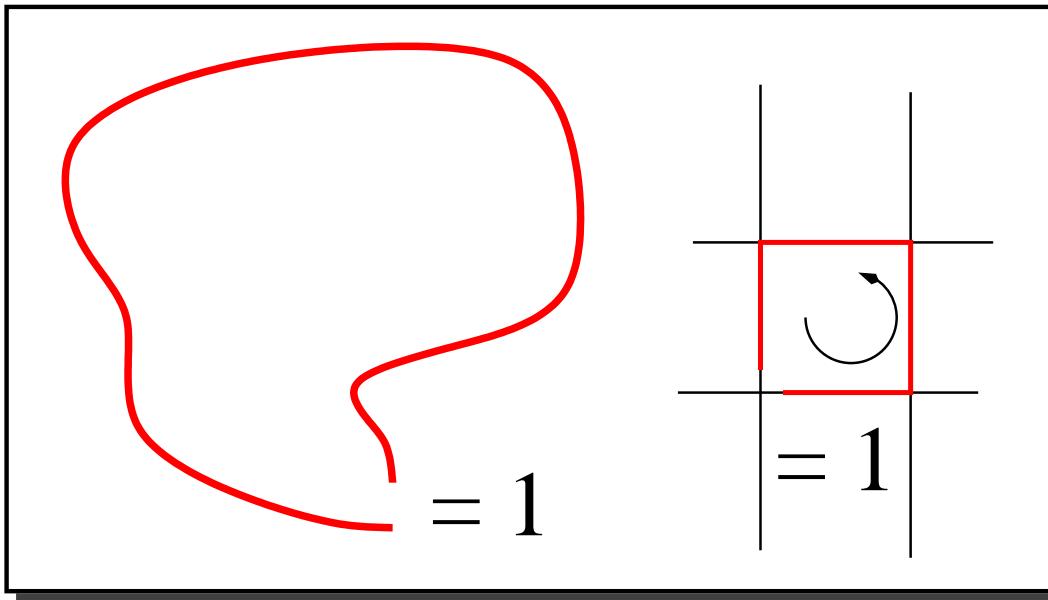


- gauge fields: links U
- matter fields: site q
- gauge transformations: site Ω



Yang-Mills moduli

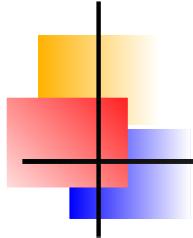
- My name is vacuum - the vacuum:
(pert.) vacuum \leftrightarrow all contractable loops are 1



example:

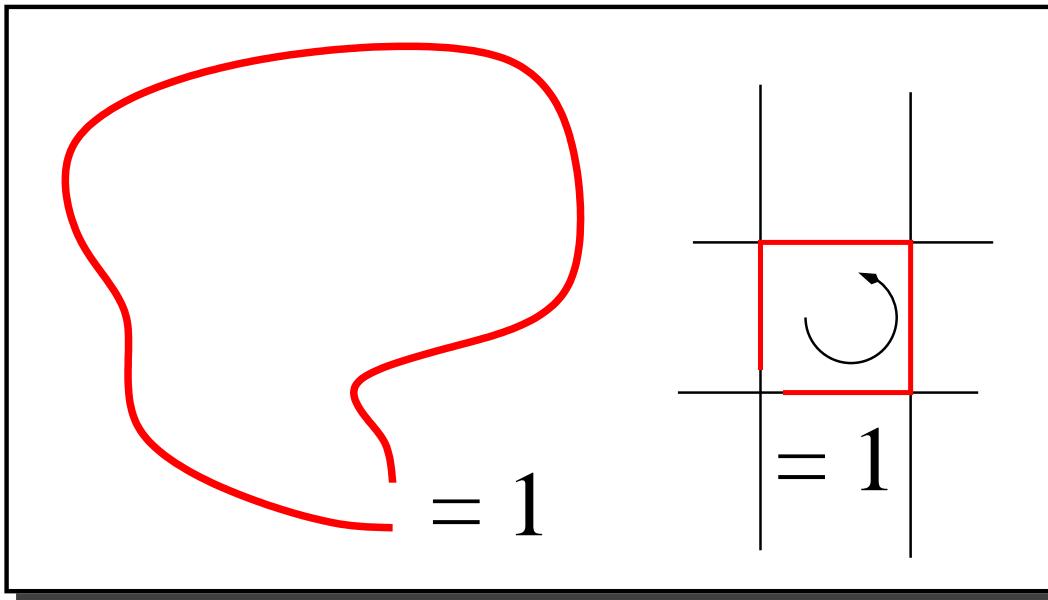
$$U_\mu(x) = 1$$

more vacua?



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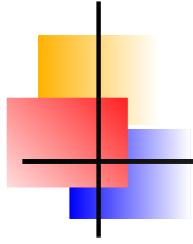


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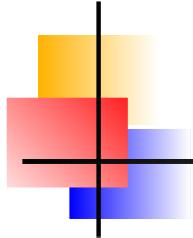
more vacua?

- constructing the moduli space
 - ⇒ need to “devide out” the gauge transformations
 - [Keurentjes, Rosly, Smilga, PRD 58 (1998) 081701]
 - [Schaden, PRD 71 (2005) 105012]
 - [Langfeld, Lages, Reinhardt, PoS LAT2005:201, 2006.]



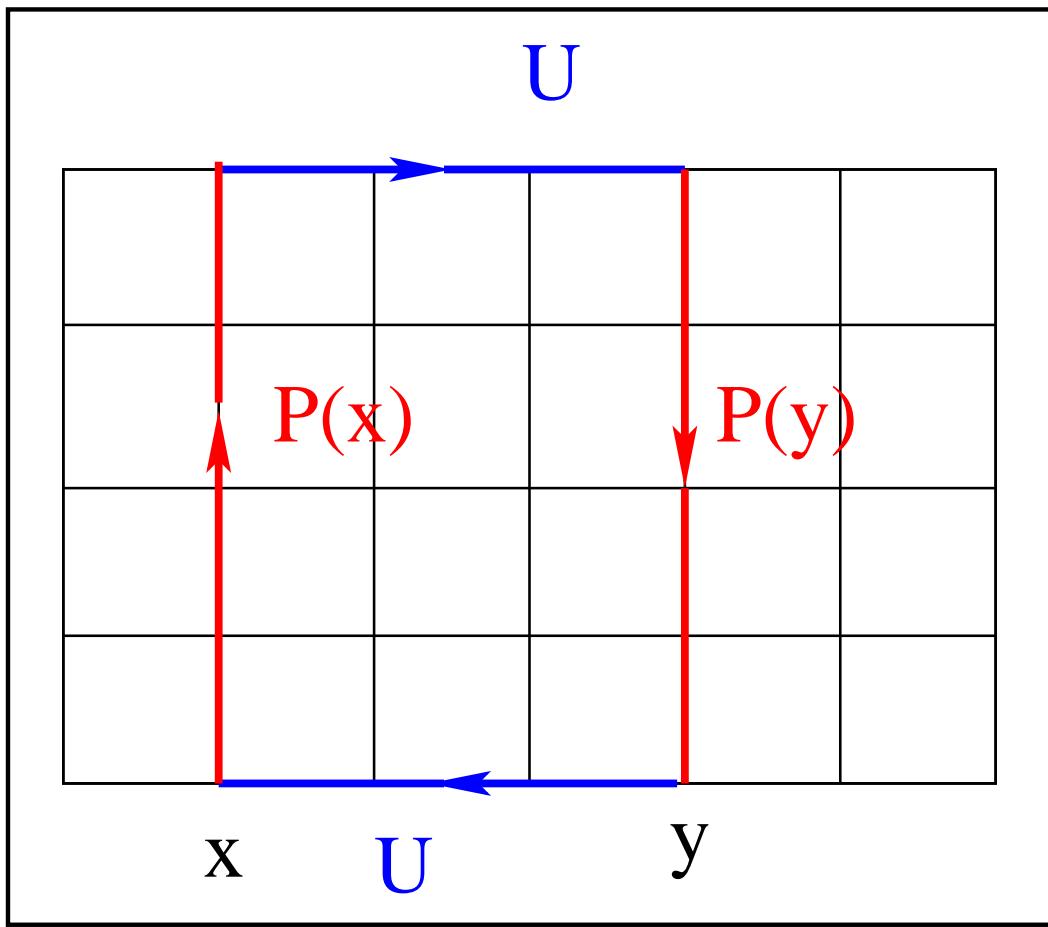
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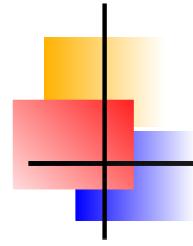
- vacuum $\Rightarrow \text{tr } P(x) = \text{tr } P(y), \quad P$: Polyakov line



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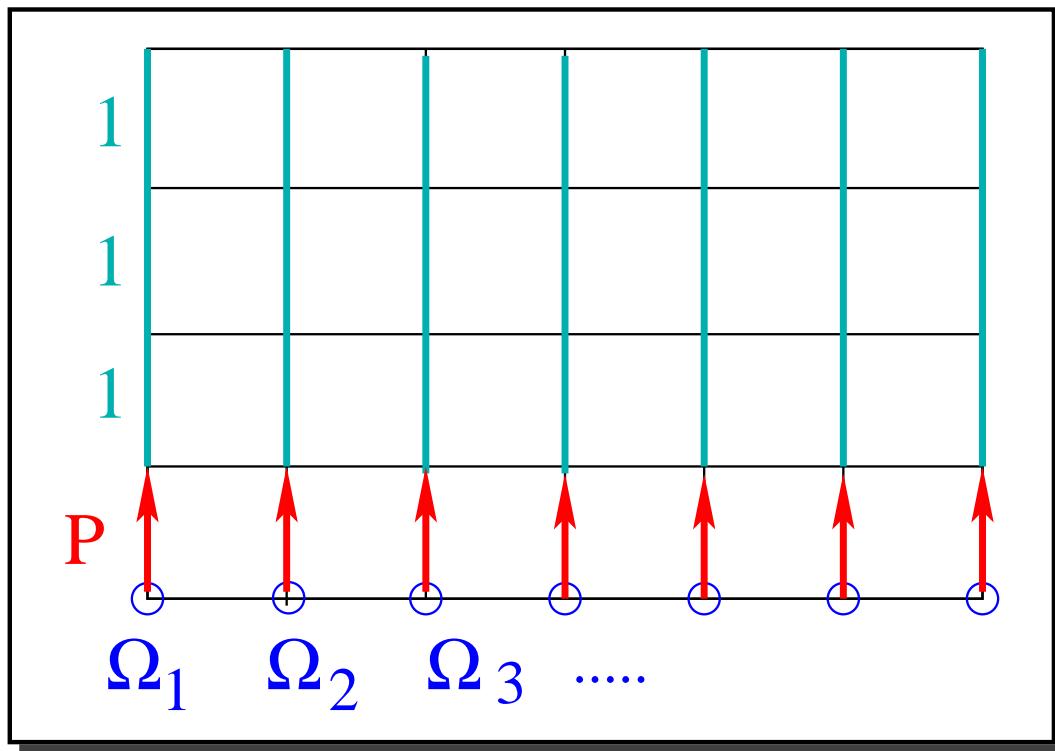
- vacuum $\Rightarrow \text{tr } P(x) = \text{tr } P(y)$, P : Polyakov line
- $U^\dagger P(x) U P^\dagger(y) = 1 \Rightarrow P(y) = U^\dagger P(x) U$ q.e.d.

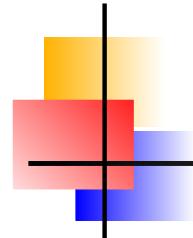




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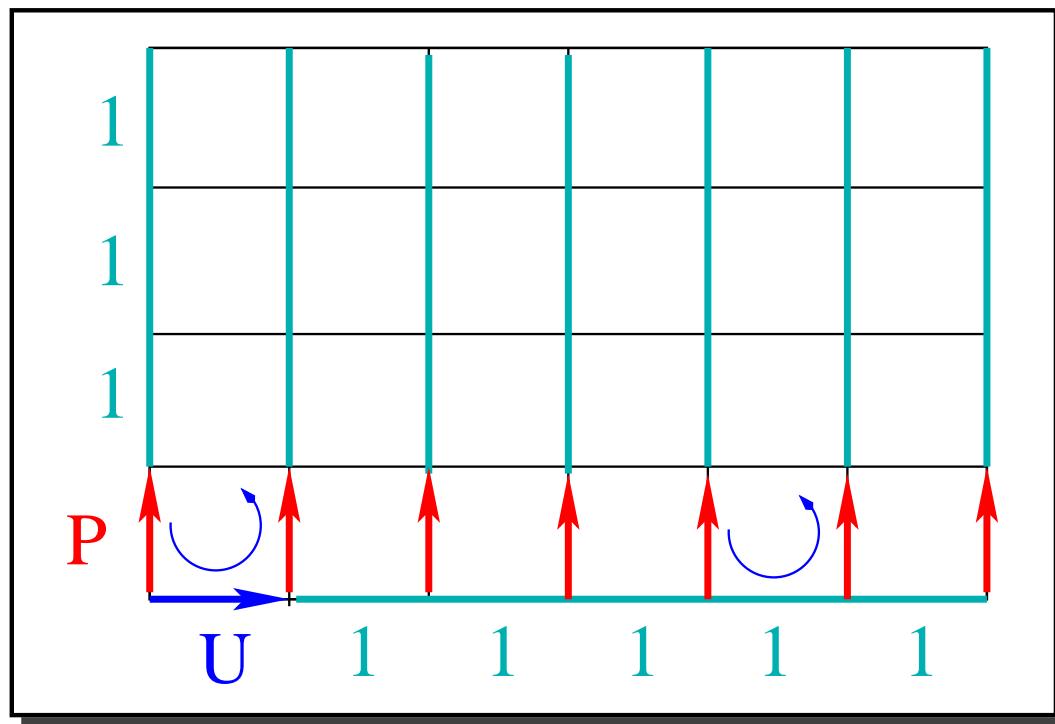
- complete gauge fixing \longrightarrow moduli space
step 1

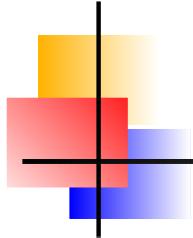




Yang-Mills moduli

- complete gauge fixing \longrightarrow moduli space
step 2

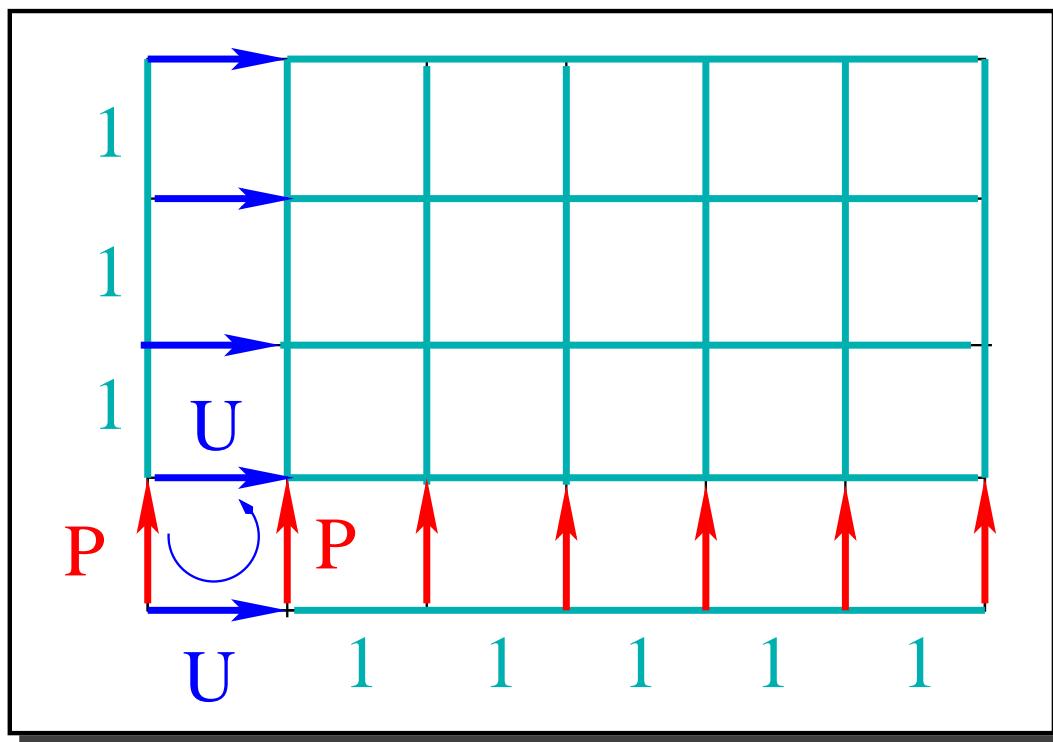


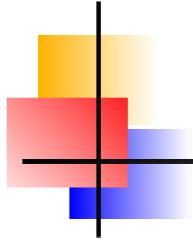


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step 3

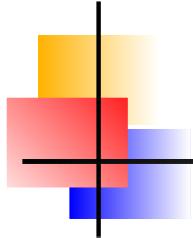
$$U P U^\dagger P^\dagger = 1$$
$$[U, P] = 0$$
$$U, P \in \text{cartan}$$





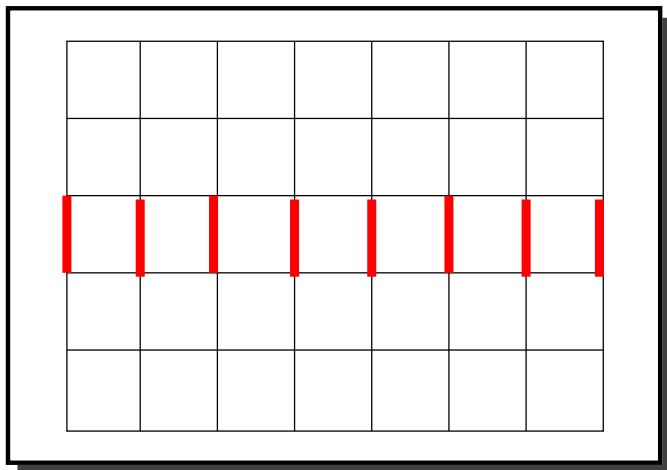
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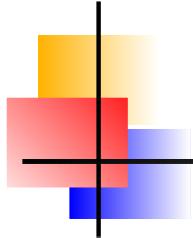
- choose $P_1, P_2 \in \text{Cartan}$ such that $\text{tr } P_1 \neq \text{tr } P_2$
found a variety of **gauge inequivalent** vacua
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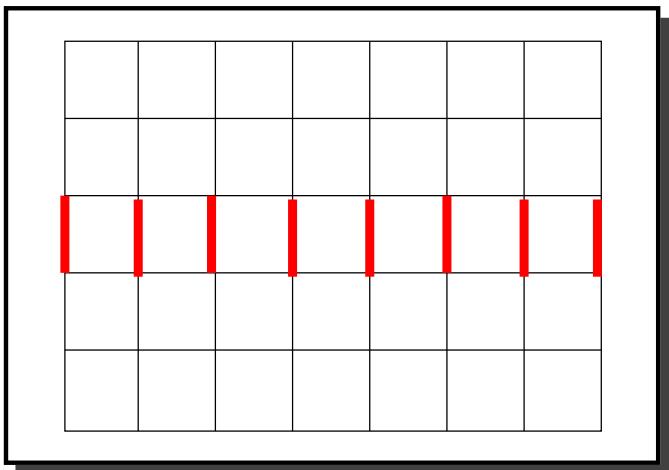
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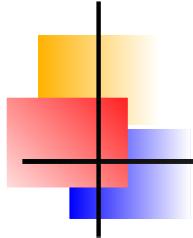


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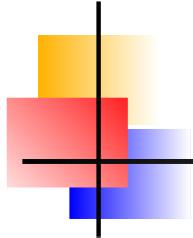
- symmetry of the action
- mediates between vacua:*
 $\text{tr } P \rightarrow z \text{tr } P$



Centre sector tunneling:

Hypothesis:

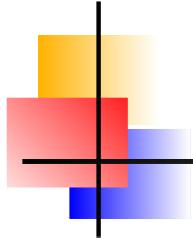
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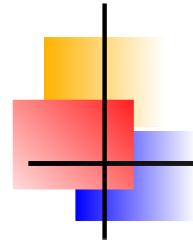
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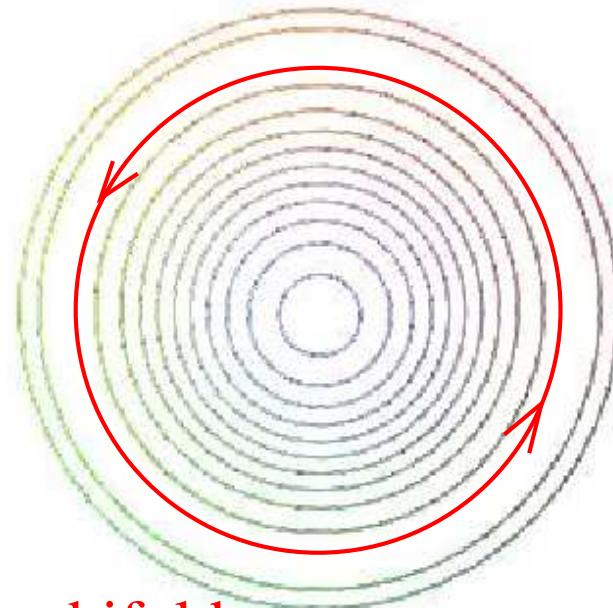
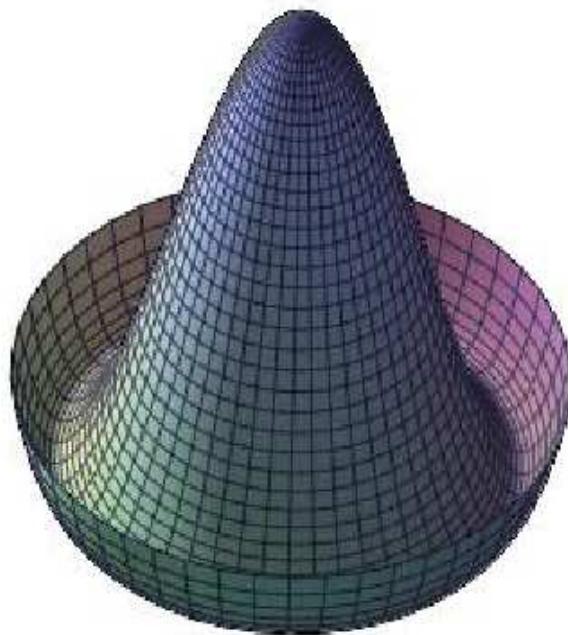
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- dynamical matter (*QCD!*): flat directions of the vacuum are lifted, but centre sector tunneling still takes place in the hadronic phase
extreme conditions: SSB of centre symmetry on top of explicit breaking



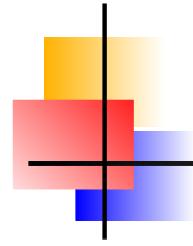
Centre sector tunneling:

Illustration:

- classical action \Rightarrow moduli space



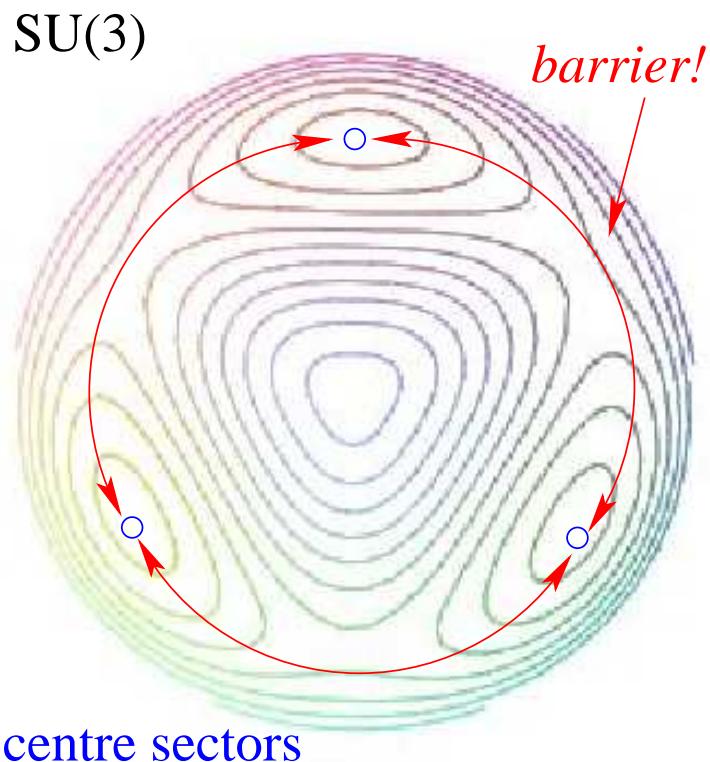
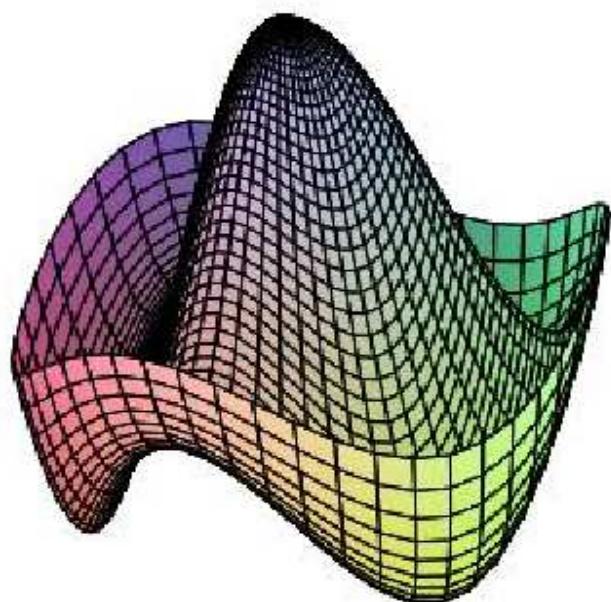
orbifold

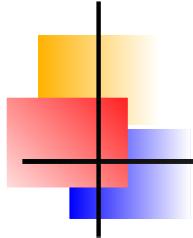


Centre sector tunneling:

Illustration:

- quantum effective action, pure YM \Rightarrow centre sectors

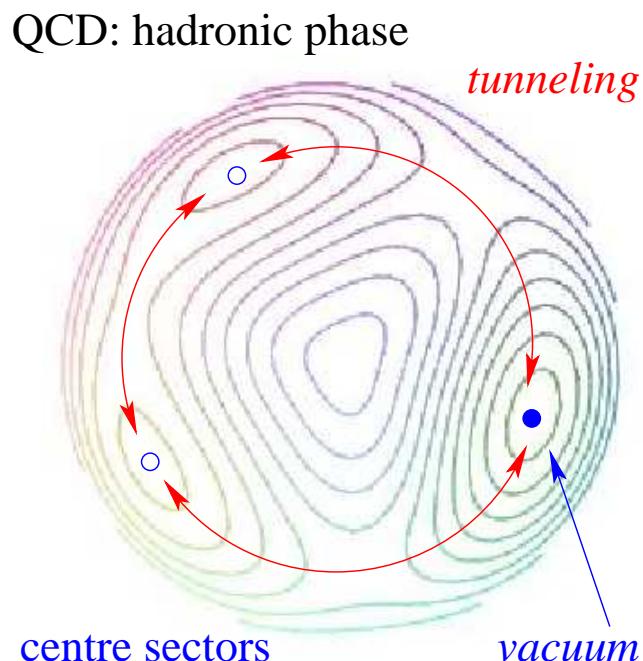


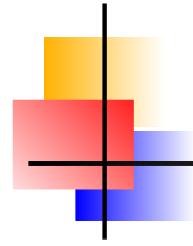


Centre sector tunneling:

Illustration:

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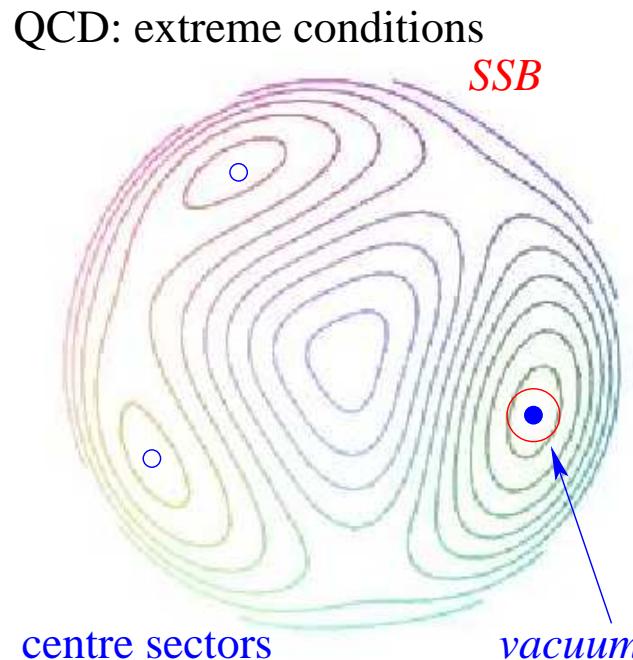
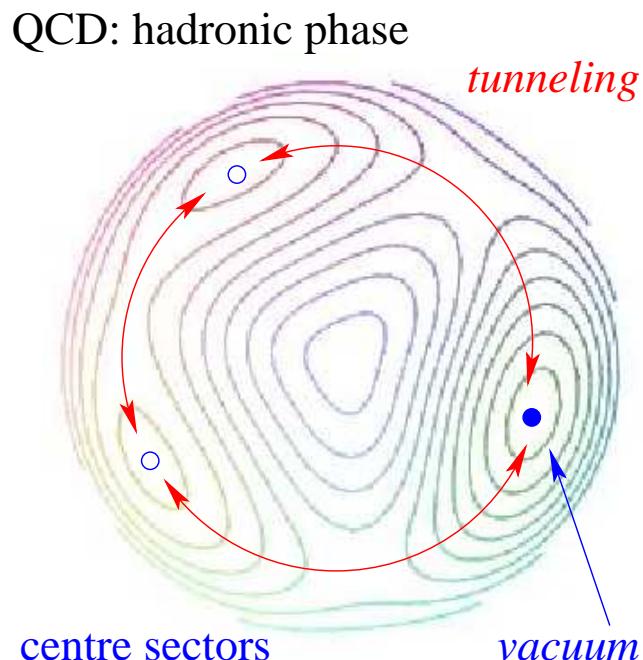


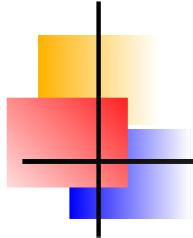


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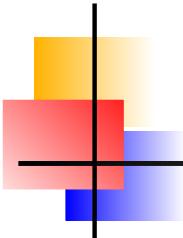
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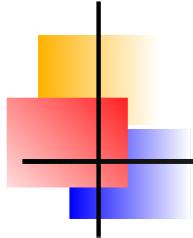
Centre sector tunneling:

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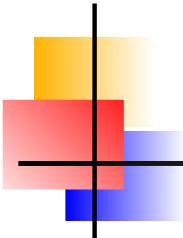
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at finite **chemical potential μ**



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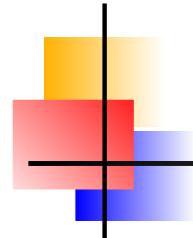
⇒ Fermi-Einstein-Condensation (FEC) ←



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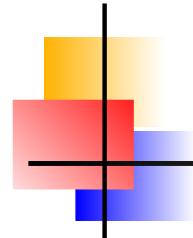
\Rightarrow Fermi-Einstein-Condensation (FEC) \Leftarrow
- ...will talk about $SU(3)$ + matter = QCD later



Fermi-Einstein-Condensation (FEC)

Model consideration:

- $q(x)$: quarks, m : mass, μ : chemical potential
- A_m : moduli fields \Rightarrow weighted sum over centre sectors



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- $q(x)$: quarks, m : mass, μ : chemical potential
- A_m : moduli fields \Rightarrow weighted sum over centre sectors
- partition function: $\exp\{iA_m\} = Z_m \in Z(N_C)$

$$Z = \sum_{m=1}^{N_c} p_m \int \mathcal{D}q \mathcal{D}\bar{q} \exp\{\bar{q}(i\partial + (A_m + i\mu)\gamma_0 + im)q\}$$

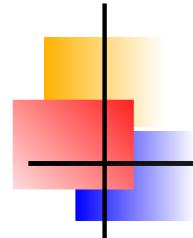
p_m : probability for centre sector m

pure YM-theory: $p_m = 1/N_c, \forall m$

high T SSB phase: $p_{N_c} = 1, p_m = 0$ for $m = 1 \dots N_c - 1$

hadronic phase: $p_{N_c} > p_m \neq 0$ for $m = 1 \dots N_c - 1$

[Langfeld, Welleghausen, Wipf, Phys. Rev. D81 (2010) 114502]

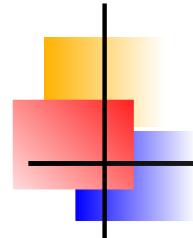


Fermi-Einstein-Condensation (FEC)

Results for baryon density:

- $$B = \frac{1}{\pi^2} \int_m^\infty dE E \sqrt{E^2 - m^2} \rho(E, T, \mu)$$

$\rho(E, T, \mu)$: density of states



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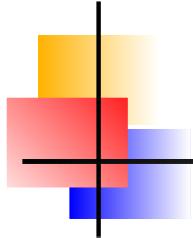
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z_m : centre phases, w_m : weights



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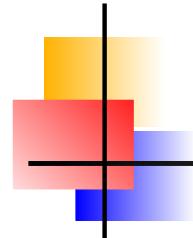
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z_m : centre phases, w_m : weights

- $$w_m = p_m \rho_m / \sum_i p_i \rho_i$$

$$\rho_i = \exp \left\{ \frac{V}{\pi^2} \int_m^\infty dE E \sqrt{E^2 - m^2} \ln(1 + z_i e^{-\frac{E-\mu}{T}}) \right\}$$



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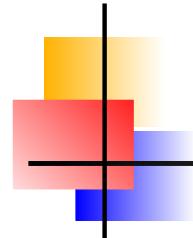
(I) high temperature phase

- remember:

sector probability: $p_{N_c} = 1, p_m = 0$ for $m = 1 \dots N_c - 1$

centre element: $z_{N_c} = 1$

weights: $w_{N_c} = \rho_{N_c}/\rho_{N_c} = 1, w_m = 0$ else

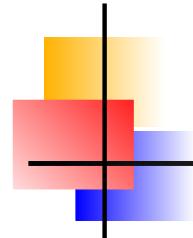


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- $$\rho(E, T, \mu) = \sum_m \frac{z_m}{e^{[E-\mu]/T} + z_m} w_m = \frac{1}{e^{[E-\mu]/T} + 1}$$

free quarks with a *Fermi surface !!*



Fermi-Einstein-Condensation (FEC)

(II) hadronic phase (N_c even)

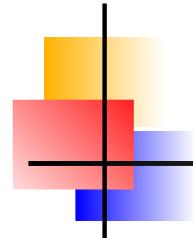
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centre element: $z_{N_c/2} = \exp\{i \frac{2\pi}{N_c} \frac{N_c}{2}\} = -1$

weights: $w_{N_c/2} \approx \frac{p_{N_c/2} \rho_{N_c/2}}{p_{N_c/2} \rho_{N_c/2}} = 1, w_m \approx 0$ else

$$\rho_{N_c/2} = \exp\left\{\frac{V}{\pi^2} \int_m^\infty dE E \sqrt{E^2 - m^2} \ln(1 - e^{-\frac{E-\mu}{T}})\right\} \rightarrow \infty$$

Cooper instability familiar from BEC !!



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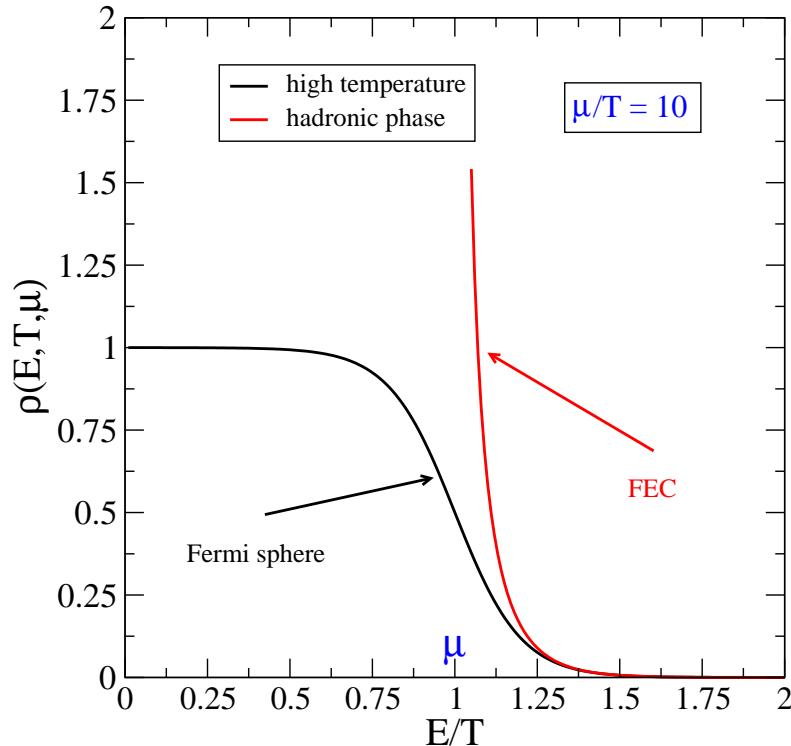
Cooper instability familiar from BEC !!

- $\rho(E, T, \mu) = \sum_m \frac{z_m}{e^{[E-\mu]/T} + z_m} w_m = \frac{-1}{e^{[E-\mu]/T} - 1}$

Fermi-Einstein Condensation (FEC) !!

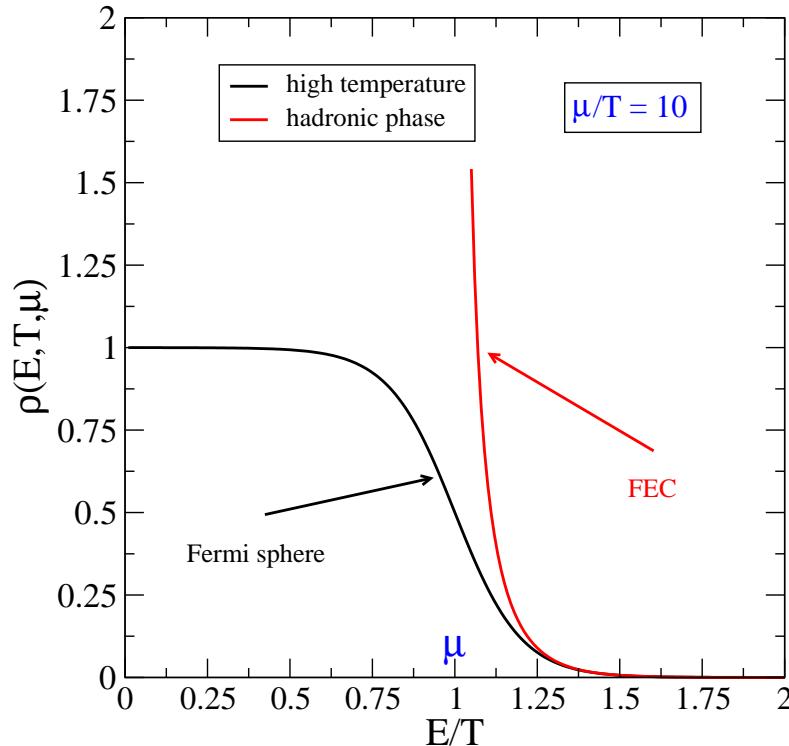
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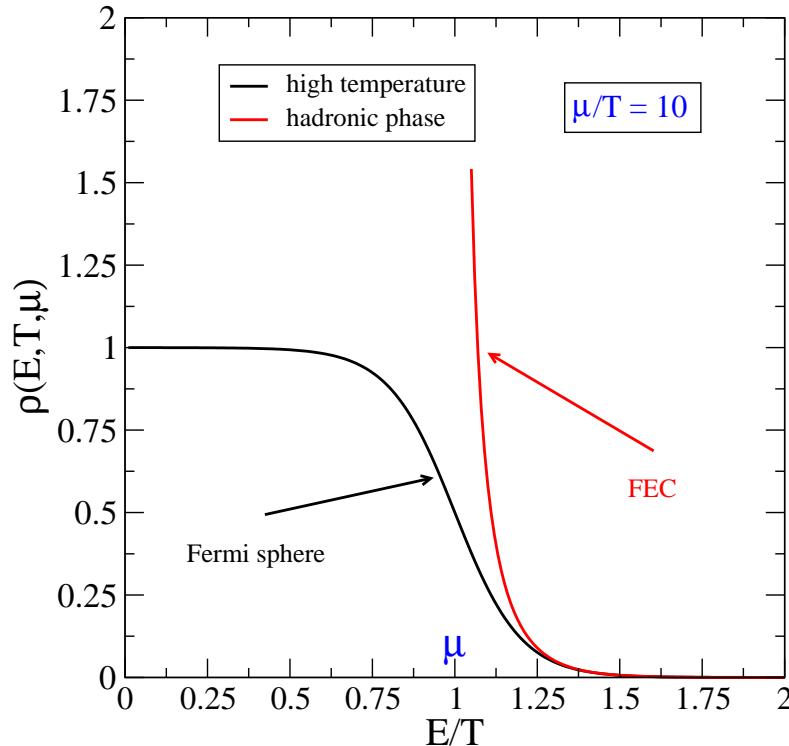
Fermi-Einstein-Condensation (FEC)



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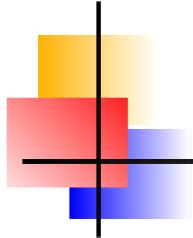
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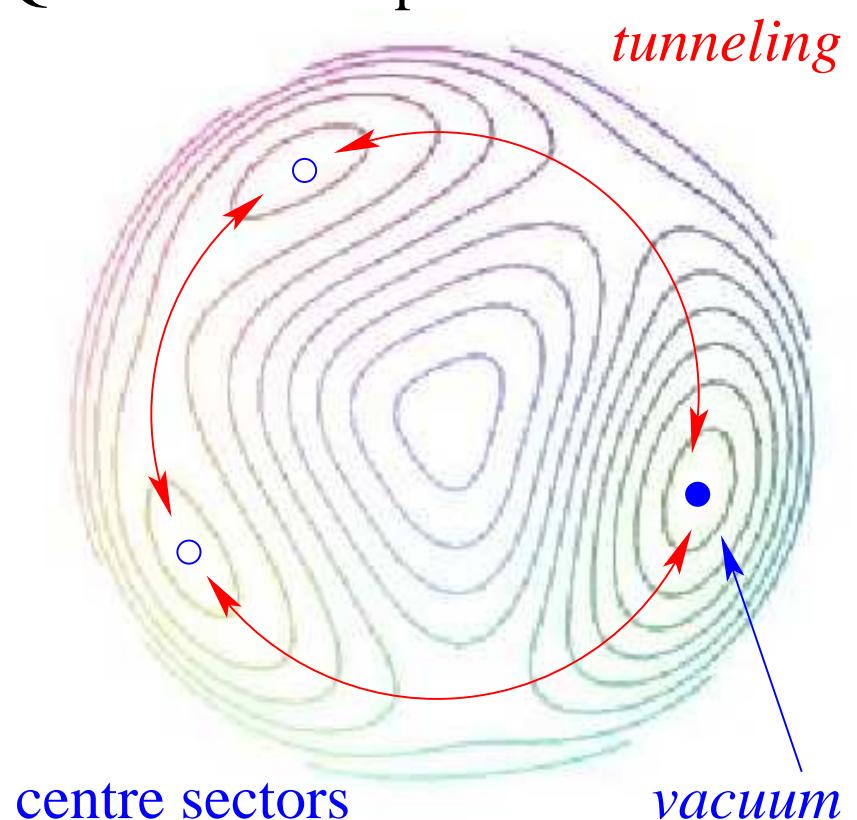
- centre dressed quarks acquire Bose statistic and condense because of a Cooper instability
- quarks are still represented by Grassmann fields but the spin-statistic theorem does not apply as long as colour is confined



Centre sector tunneling:

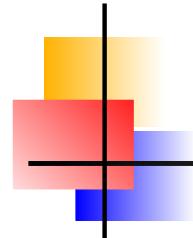
- Does centre sector tunneling take place in the hadronic phase of QCD ?

QCD: hadronic phase



centre sectors

vacuum



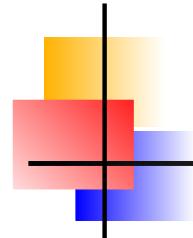
Model considerations:

The SU(2) - qHiggs model

- Degrees of freedom:

gluons $\longrightarrow U_\mu(x)$

Higgs (scalar, fundamental rep.) $\longrightarrow \phi(x)$



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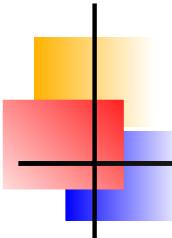
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- Coupling constants:

bare gauge coupling $\longrightarrow g \leftrightarrow \beta$

Higgs mass $\longrightarrow \kappa$: mass $\propto 1/\kappa$

NO Higgs quartic coupling (quadratic Higgs)



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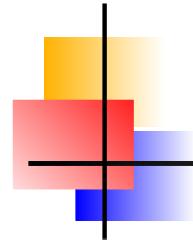
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- need good ergodicity:

1HMC for the gluon sector

HMC for the Higgs sector

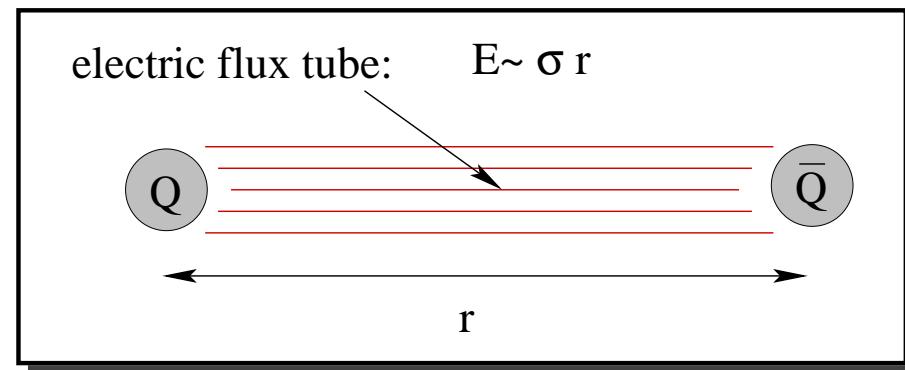
simulations at the HPCC, Plymouth



Model consideration:

String breaking

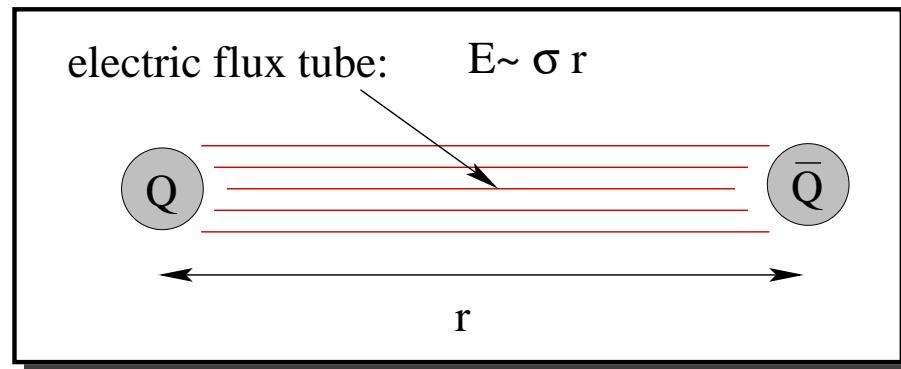
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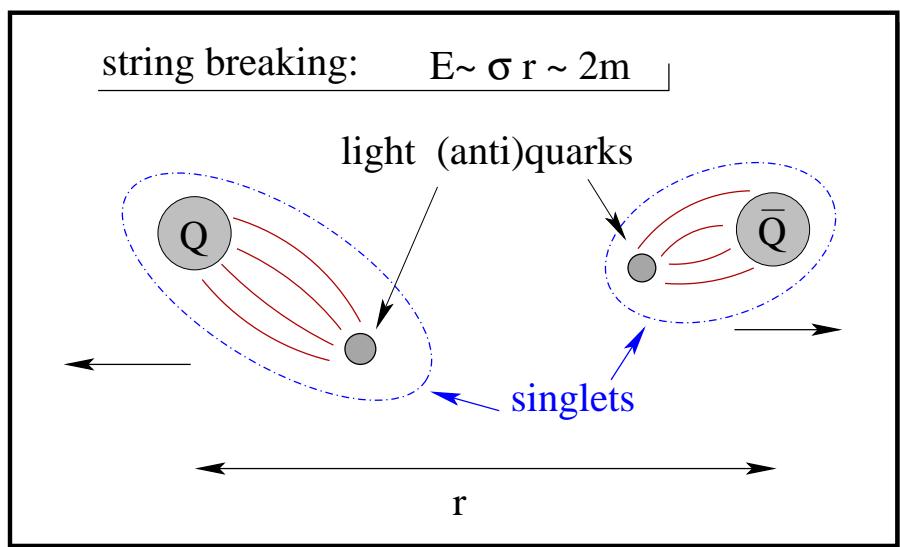
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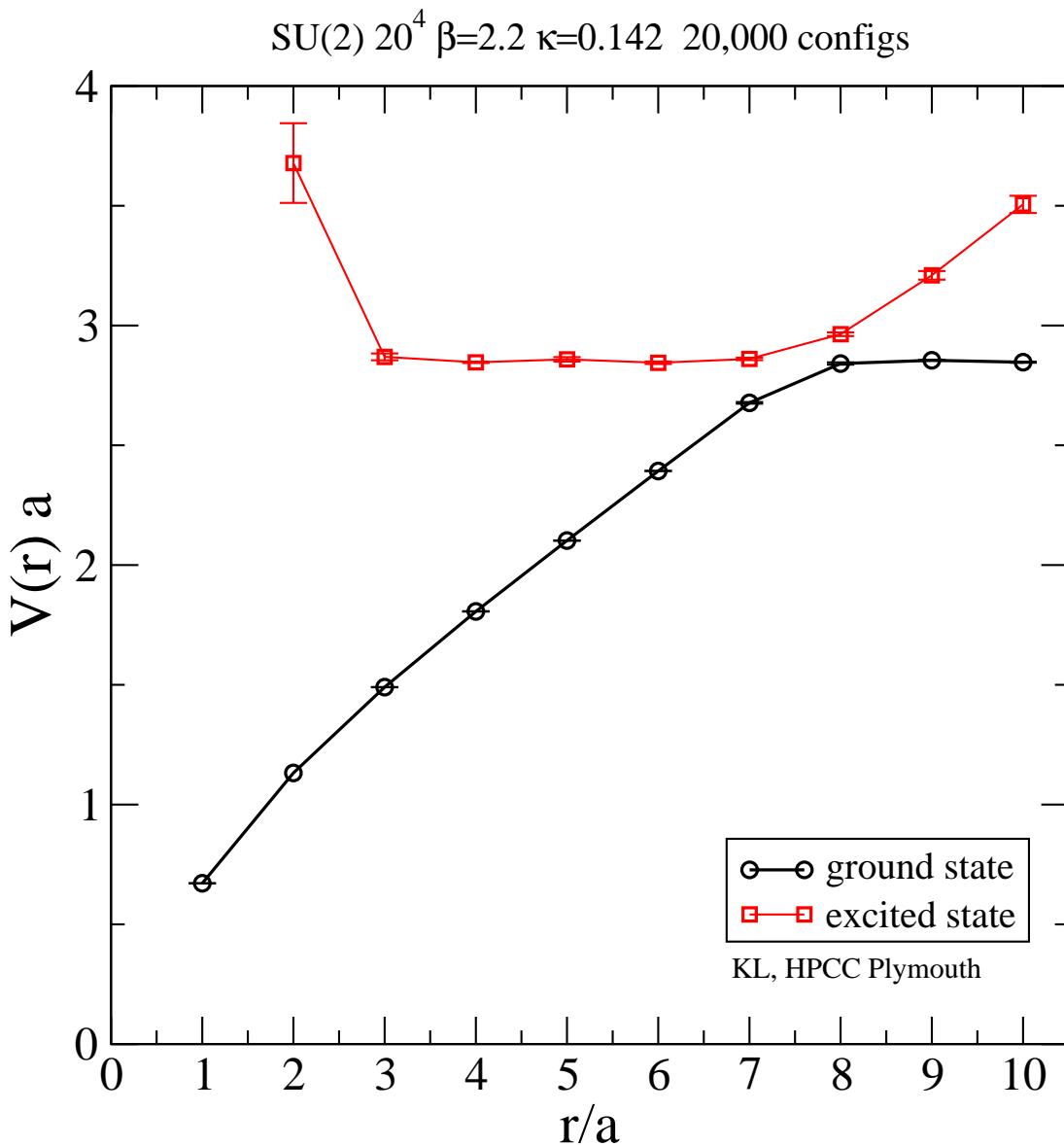
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- QCD (with dynamical quarks)

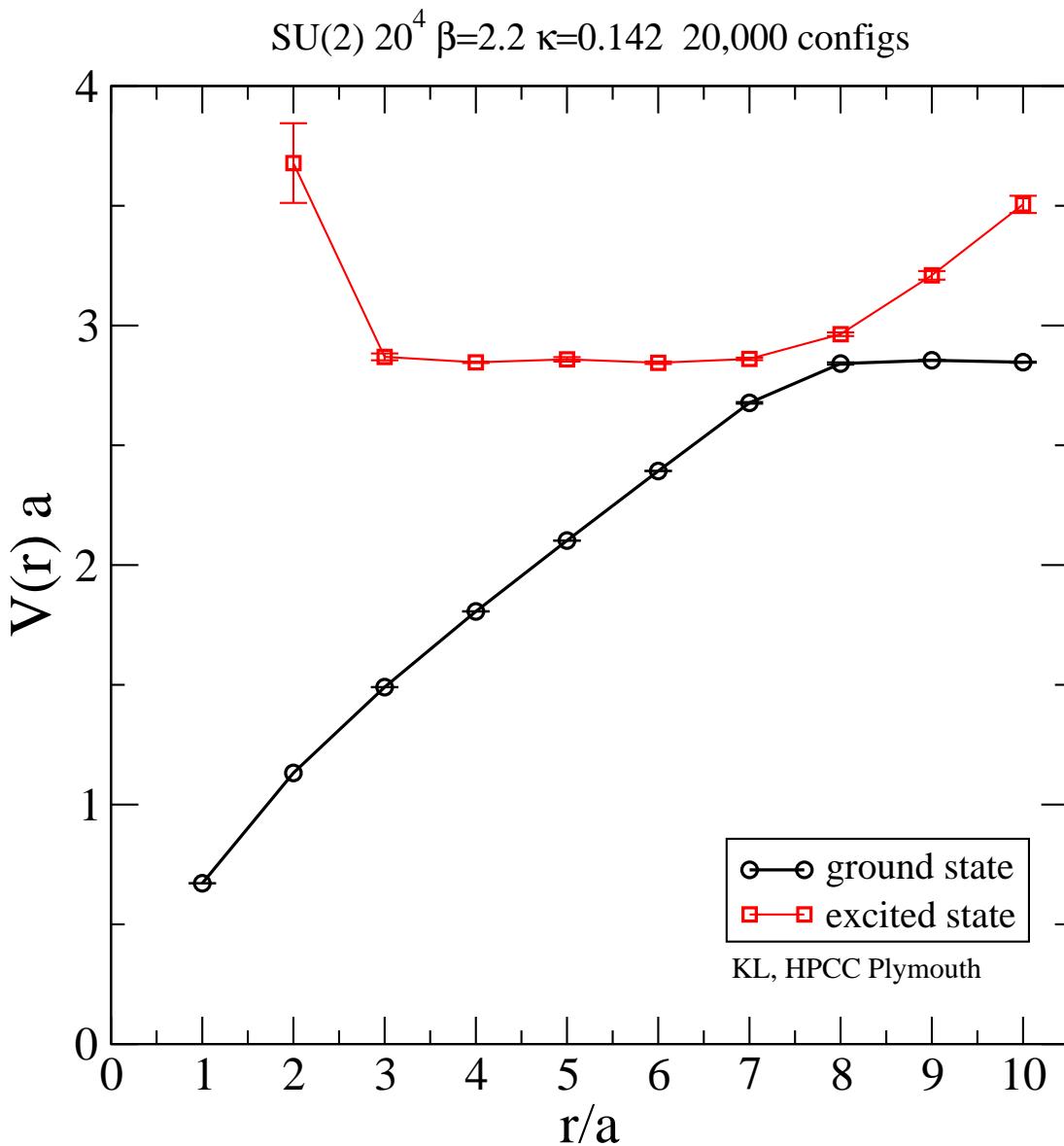


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carefully tune the
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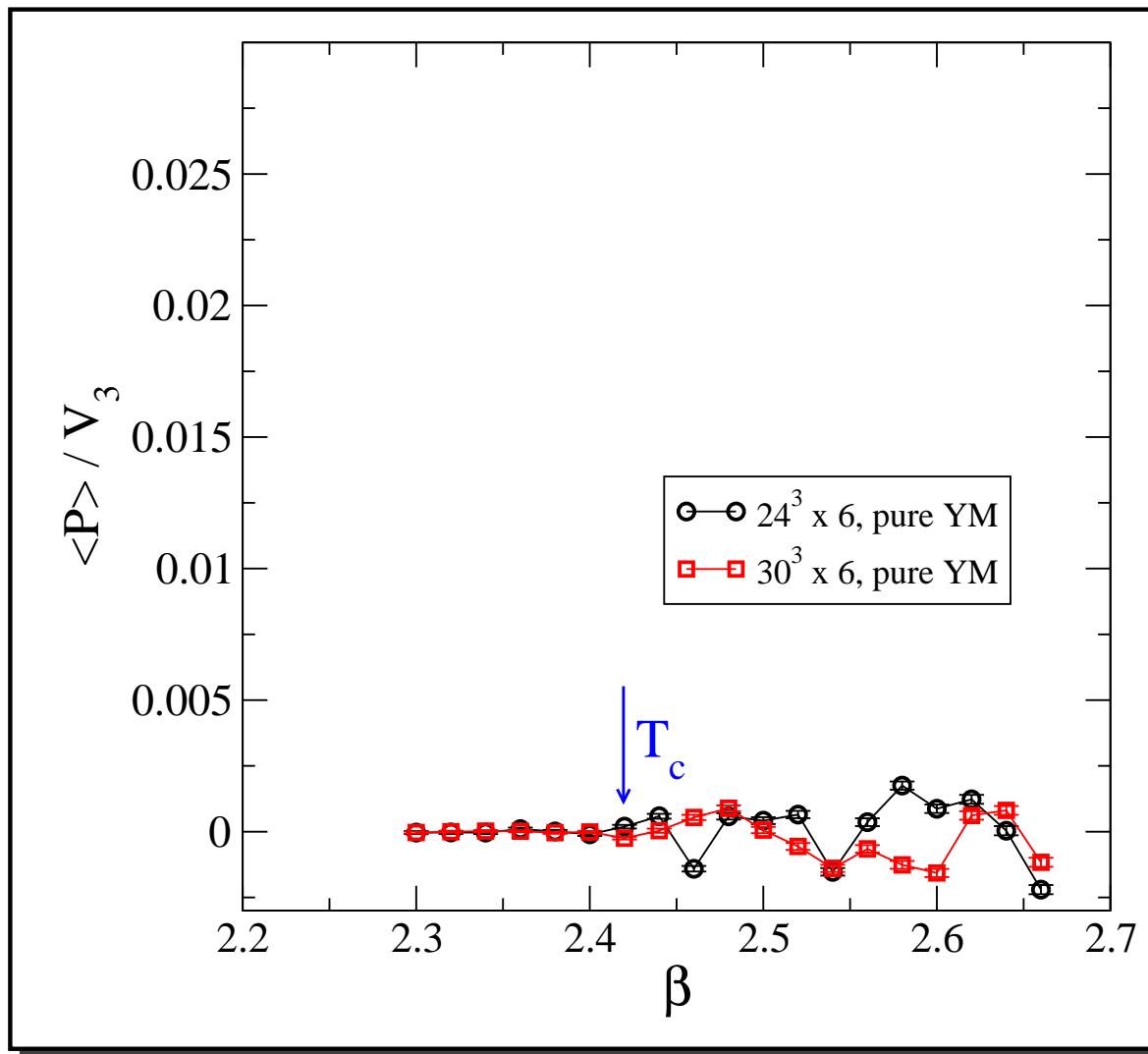
SU(2) + qHiggs:

carefully tune the
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shows string breaking
breaks center symmetry
abundance of configs

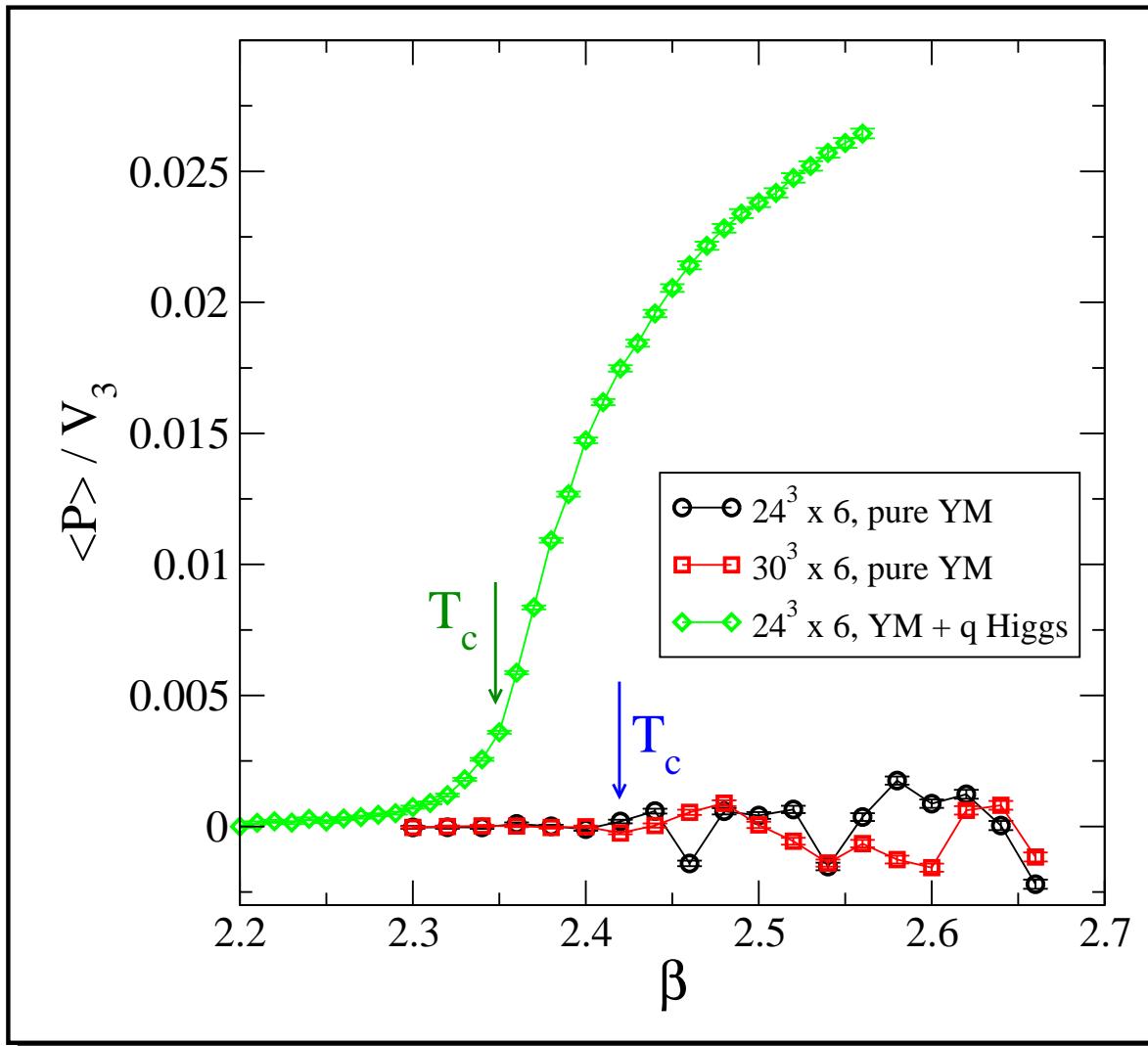
Centre sector tunneling: YM + qHiggs

- Polyakov line: $P = \left\langle \sum_{\vec{x}} \prod_t U_0(\vec{x}, t) \right\rangle$

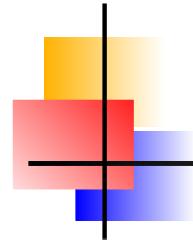


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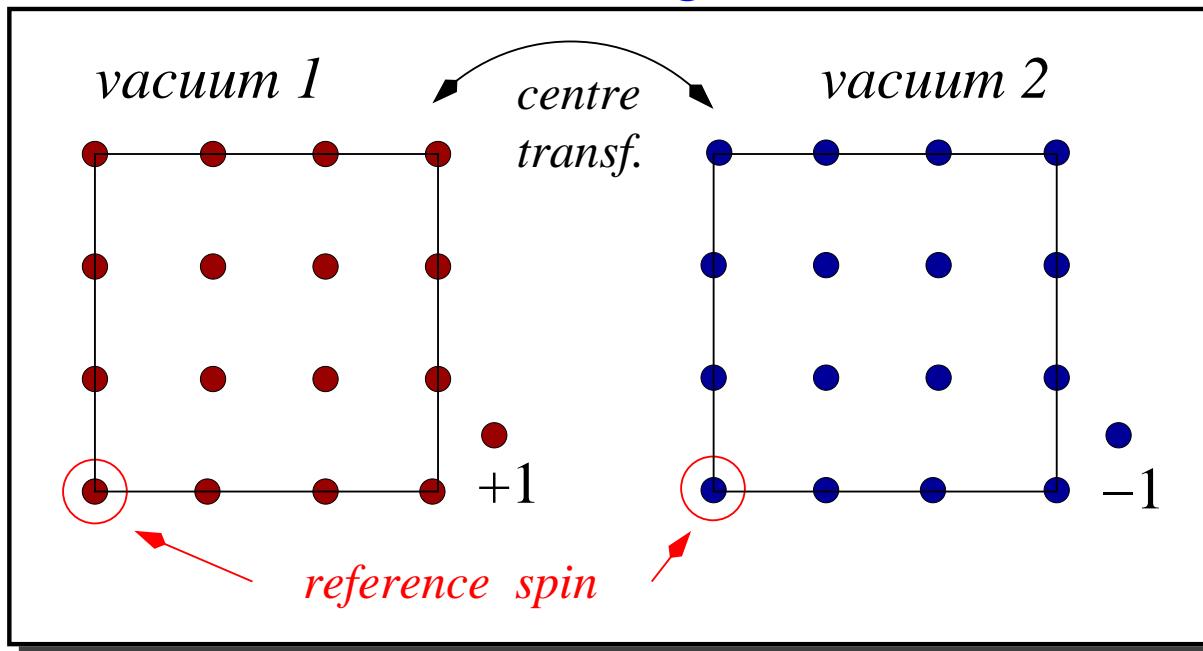


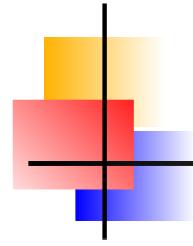
shift of T_c :
 $300 MeV \longrightarrow 170 MeV$



Centre sector tunneling: YM + qHiggs

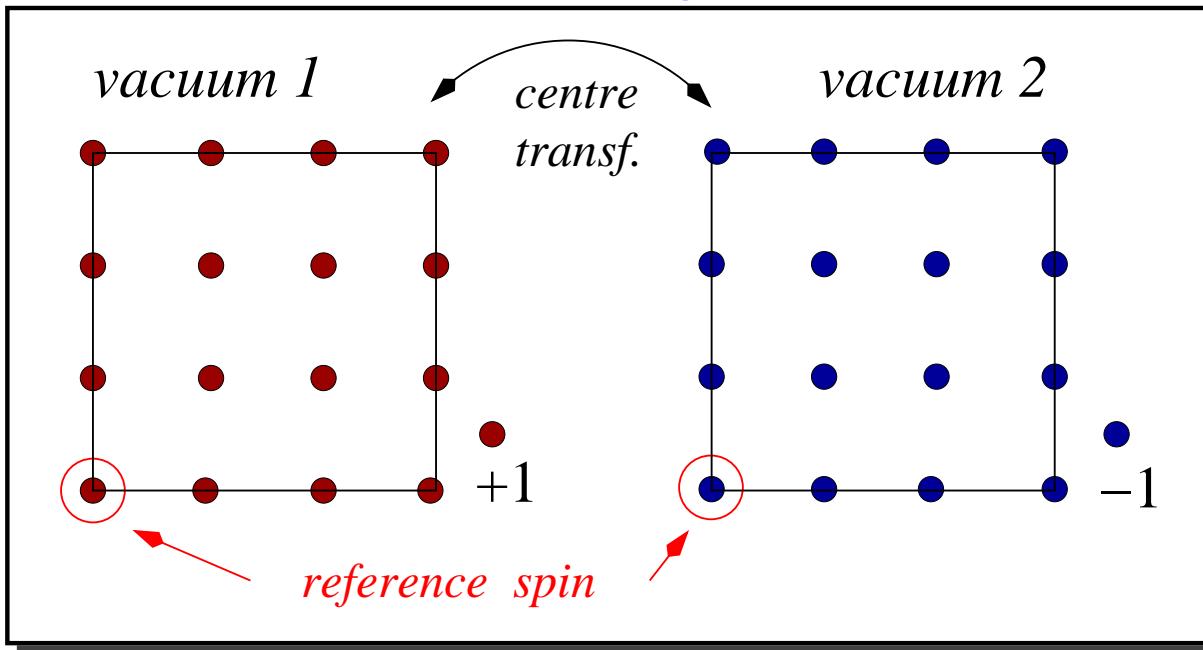
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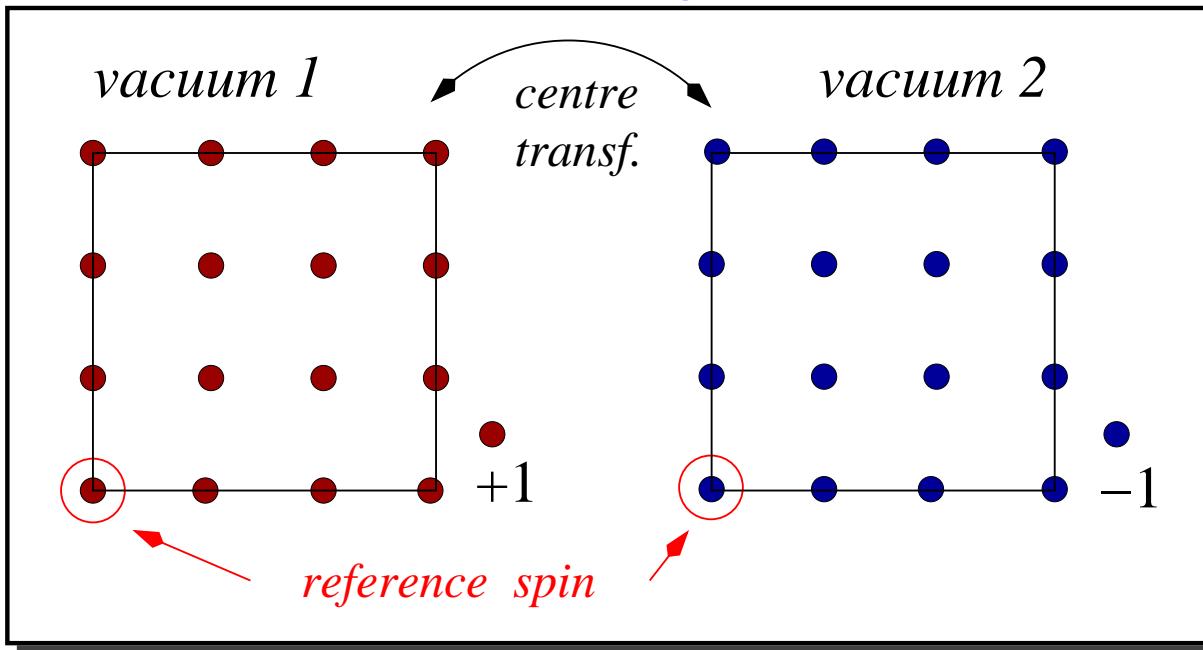
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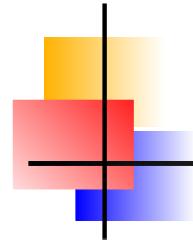
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- magnetisation: $\langle M \rangle = 0$ (Swendsen-Wang cluster alg)
- use: $\left\langle M(0) \sum_x M(x) \right\rangle \propto \begin{cases} \xi^3 = \text{finite, symm. phase} \\ V_3 \text{ SSB phase} \end{cases}$

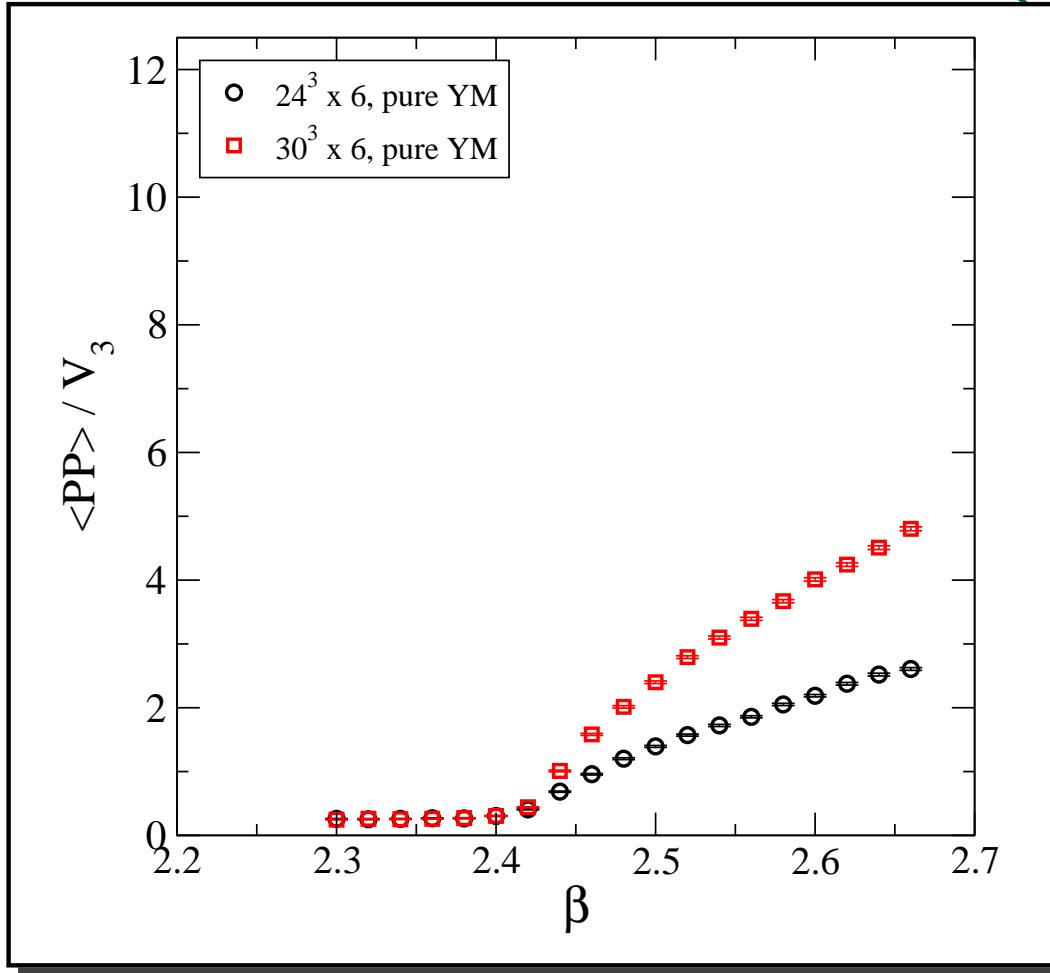


Centre sector tunneling: YM + qHiggs

- Polyakov line: $\left\langle p(0) \sum_x p(x) \right\rangle \propto \begin{cases} \sigma^{-3} & \text{confinement phase} \\ V_3 & \text{high } T \text{ phase} \end{cases}$

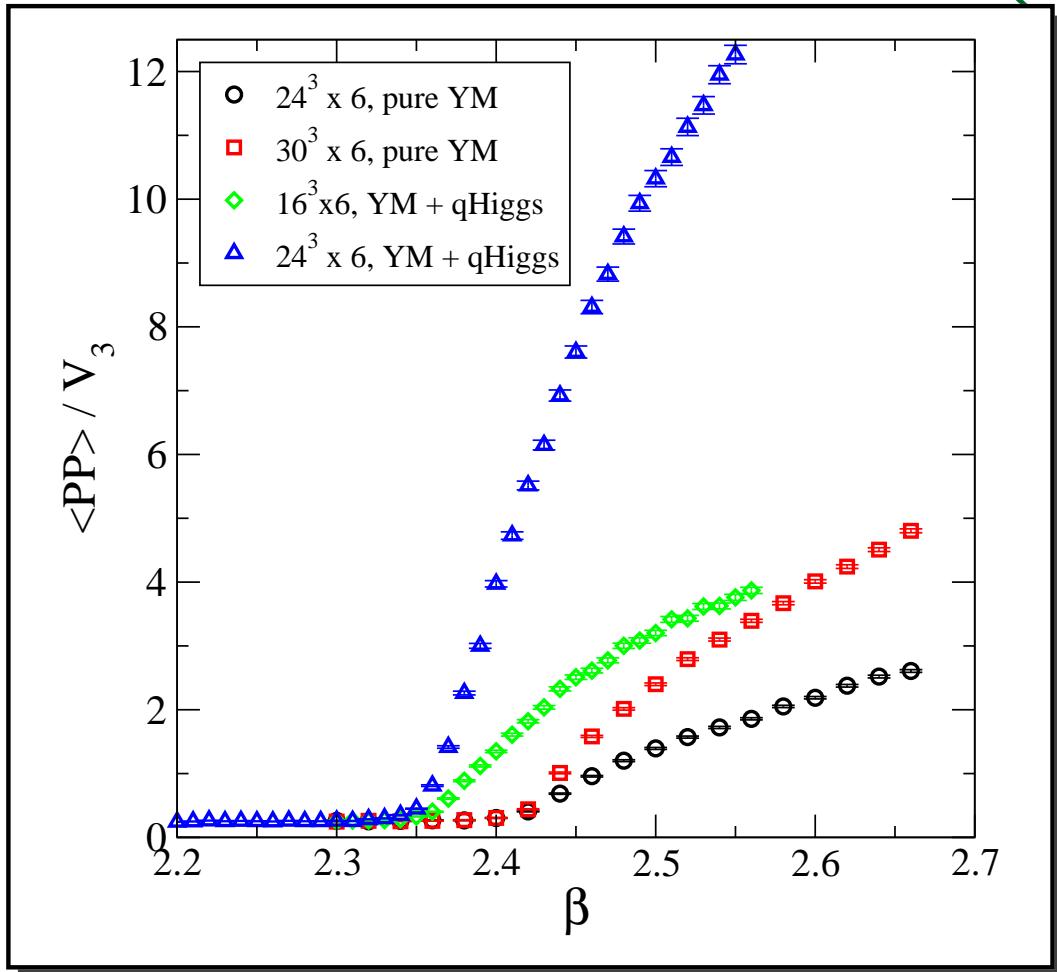
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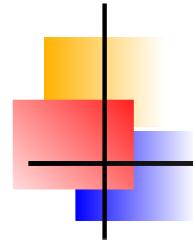


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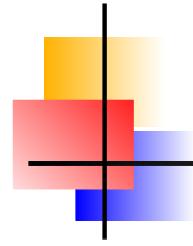


data suggest tunneling
for $\beta < \beta_c$
despite of explicit
centre symm. breaking



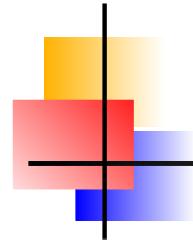
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Wigner Weyl: $\langle\psi|^Z\psi\rangle = 1$
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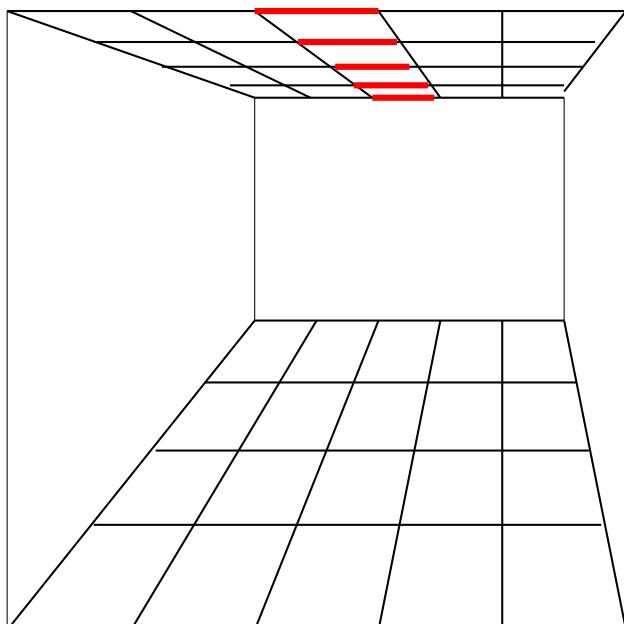
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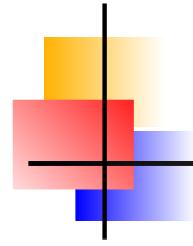
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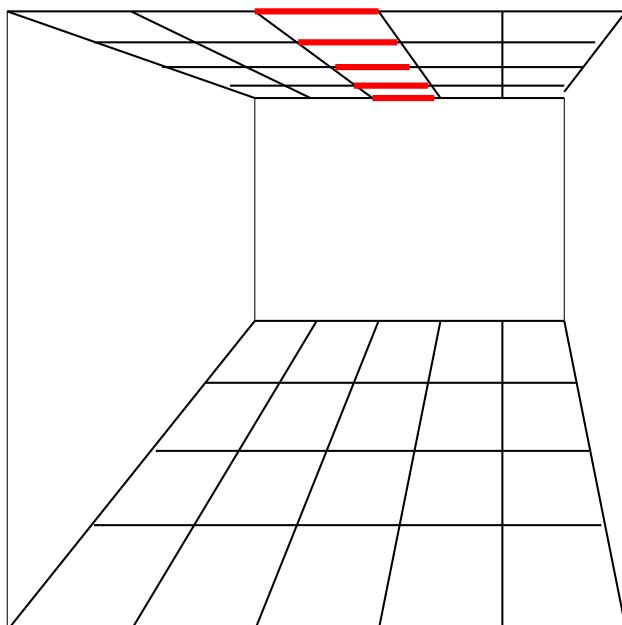
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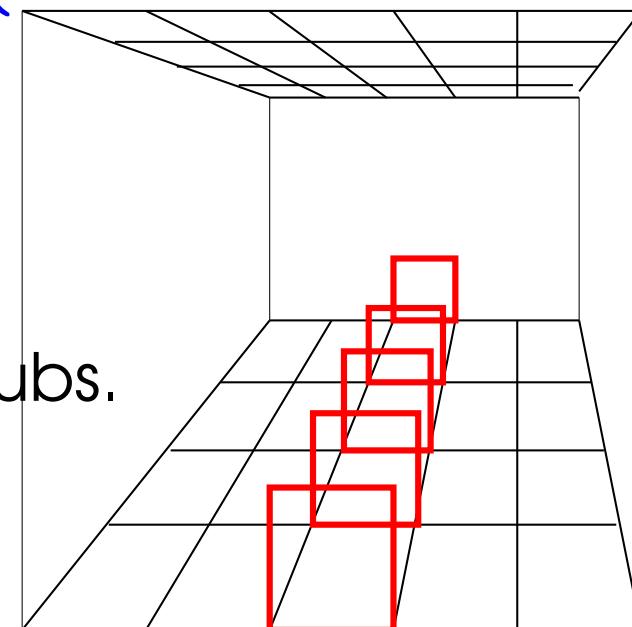
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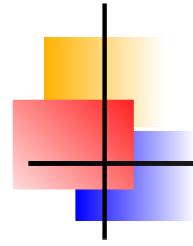
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variable subs.



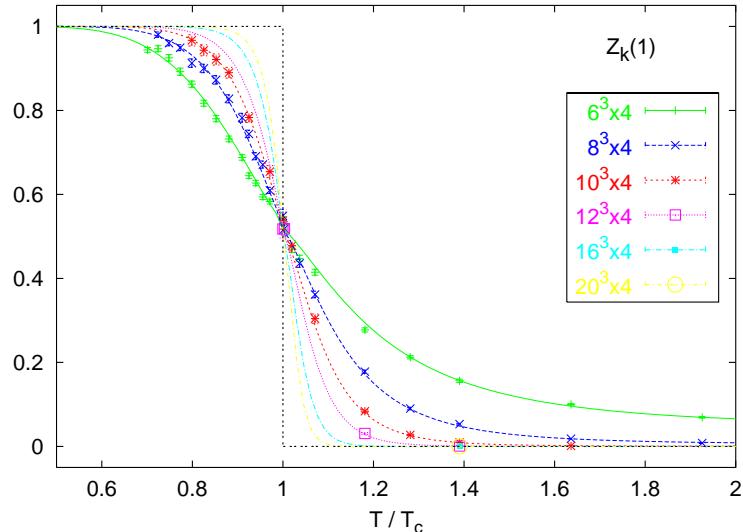
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Centre sector tunneling: 't Hooft loop

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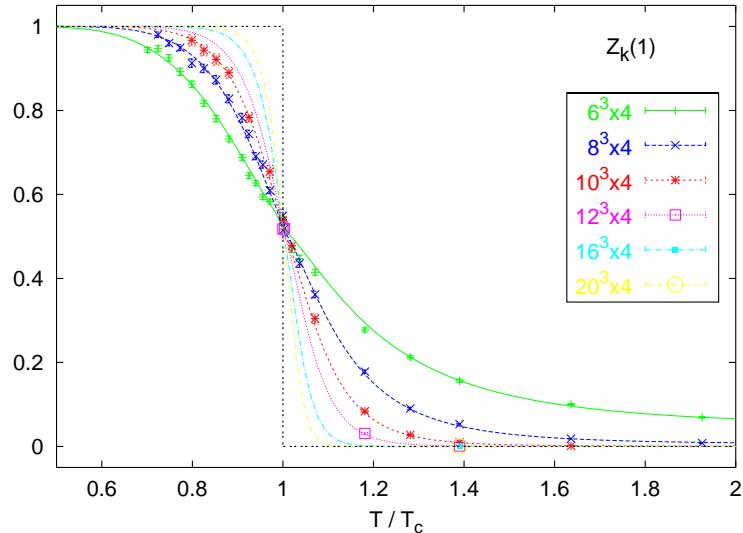
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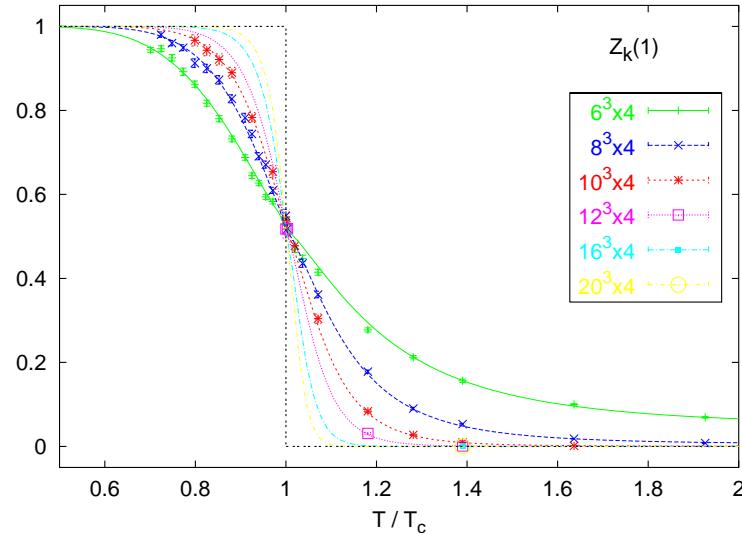
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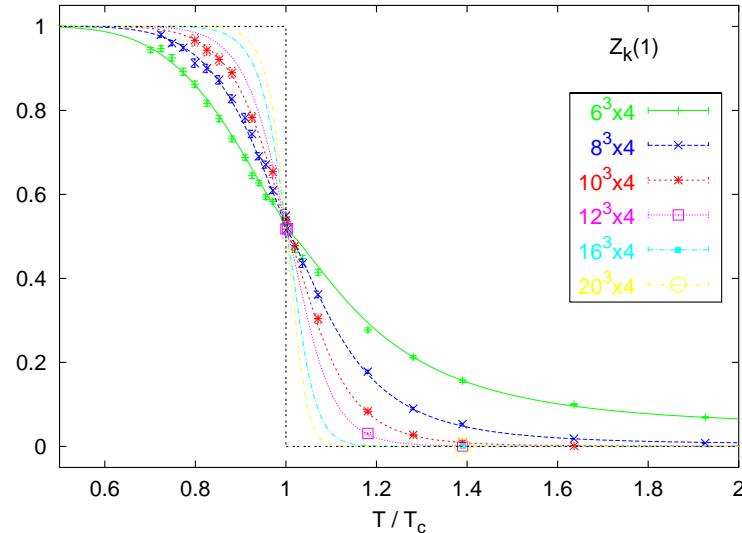
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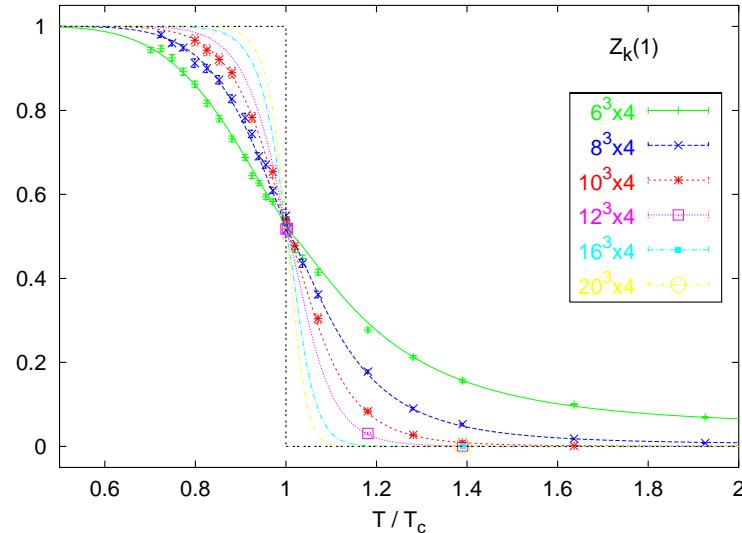
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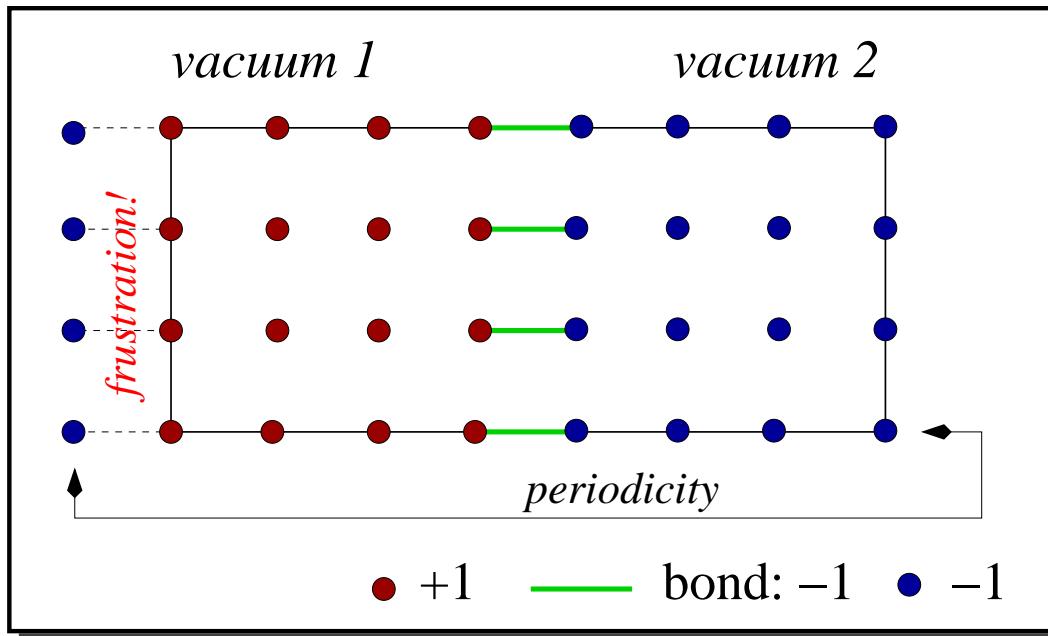
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- need something else....

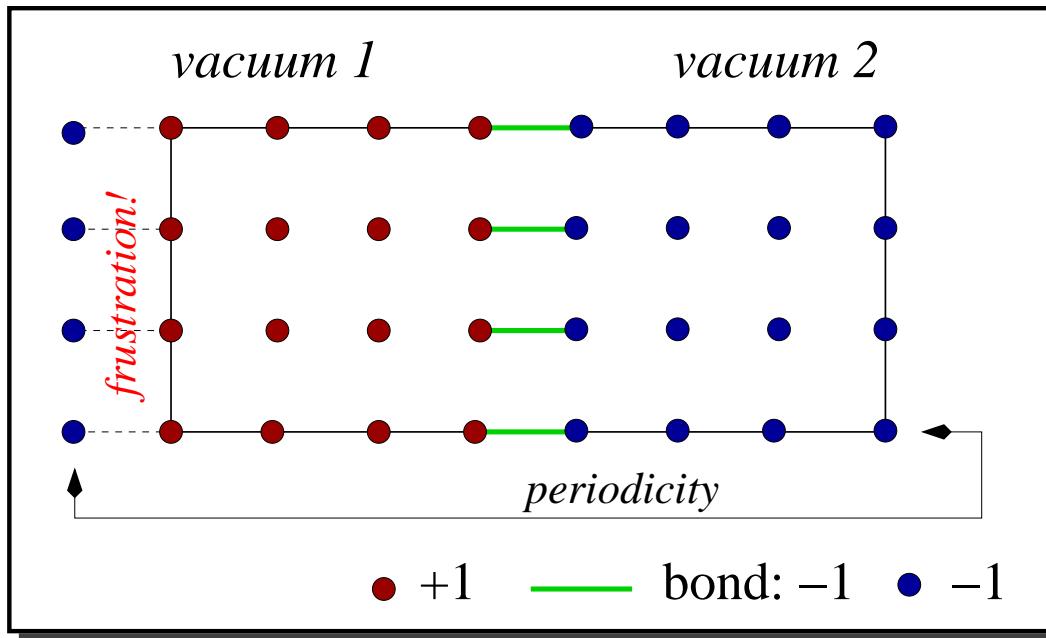
Interface tension: YM + qHiggs

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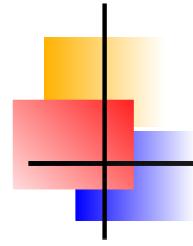


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$$\exp\{-F/T\} = \frac{\text{partition function}}{\text{partition function}}$$

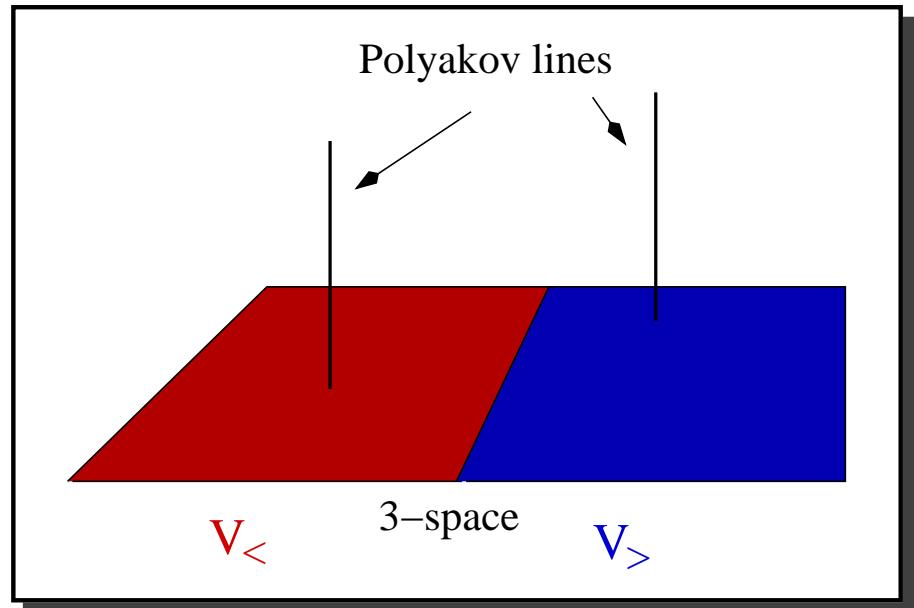


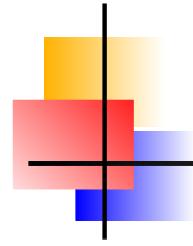
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$$P_< = \sum_{V_<} \text{tr} \prod_t U_0(x)$$

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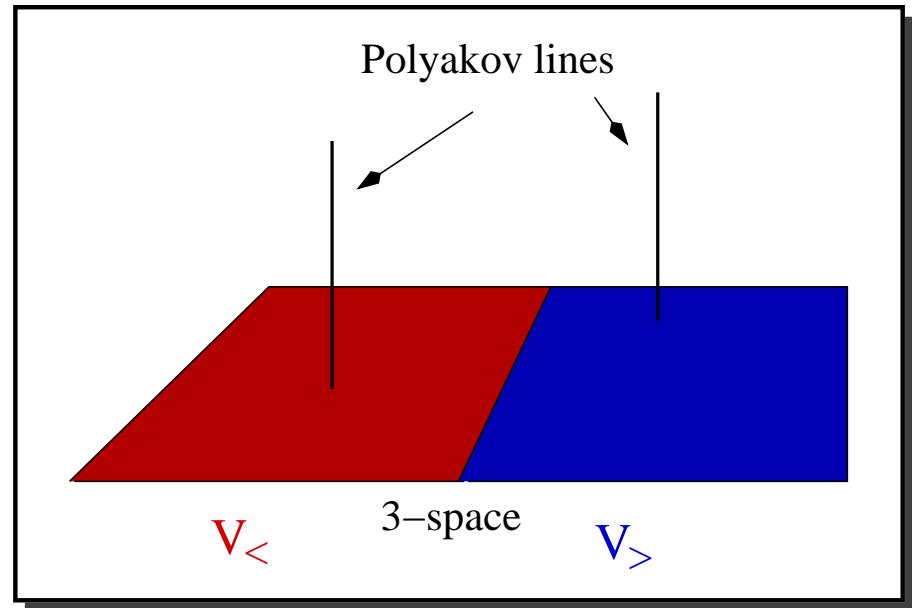
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$$C = n : P \in \mathbb{C} \rightarrow z_n$$



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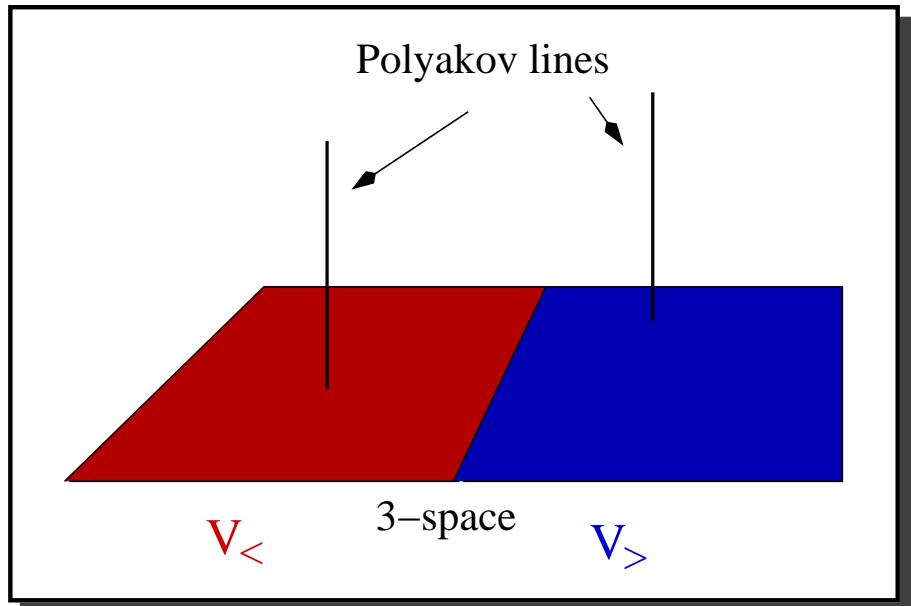
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$$\exp\{-F_{nm}/T\} = \frac{1}{N} \int \mathcal{D}U \mathcal{D}\phi \delta(m, C(P_<)) \delta(n, C(P_>)) e^S$$



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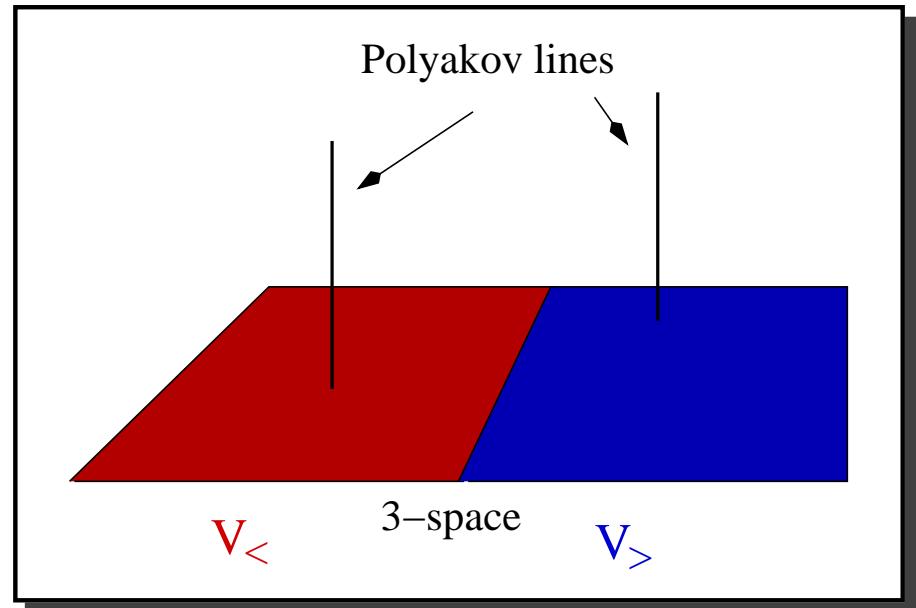
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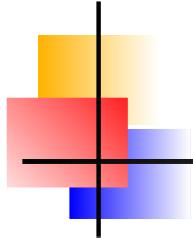
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- **Tunneling coefficient:**

probability $(C(P_<) \text{ and } C(P_>))$ are different





Interface tension: YM + qHiggs

- tunneling coefficient for SU(2):

$C(P_<)$	-1	-1	+1	+1
$C(P_>)$	-1	+1	-1	+1

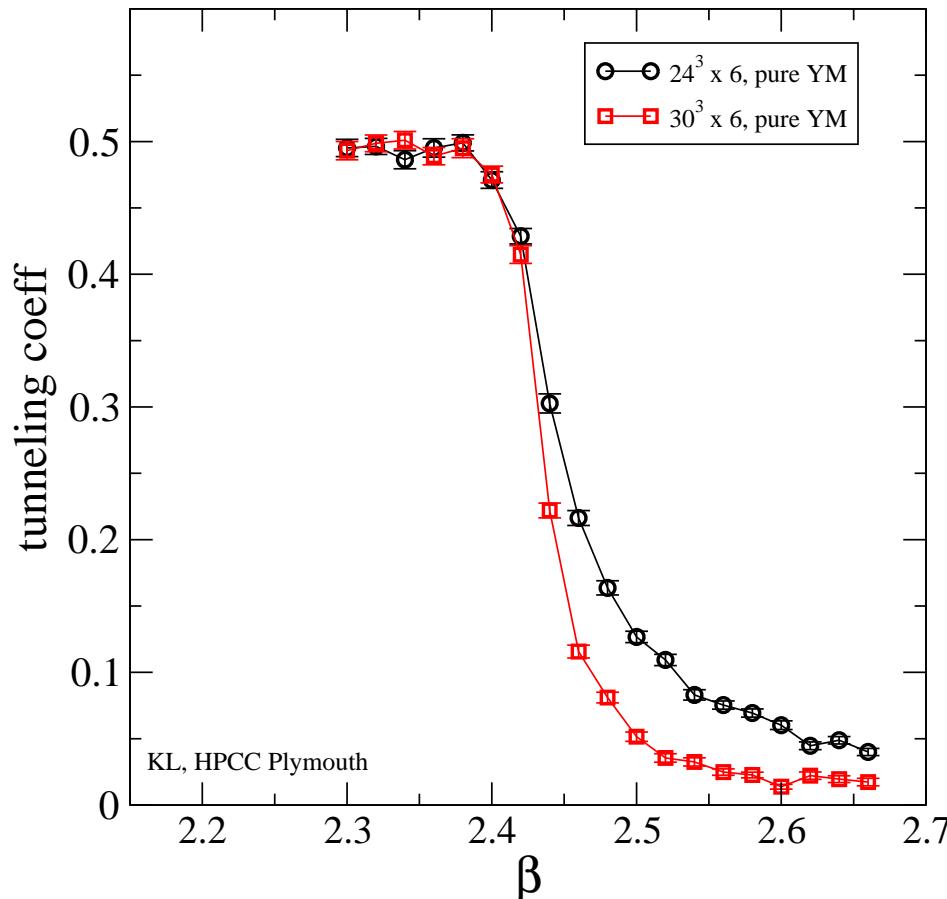
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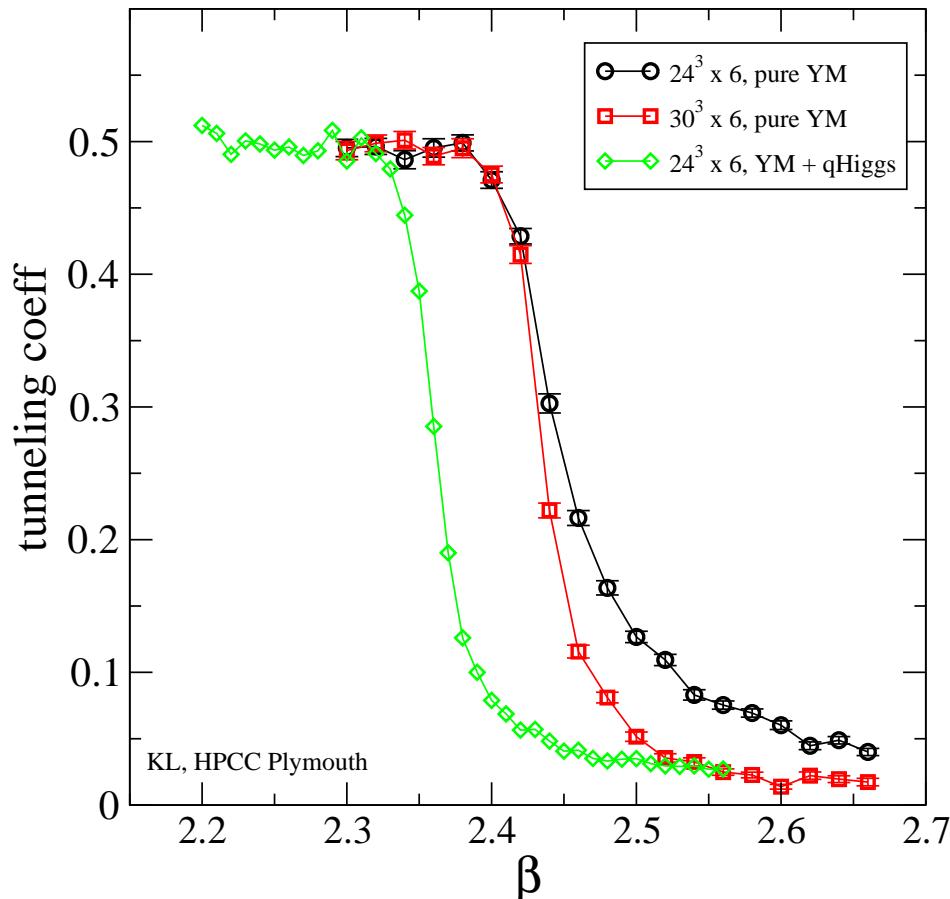


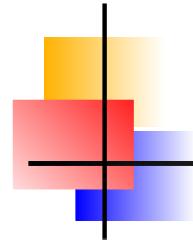
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Interface tension: $SU(3)$ + quarks

- tunneling coefficient for $SU(3)$:

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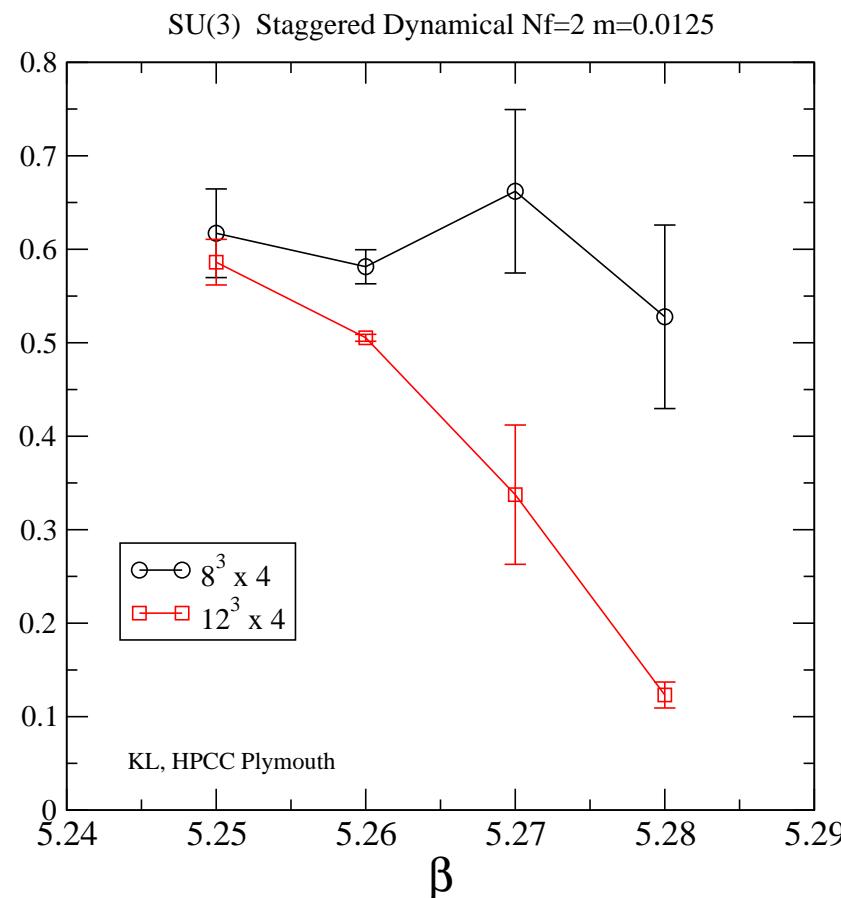
tunneling $\rightarrow 2/3$
SSB $\rightarrow 0$

Interface tension: $SU(3) + \text{quarks}$

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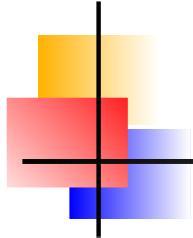
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using MILC configs
 $N_f = 2, m a = 0.0125$

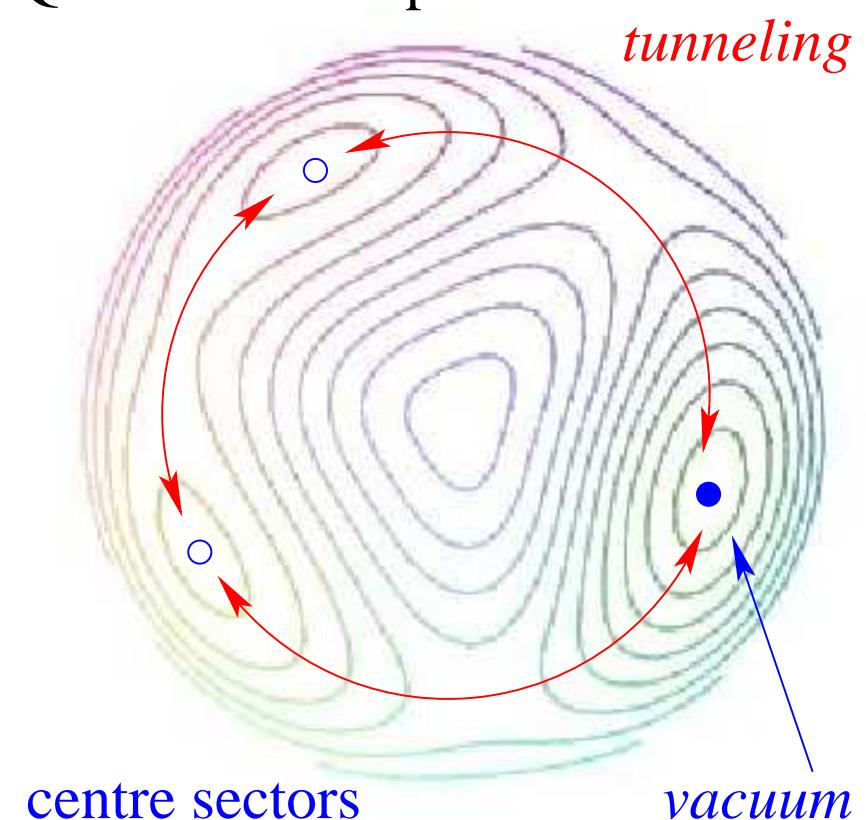
$$\beta_c \approx 5.26$$



Centre sector tunneling:

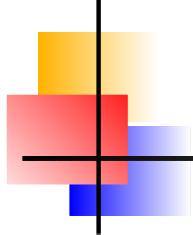
- evidence that centre sector tunneling take places in
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QCD: hadronic phase



centre sectors

vacuum



Centre sector tunneling:

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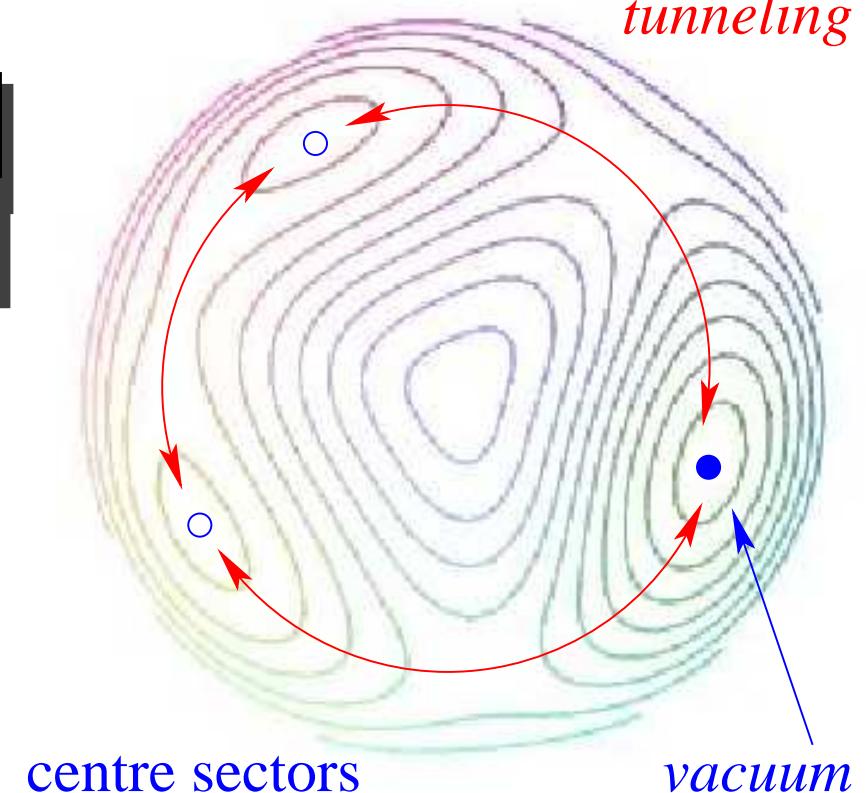
tunneling

- is there FEC for $SU(N)$,
 N odd, such as $SU(3)$?

$$Z = -\frac{1}{2} + i \frac{\sqrt{3}}{2}$$

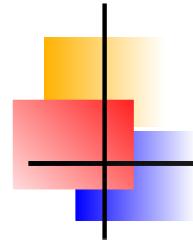
$$Z = 1$$

$$Z = -\frac{1}{2} - i \frac{\sqrt{3}}{2}$$



centre sectors

vacuum

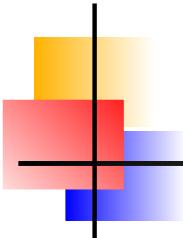


Fermi-Einstein-Condensation (FEC)

Model consideration:

$SU(3)$

- $q(x)$: quarks, m : mass, μ : chemical potential
 A_m : moduli fields \Rightarrow weighted sum over centre sectors



Fermi-Einstein-Condensation (FEC)

Model consideration:

$SU(3)$

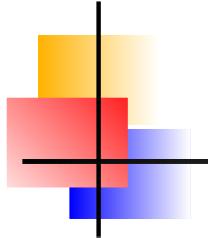
- $q(x)$: quarks, m : mass, μ : chemical potential
- A_m : moduli fields \Rightarrow weighted sum over centre sectors
- partition function: $\exp\{iA_m\} = Z_m \in Z(N_C)$

$$Z = \sum_{m=1}^{N_c} p_m \int \mathcal{D}q \mathcal{D}\bar{q} \exp\{\bar{q}(i\partial + (\textcolor{red}{A}_m + i\mu)\gamma_0 + im)q\}$$

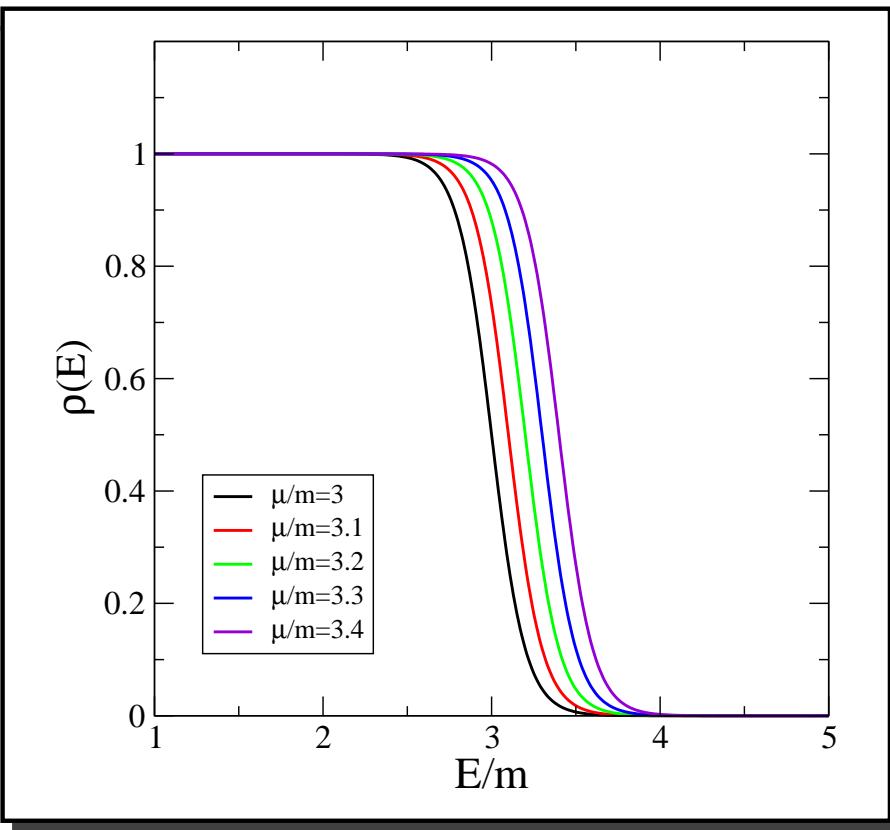
p_m : probability for centre sector m

hadronic phase: $p_m \approx 1/N_c, \forall m$

high T SSB phase: $p_{N_c} = 1, p_m = 0$ for $m = 1 \dots N_c - 1$

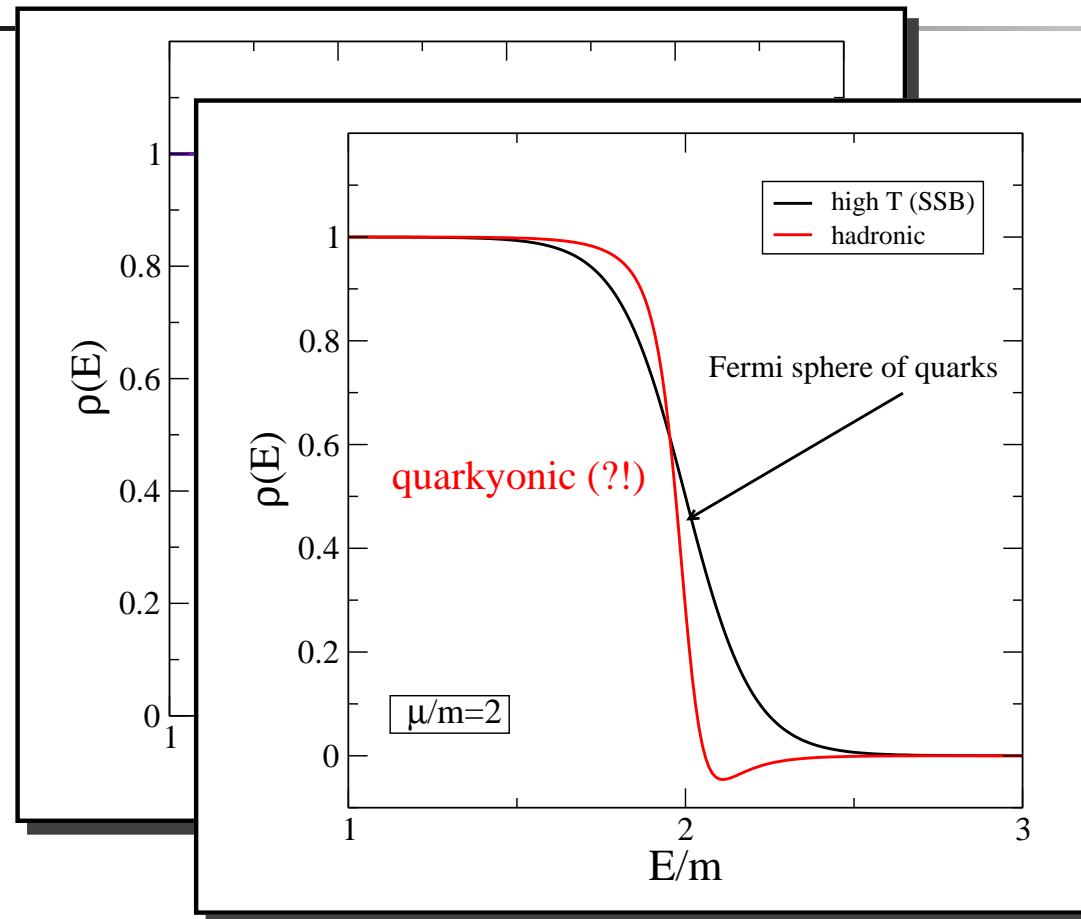


Fermi-Einstein-Condensation (FEC)



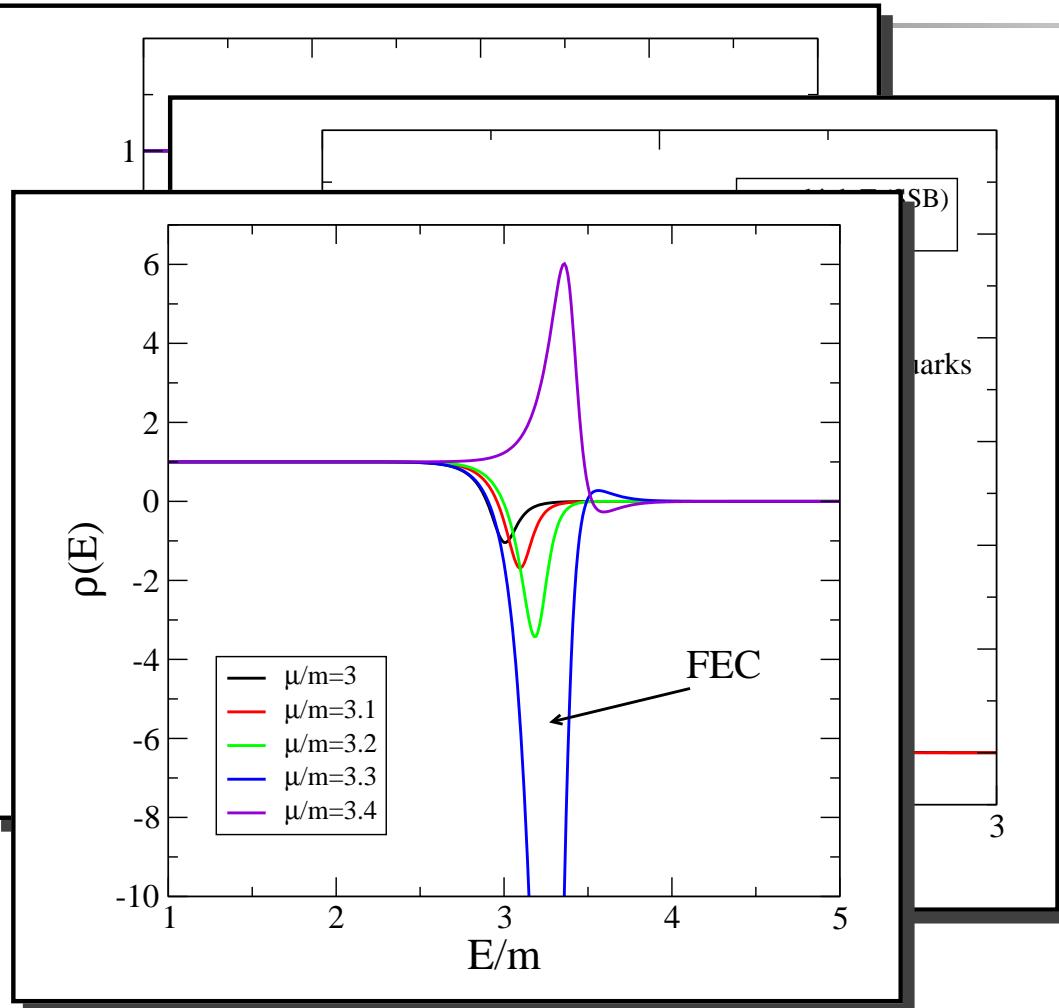
high- T SSB phase

Fermi-Einstein-Condensation (FEC)



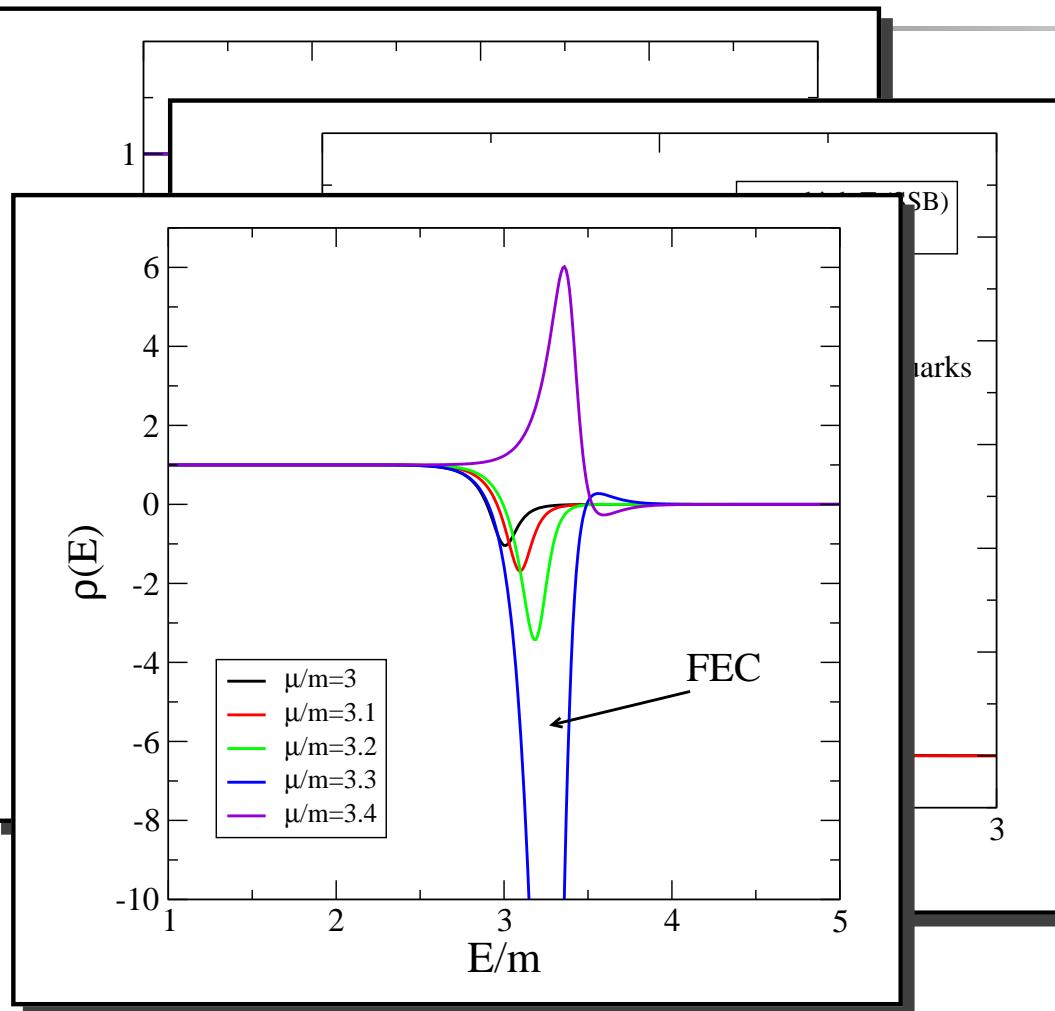
high- T SSB phase
quarkyonic phase

Fermi-Einstein-Condensation (FEC)



high- T SSB phase
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hadronic phase

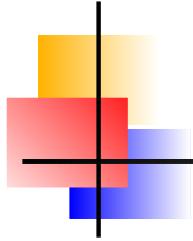
Fermi-Einstein-Condensation (FEC)



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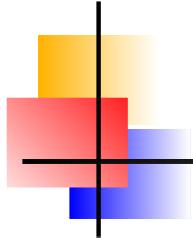
- origin of the Cooper instability:

$$\exp\{-F/T\} = Z = \sum_{m=1}^{N_c} p_m \int \mathcal{D}q \mathcal{D}\bar{q} \dots \rightarrow 0$$



Conclusions:

- Sum over Yang-Mills moduli \Rightarrow confinement
Yes, there is a QCD perturbation theory with confinement!
quantum level: centre-sector tunneling \Rightarrow confinement

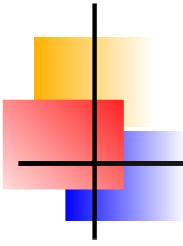


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Fermi-Einstein condensation in $SU(N)$ QCD-like theories

N even [analytical] ✓, N odd [numerical] ✓



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Fermi-Einstein condensation in $SU(N)$ QCD-like theories

- N even [analytical] ✓, N odd [numerical] ✓
- Evidence that centre-sector-tunneling takes place in the hadronic phase:
lattice gauge simulations: $SU(2) + q\text{Higgs}$, $SU(3) + N_f = 2$
 \rightarrow tunneling coefficient
generalised 't Hooft loop ($\langle \psi |^Z \psi \rangle$) with dynamical matter...