

Meson Production on the Nucleon in the Giessen K-Matrix Approach

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Agenda:

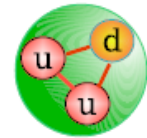
- Probing Baryons by Resonance excitation
- Unitary K-Matrix Approach to Meson Production
- Photoproduction on the Nucleon: ωN and $K\Lambda$
- Outlook: Hadronic and Leptonic Production on Nuclei
- Summary

Probing Baryons by Hadronic and Electromagnetic Probes

pion



nucleon



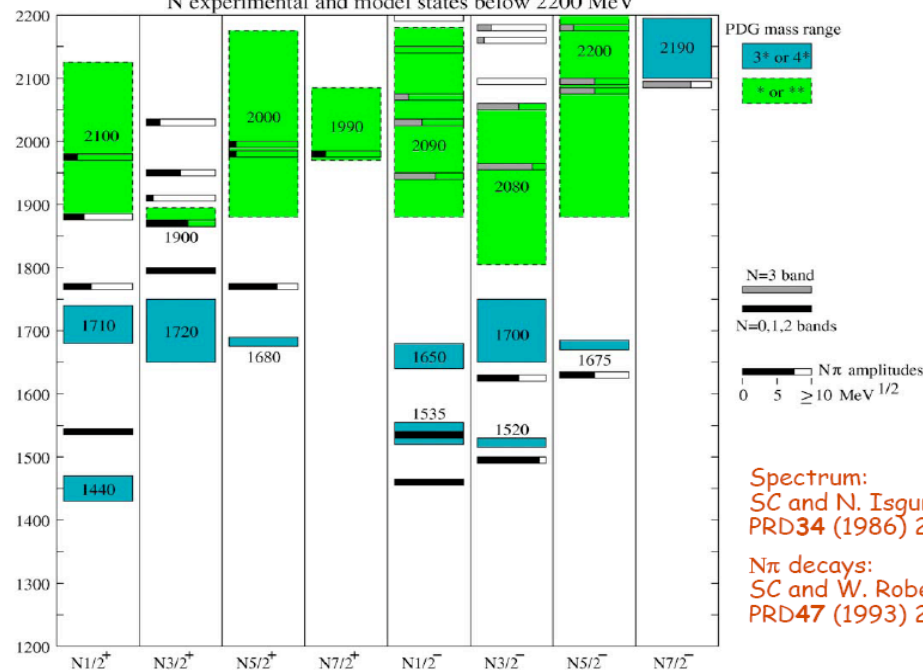
First resonance: $\Delta^{\frac{3}{2}}(1232)$

=> first indication to nucleon structure

both in πN and γN scattering!

Nucleon model states and $N\pi$ couplings

N experimental and model states below 2200 MeV



Problem:
Missing Resonances!



Confronting Theory with Reality:

Hunting missing resonances: photoproduction.

Recent years:
new photoproduction data, new opportunity

Motivation

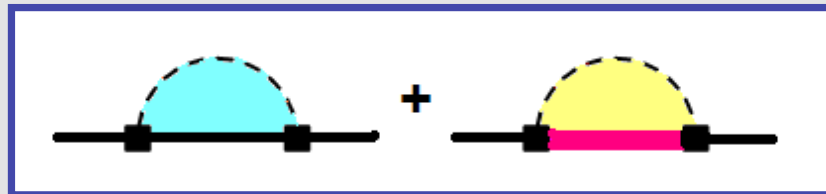
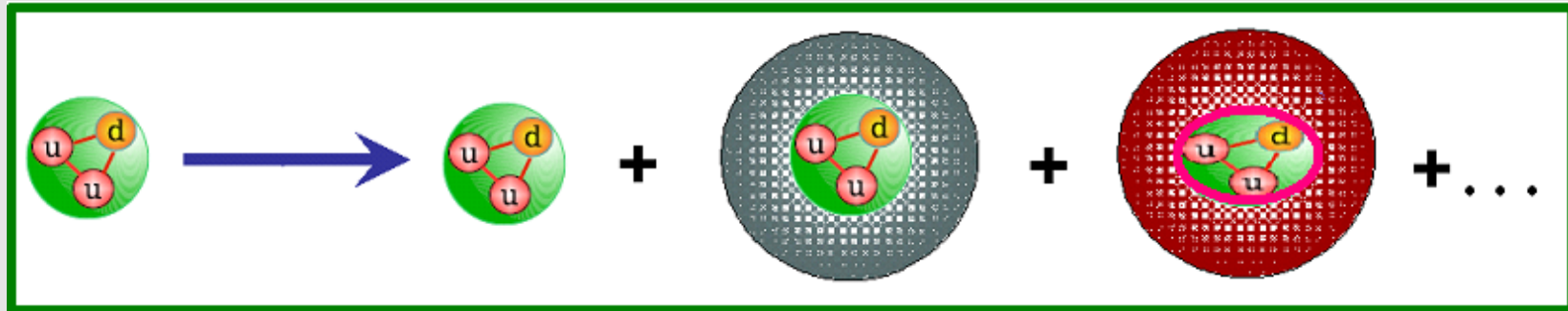
- $\gamma N \leftrightarrow \omega N$ (SAPHIR)

Search for nucleon resonances which couples strongly to ωN channel

- $\gamma N \leftrightarrow K\Lambda$ (SAPHIR, CLAS, SPring-8)

New physical domain:
search for 'missing' states which couple strongly to the associated strangeness channel.

Baryons: Quark Model vs. Dynamical Models



Our Choice of Degrees of Freedom (besides photon):

- asymptotic hadronic states: mesons and baryons
- empirical hadron masses
- Lagrangian model with phenomenological coupling constants
- regularization by form factors
- unitary coupled channels approach

Structure of the Lagrangian:

Born and t -channel terms:

$$\begin{aligned} \mathcal{L}_{Born} + \mathcal{L}_t = & -\bar{u}_{B'}(p') \left[\frac{g_{\tilde{\varphi}}}{m_B + m_{B'}} \gamma_5 \gamma_\mu (\partial^\mu \tilde{\varphi}) + g_\eta i \gamma_5 \eta + g_S \Phi_S \right. \\ & \left. + g_V \left(\gamma_\mu V^\mu + \frac{\kappa_V}{2m_N} \sigma_{\mu\nu} V^{\mu\nu} \right) \right] u_B(p) \\ & - \frac{g_S}{2m_\pi} (\partial_\mu \varphi') (\partial^\mu \varphi) \Phi_S - g_V \varphi' (\partial_\mu \varphi) V^\mu - \frac{g_V}{4m_\varphi} \epsilon_{\mu\nu\rho\sigma} V^{\mu\nu} V'^{\rho\sigma} \varphi. \end{aligned}$$

positive-parity **spin- $\frac{1}{2}$** resonances, PV coupling is used:

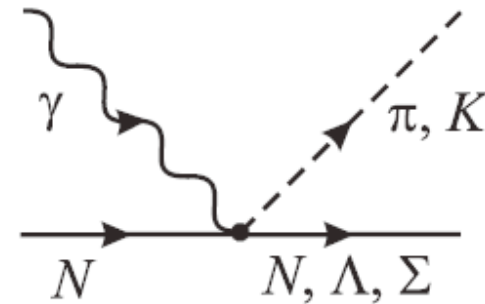
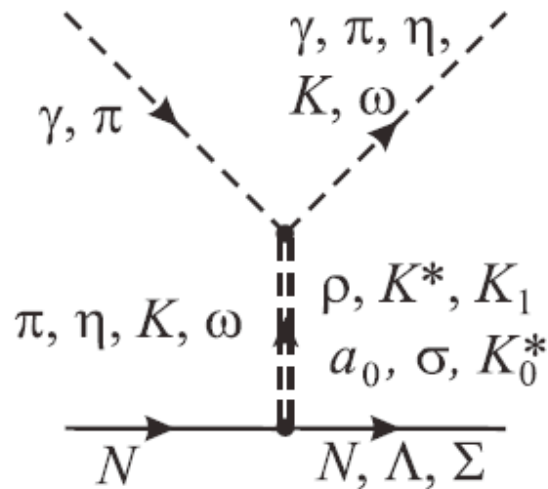
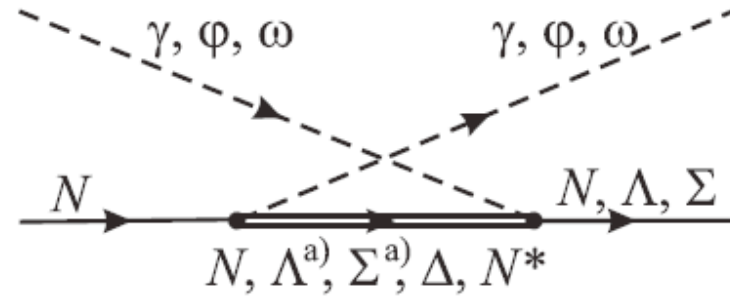
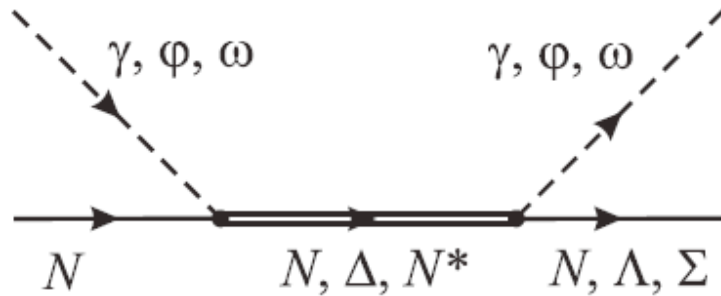
$$\mathcal{L}_{\frac{1}{2}B\varphi}^{PV} = -\frac{g_{RB\varphi}}{m_R \pm m_B} \bar{u}_R \begin{pmatrix} \gamma_5 \\ i \end{pmatrix} \gamma_\mu u_B \partial^\mu \varphi.$$

π, η, K

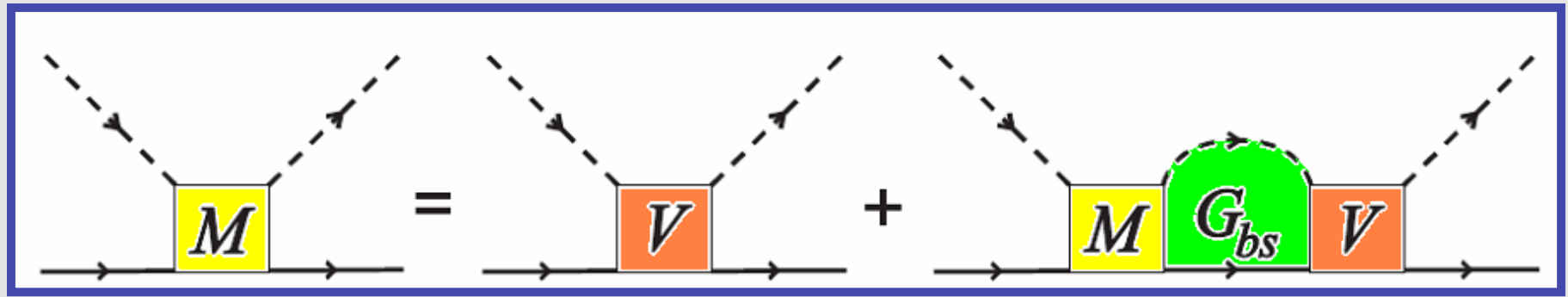
negative-parity **spin- $\frac{1}{2}$** resonances, PS coupling is used:

$$\mathcal{L}_{\frac{1}{2}B\varphi}^{PS} = -g_{RB\varphi} \bar{u}_R \begin{pmatrix} 1 \\ -i\gamma_5 \end{pmatrix} u_B \varphi.$$

Tree-Level Diagrams in the BS Equation



The Giessen Model: Coupled Channels K-Matrix



$$G_{bs} \equiv \frac{P}{H - \omega} + i\pi\delta(H - \omega):$$

$$\begin{aligned} \mathcal{K} &= \mathcal{V} + \mathcal{V} \frac{P}{H - \omega} \mathcal{K} \sim \mathcal{V} \\ \mathcal{M} &= \mathcal{K} + i\mathcal{K} \delta(H - \omega) \mathcal{M} \end{aligned}$$

$$\mathcal{M}_{\alpha\beta} \equiv \left[\frac{\mathcal{K}}{1 - i\mathcal{K}} \right]_{\alpha\beta} \sim \left[\frac{\mathcal{V}}{1 - i\mathcal{V}} \right]_{\alpha\beta}$$

- Widths Γ_{N^*} are generated dynamically
- Unitarity for $\mathcal{K} = \mathcal{K}^+$, Resonances and background!

The Giessen Model: Multi-Channel T-Matrix

Coupled Channel Picture of Resonance Excitation

- Each resonance can be reached through each asymptotic channel
- T matrix provides unitary, analytic structure
- all channels (e.g. πN , ρN) couple to all other channels in intermediate state
- photon multipoles ($E_{l\pm}$, $M_{l\pm}$) directly related to \mathbf{T}

$$T(J^\pi) = \begin{pmatrix} T_{\pi N \rightarrow \pi N} & T_{\eta N \rightarrow \pi N} & T_{\gamma N \rightarrow \pi N} & T_{\rho N \rightarrow \pi N} & T_{\sigma N \rightarrow \pi N} & T_{K\Lambda \rightarrow \pi N} & T_{K\Sigma \rightarrow \pi N} \\ T_{\pi N \rightarrow \eta N} & T_{\eta N \rightarrow \eta N} & T_{\gamma N \rightarrow \eta N} & T_{\rho N \rightarrow \eta N} & T_{\sigma N \rightarrow \eta N} & T_{K\Lambda \rightarrow \eta N} & T_{K\Sigma \rightarrow \eta N} \\ T_{\pi N \rightarrow \gamma N} & T_{\eta N \rightarrow \gamma N} & T_{\gamma N \rightarrow \gamma N} & T_{\rho N \rightarrow \gamma N} & T_{\sigma N \rightarrow \gamma N} & T_{K\Lambda \rightarrow \gamma N} & T_{K\Sigma \rightarrow \gamma N} \\ T_{\pi N \rightarrow \rho N} & T_{\eta N \rightarrow \rho N} & T_{\gamma N \rightarrow \rho N} & T_{\rho N \rightarrow \rho N} & T_{\sigma N \rightarrow \rho N} & T_{K\Lambda \rightarrow \rho N} & T_{K\Sigma \rightarrow \rho N} \\ T_{\pi N \rightarrow \sigma N} & T_{\eta N \rightarrow \sigma N} & T_{\gamma N \rightarrow \sigma N} & T_{\rho N \rightarrow \sigma N} & T_{\sigma N \rightarrow \sigma N} & T_{K\Lambda \rightarrow \sigma N} & T_{K\Sigma \rightarrow \sigma N} \\ T_{\pi N \rightarrow K\Lambda} & T_{\eta N \rightarrow K\Lambda} & T_{\gamma N \rightarrow K\Lambda} & T_{\rho N \rightarrow K\Lambda} & T_{\sigma N \rightarrow K\Lambda} & T_{K\Lambda \rightarrow K\Lambda} & T_{K\Sigma \rightarrow K\Lambda} \\ T_{\pi N \rightarrow K\Sigma} & T_{\eta N \rightarrow K\Sigma} & T_{\gamma N \rightarrow K\Sigma} & T_{\rho N \rightarrow K\Sigma} & T_{\sigma N \rightarrow K\Lambda} & T_{K\Lambda \rightarrow K\Sigma} & T_{K\Sigma \rightarrow K\Sigma} \end{pmatrix}$$

ωN photoproduction

ωN final state: Hard to analyze

- Vector mesons:

γN : different helicity states $\frac{1}{2} - 1, \frac{1}{2} + 1$

ωN : different helicity states $\frac{1}{2} - 1, \frac{1}{2} + 1, \frac{1}{2} + 0$

=> different resonance coupling to helicity states.

- photoproduction

analogy with $\gamma N \rightarrow \pi N$: needs $\pi N \rightarrow \pi N$ as an input.

Same for $\gamma N \rightarrow \omega N$ photoproduction:

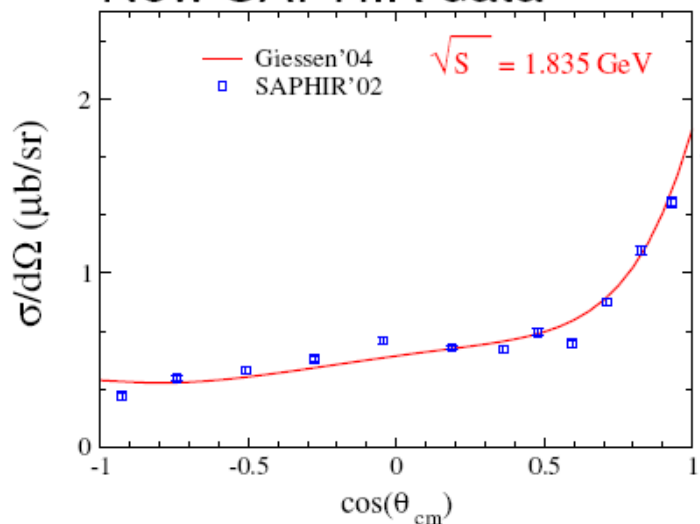
knowledge about $\pi N \rightarrow \omega N$ channel is mandatory.

New channel: include also $j=5/2$ states

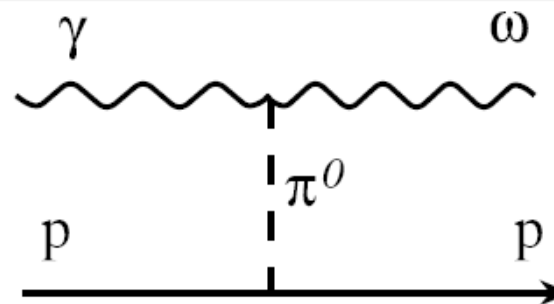
Production Dynamics of ω Photoproduction

ω photoproduction

New SAPHIR data

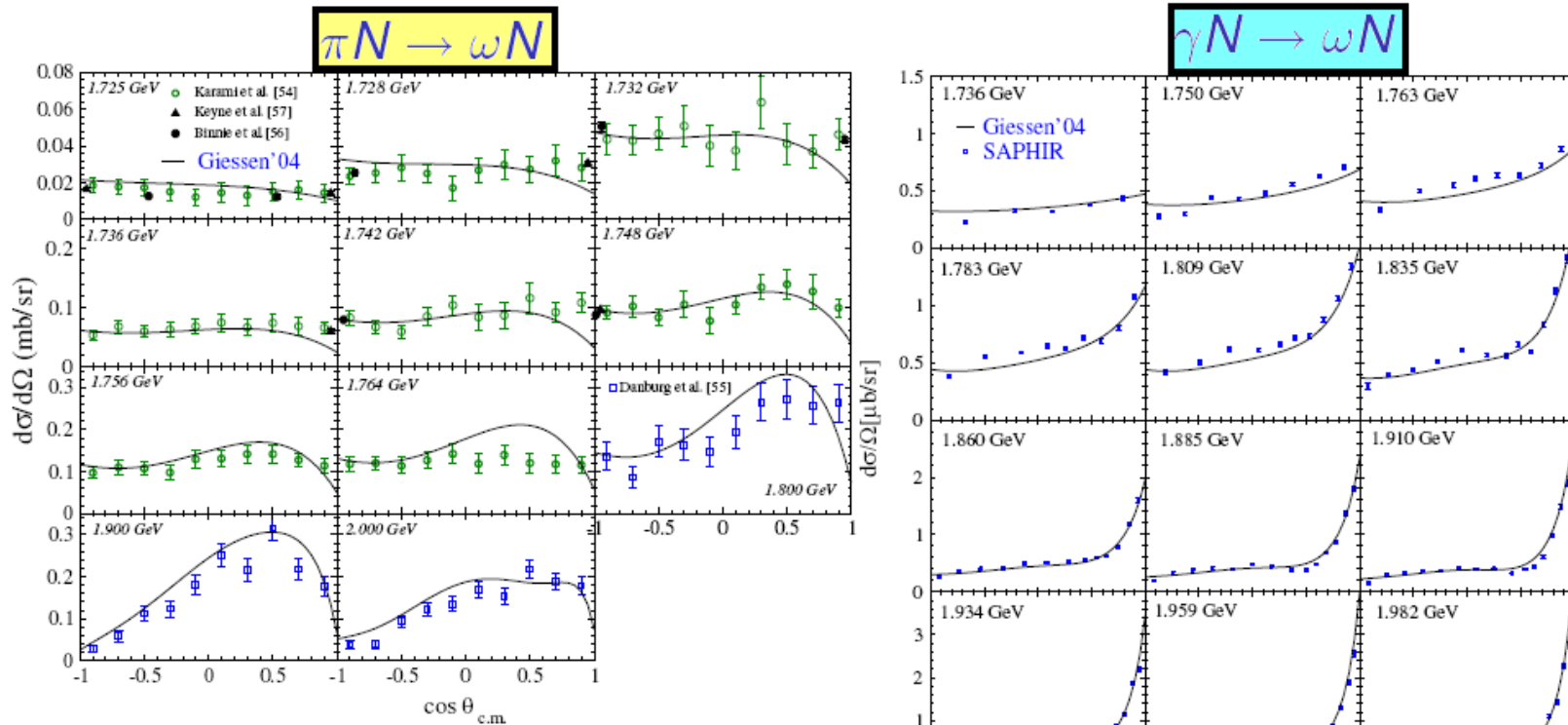


Friman and Soyeur NPA600:477
Forward directions: π^0 exchange !



$\Rightarrow \pi^0$ -exchange generates strong background contributions

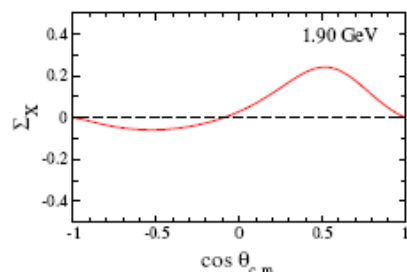
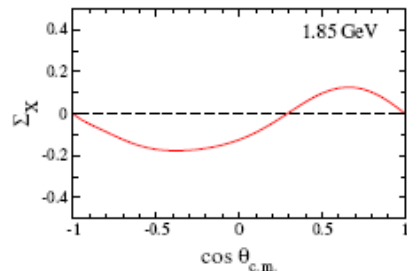
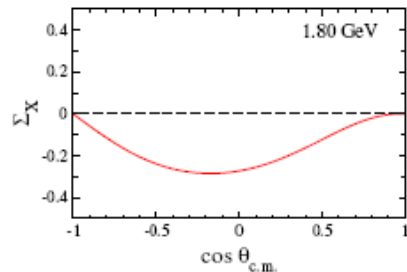
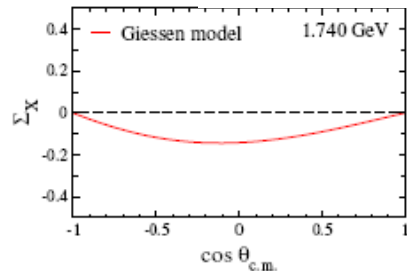
Giessen model. Results for the $(\pi, \gamma)N \rightarrow \omega N$ reactions



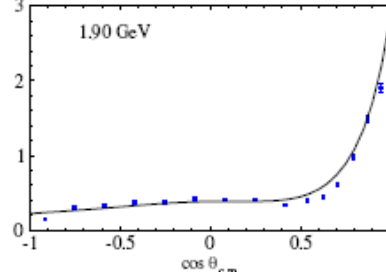
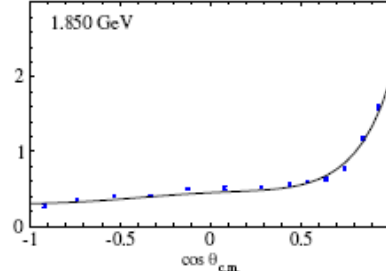
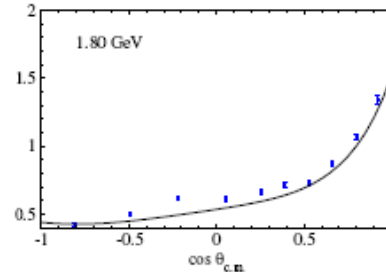
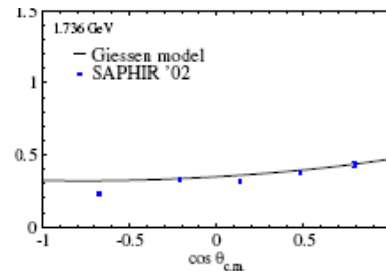
Combined analysis \implies more constraints on resonance properties. V. Shklyar et al. PRC71, 055206 (2005)

$\pi N \rightarrow \omega N$: More precise data needed

Giessen model. Results for $(\pi, \gamma)N \rightarrow \omega N$



Σ_X



$d\sigma/d\Omega$

A. Titov, T.-S.H. Lee
PRC66:015204

and

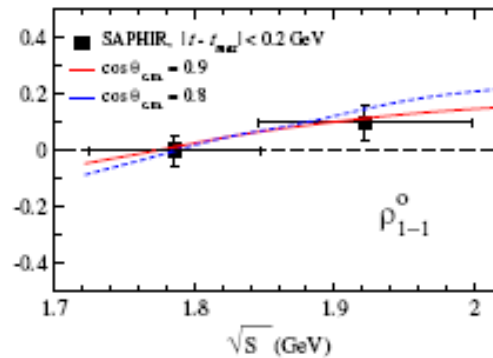
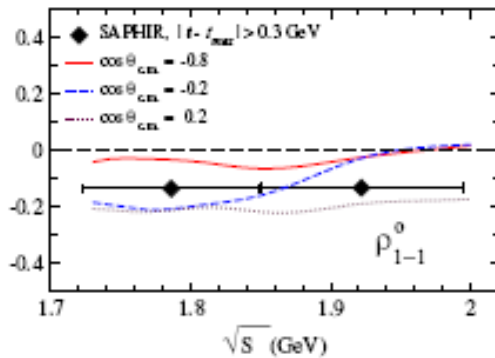
Q. Zhao PRC63:025203

Asymmetry is sensitive to the
production mechanism !

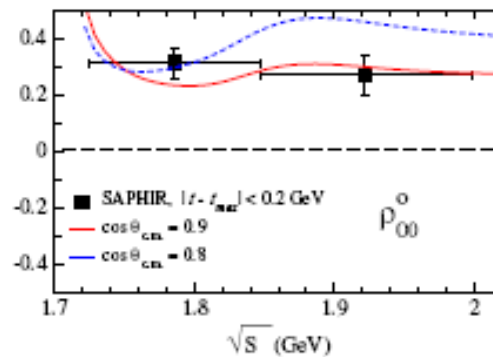
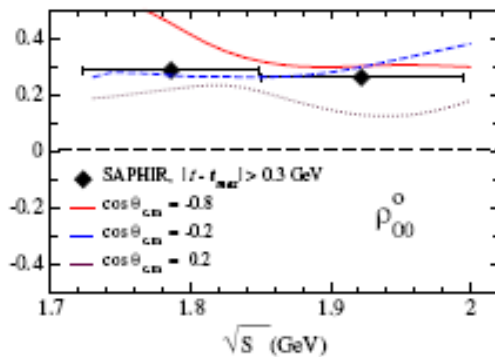
Beam asymmetry:

Changes sign with
increasing energy:
t-channel dominates the
forward directions.

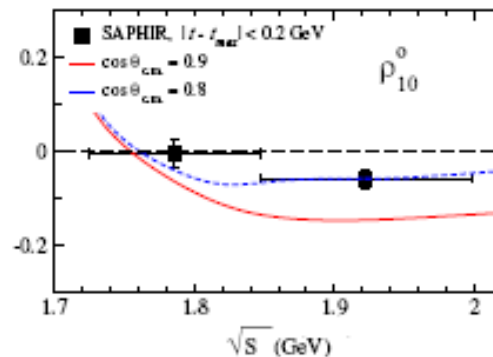
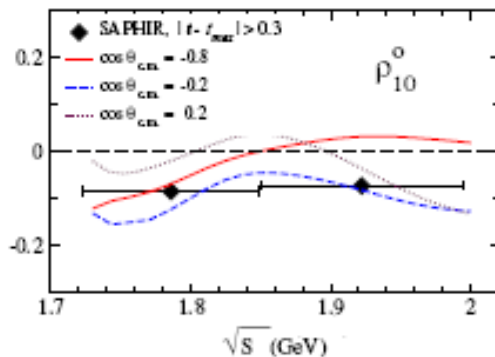
$\gamma N \rightarrow \omega N$ Spin Density Matrix



V. Shklyar, H. Lenske, U. Mosel, G. Penner,
PRC71, 055206 (2005)



$$\rho_{\lambda\omega\lambda'\omega'} = \frac{\sum_{S_i S_f \lambda_\gamma} T_{S_i \lambda_\gamma S_f \lambda_\omega} T_{S_i \lambda_\gamma S_f \lambda'_\omega}}{\sum_{S_i S_f \lambda_\gamma \lambda_\omega} T_{S_i \lambda_\gamma S_f \lambda_\omega} T_{S_i \lambda_\gamma S_f \lambda_\omega}}$$

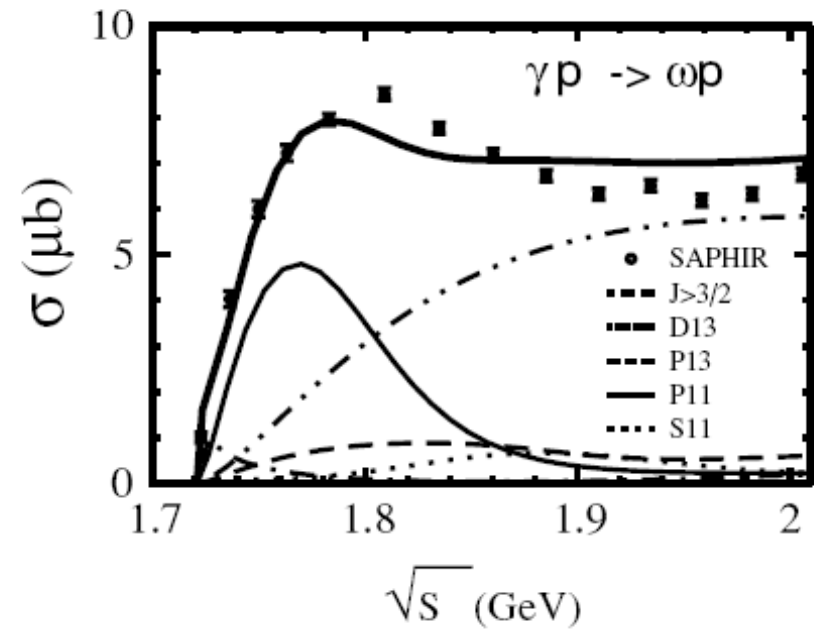
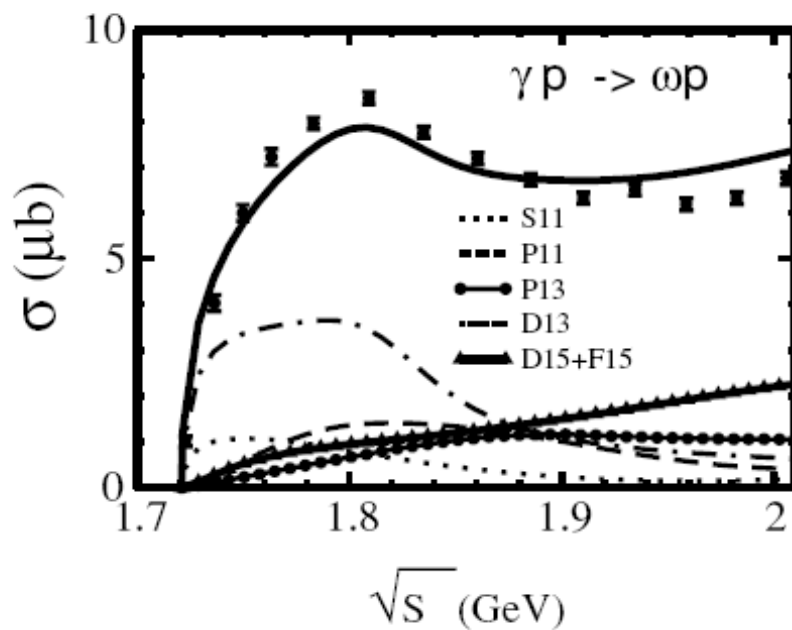


Spin density matrix provides an additional constraint on the reaction amplitude.

Sensitivity on Reaction Input: Spin Density Data and Cross Sections

with ω spin density matrix data

without ω spin density matrix data



$$\chi^2 \approx 4.2$$

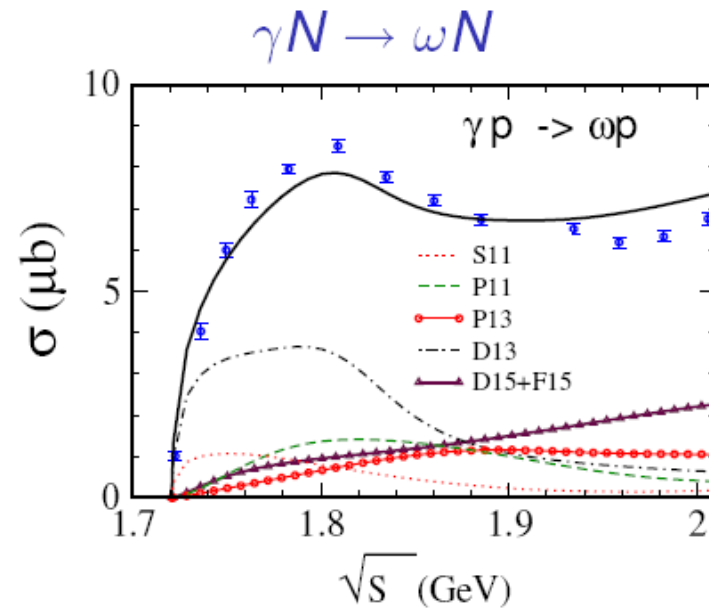
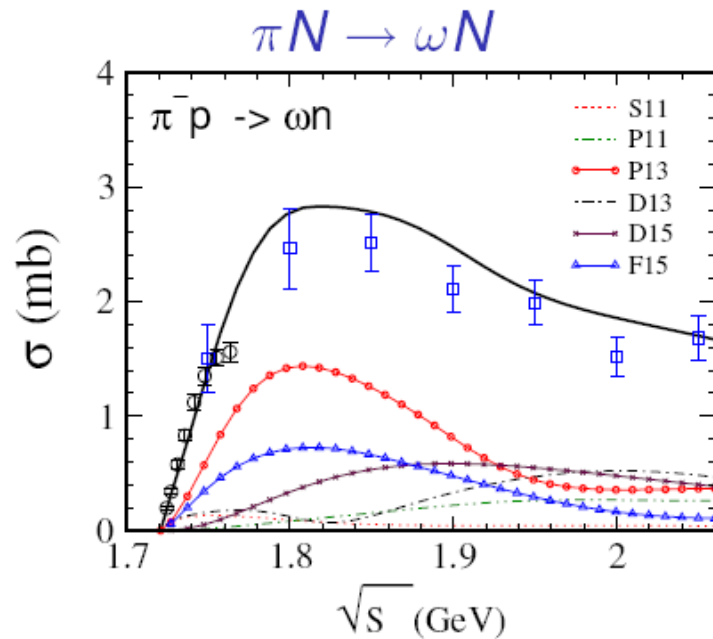
(SAPHIR Data)

$$\chi^2 \approx 6.5$$

V. Shklyar et al., PRC71 055206 (2005)

G. Penner et al., PRC66 055211 (2002)

Giessen model. Results for $(\pi, \gamma)N \rightarrow \omega N$

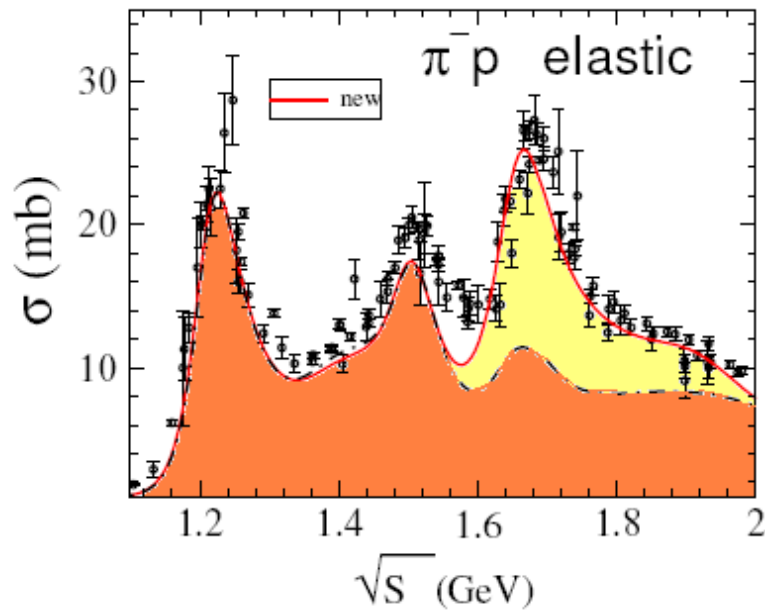


- P_{13} : interference between resonance and background
- strong $N^*(\frac{5}{2})$ coupling to ωN
- D_{13} shows minor influence

- strong Born and π^0 -exchange contributions
- D_{13} is due to π^0 -exchange

Phys.Rev.C71:055206,2005

Contributions of $j > 3/2$ Resonances to πN ?



New results:

with spin- $\frac{5}{2}$ resonances !

V. Shklyar et al .PRC71,
055206 (2005)

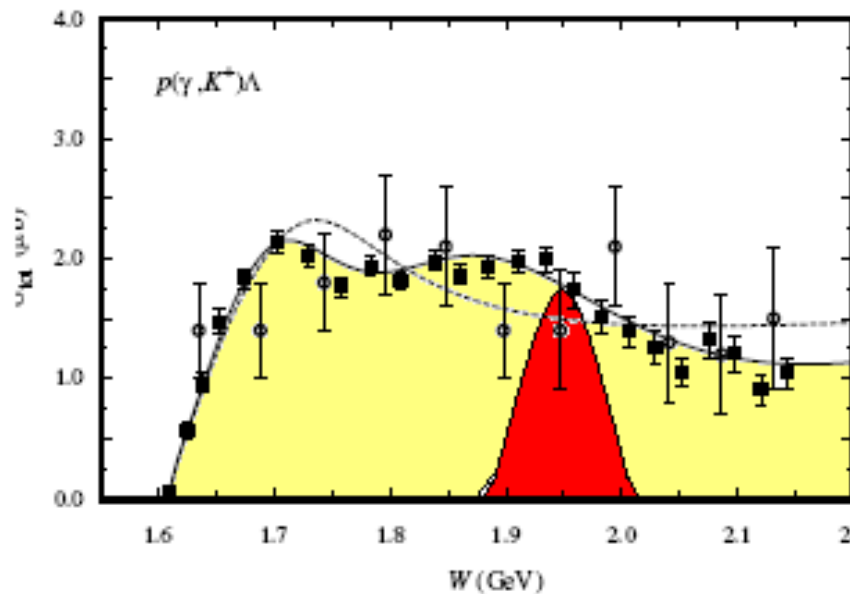
Optical theorem:

$$\text{Im}T_{\pi N \rightarrow \pi N} \sim \sigma_{\pi N \rightarrow \omega N} + \dots$$

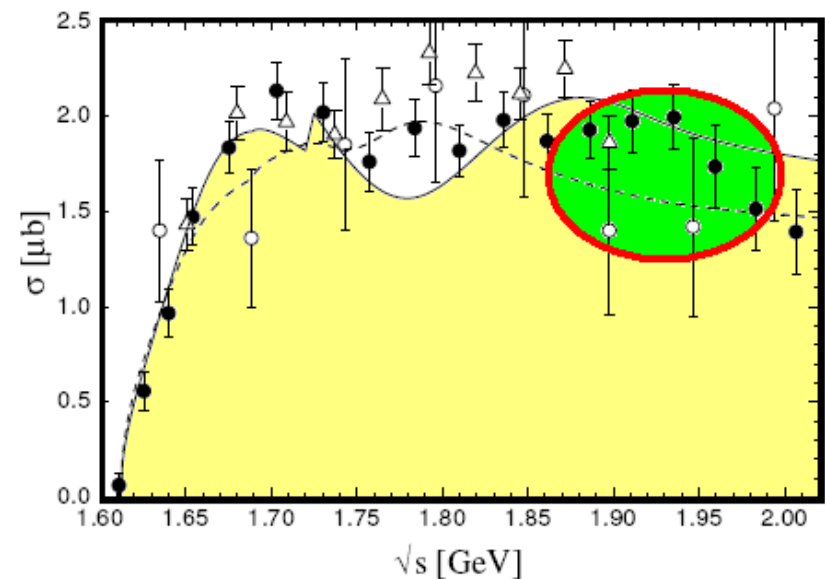
πN and ωN :
 $D_{15}(1675)$, $F_{15}(1680)$, $F_{15}(2000)$

New Resonances in $\gamma + N \rightarrow K + \Lambda$?

SAPHIR, Tran et al. (1998): structure at about 1.9 GeV in $\gamma + N \rightarrow K + \Lambda$

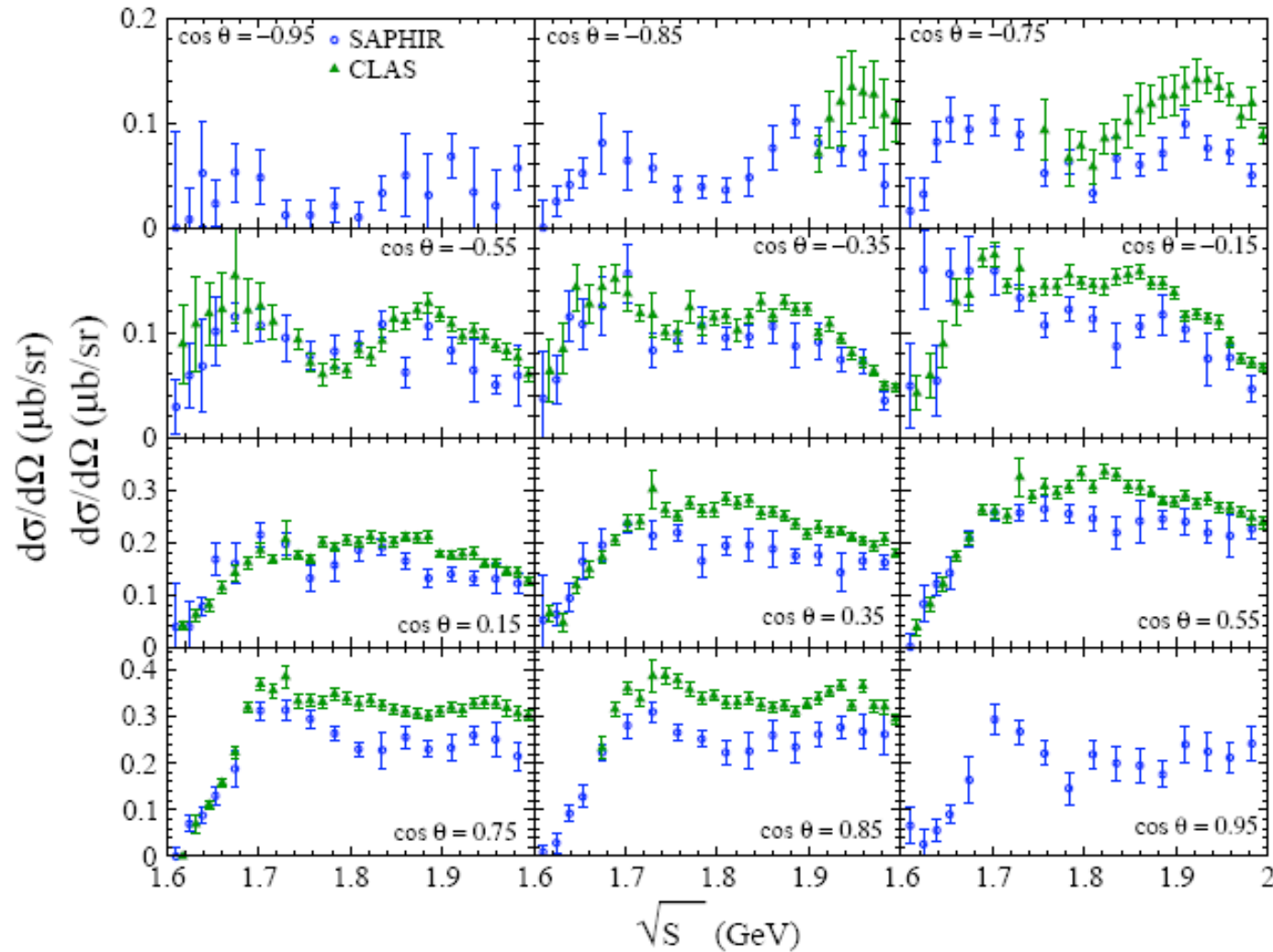


Mart, Bennhold (PRC 61: 012201): claim for a new $D_{13}(1960)$ resonance



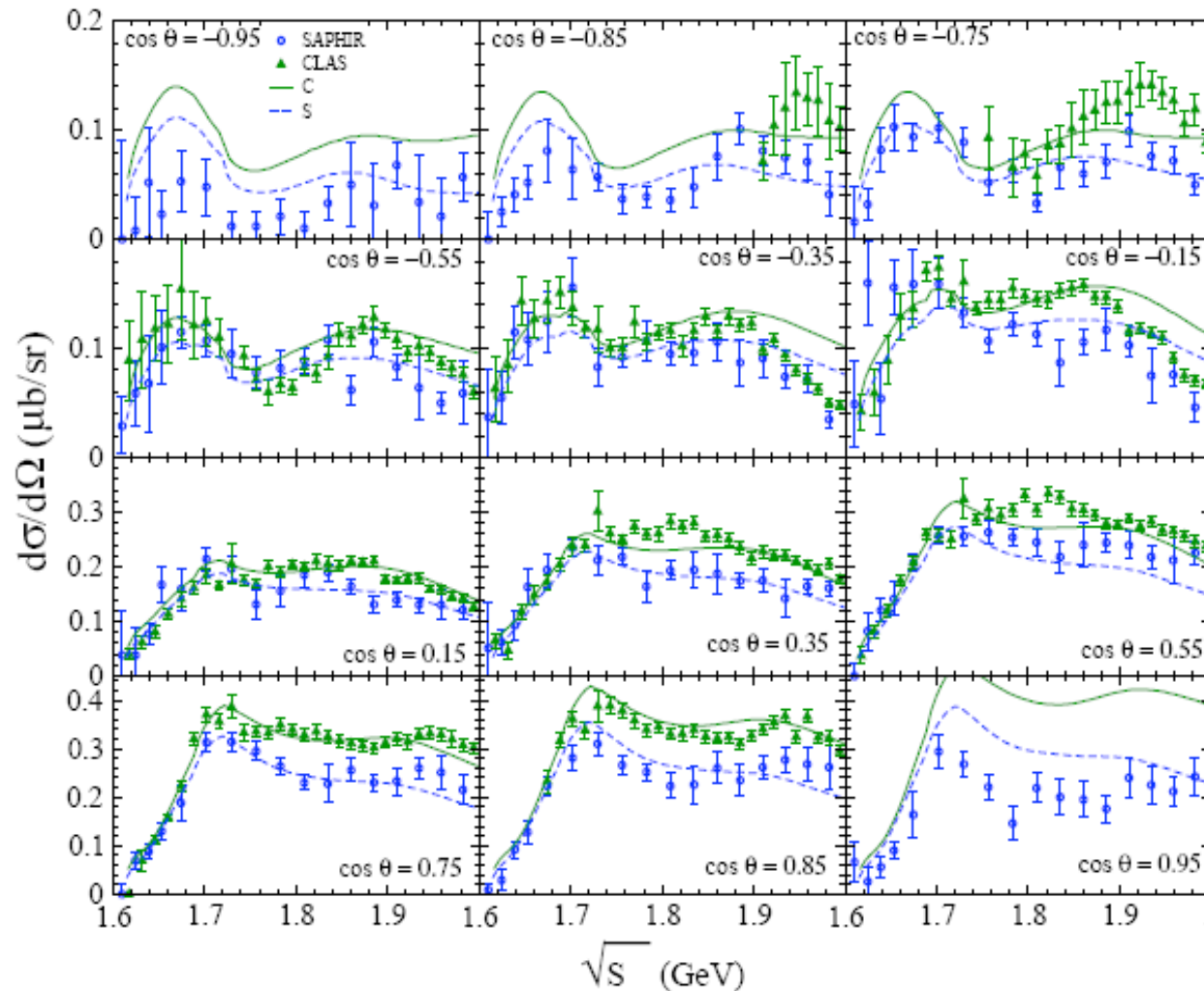
Penner, Mosel (PRC66: 055212): Born term and t -channel interference

New $\gamma N \rightarrow K\Lambda$ Data from SAPHIR and CLAS

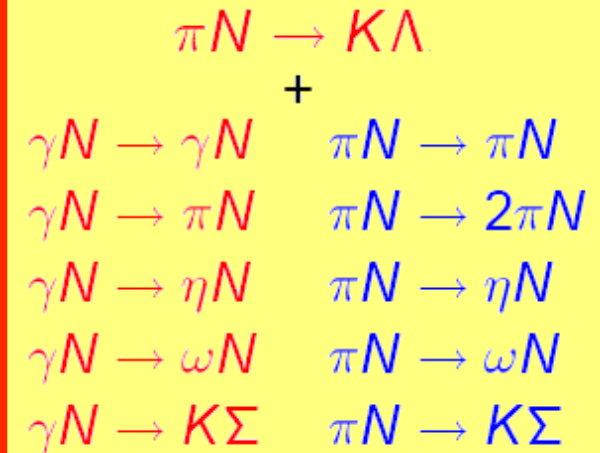


agree in the interesting energy near 1.9 GeV.
 data: position of second peak is not fixed.

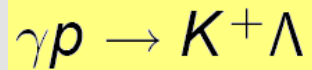
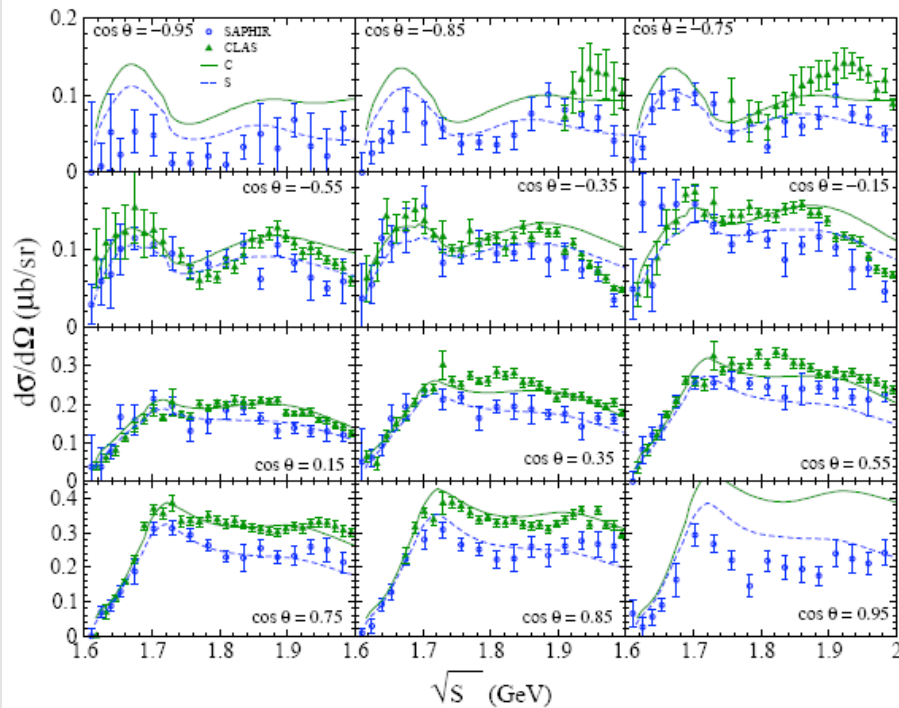
$\gamma N \rightarrow K\Lambda$ Results from the Giessen Model



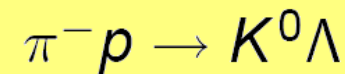
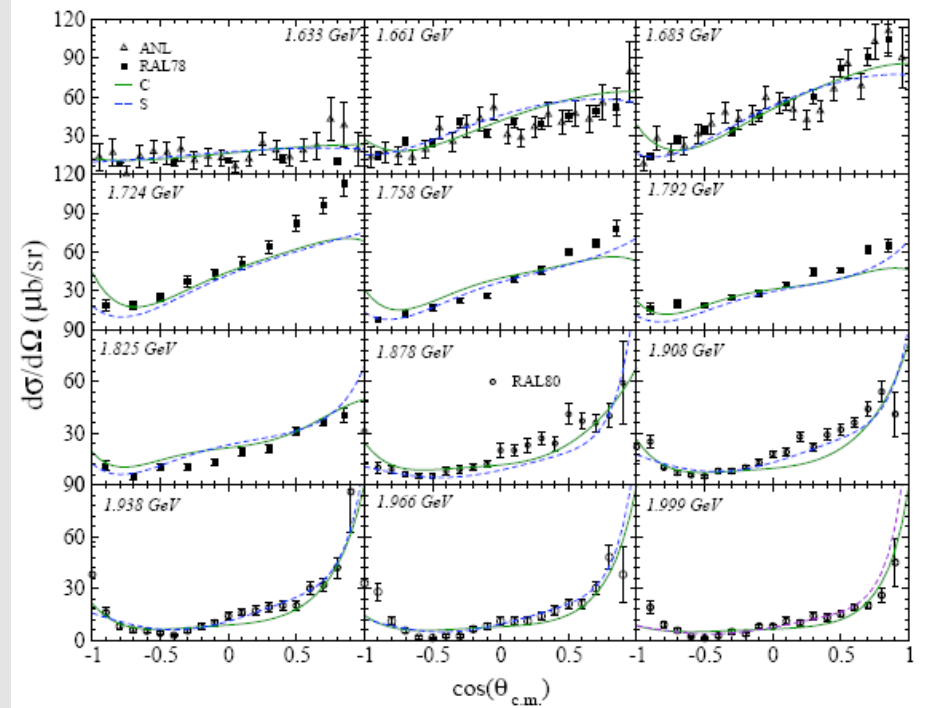
- separate fits to SAPHIR and CLAS data
- constraints from other hadronic and γ channels



$\gamma N \rightarrow K\Lambda$: Influence on other Channels



separate fits to **SAPHIR** and **CLAS** data



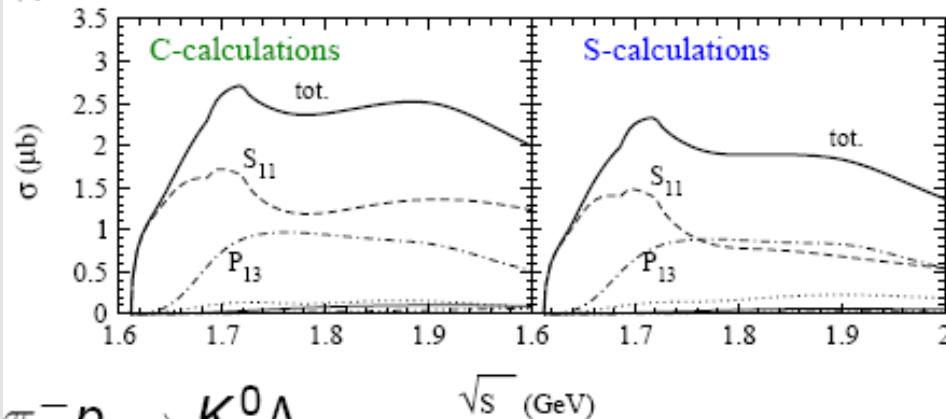
Disagreement between the CLAS and SAPHIR data does not affect $\pi^- p \rightarrow K^0 \Lambda$

Consequences on Reaction Dynamics of $K\Lambda$ Channels from CLAS/SAPHIR Solutions

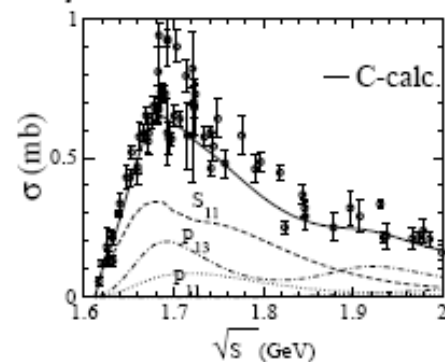
Resonance contributions:

$S_{11}(1650)$ $P_{13}(1720)$ and $P_{13}(1900)$

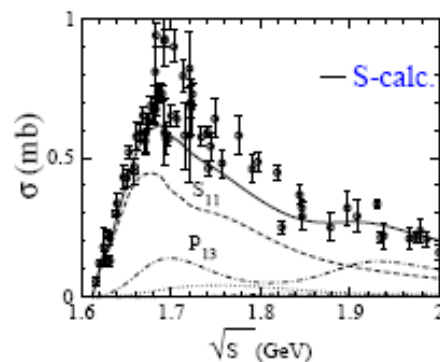
$\gamma p \rightarrow K^+ \Lambda$



$\pi^- p \rightarrow K^0 \Lambda$

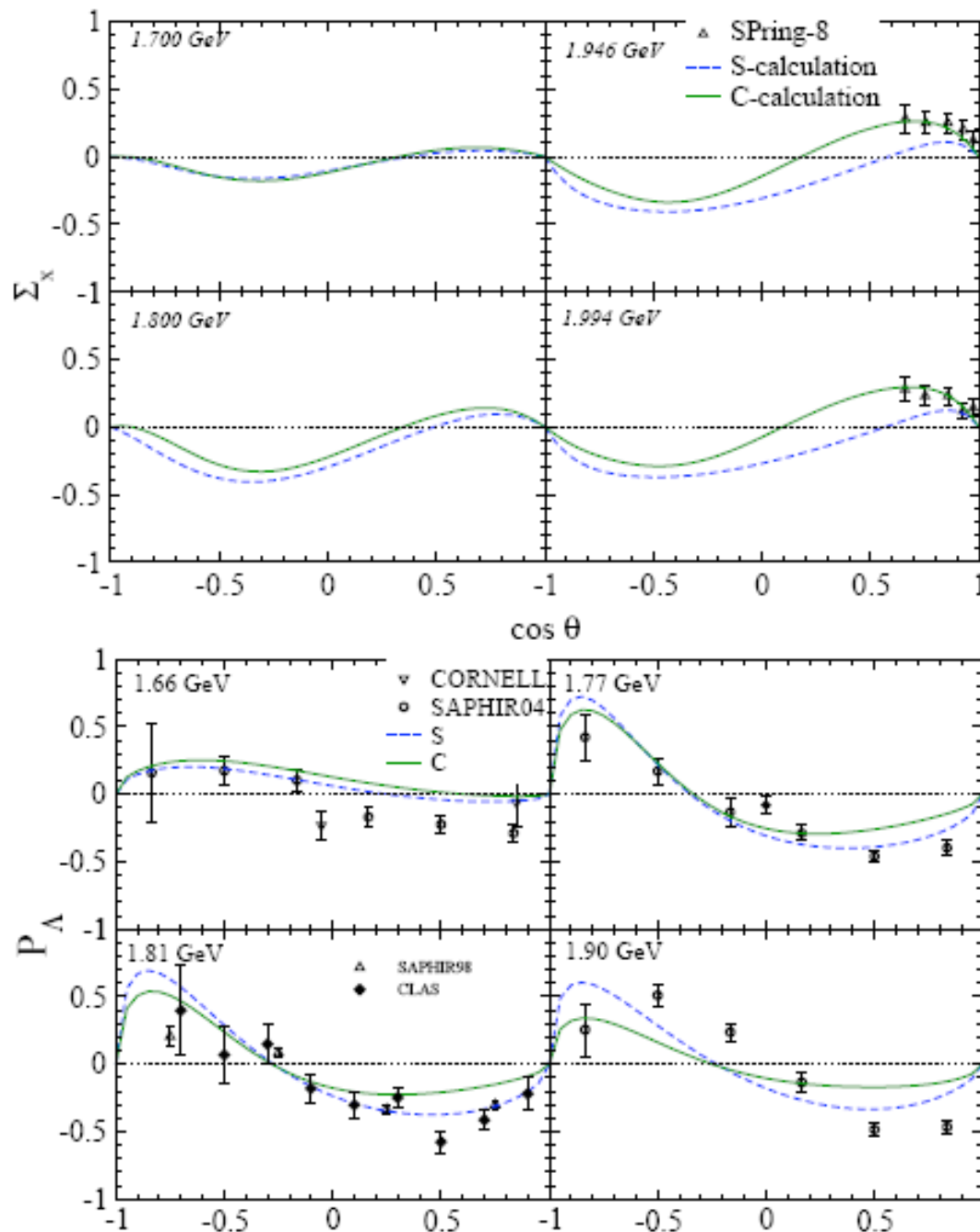


\sqrt{s} (GeV)



$L_{21,2S}$	$R_{K\Lambda}(C)$	$R_{K\Lambda}(S)$
$S_{11}(1535)$	1.3^b	1.26^b
$S_{11}(1650)$	$3.2(+)$	$4.6(+)$
$P_{11}(1440)$	1.48^b	-0.71^b
$P_{11}(1710)$	$6.8(+)$	$3.1(+)$
$P_{13}(1720)$	$4.6(+)$	$4.0(+)$
$P_{13}(1900)$	$2.4(+)$	$2.3(+)$
$D_{13}(1520)$	-0.58^b	-0.33^b
$D_{13}(1950)$	$0.1(+)$	$0.1(-)$
$D_{15}(1675)$	$0.2(+)$	$0.1(+)$
$F_{15}(1680)$	$0.0(+)$	$0.0(+)$
$F_{15}(2000)$	$0.0(+)$	$0.2(-)$

N^* decay ratios to $K\Lambda$ for the C and S calculations.



Predicted photon beam asymmetry.

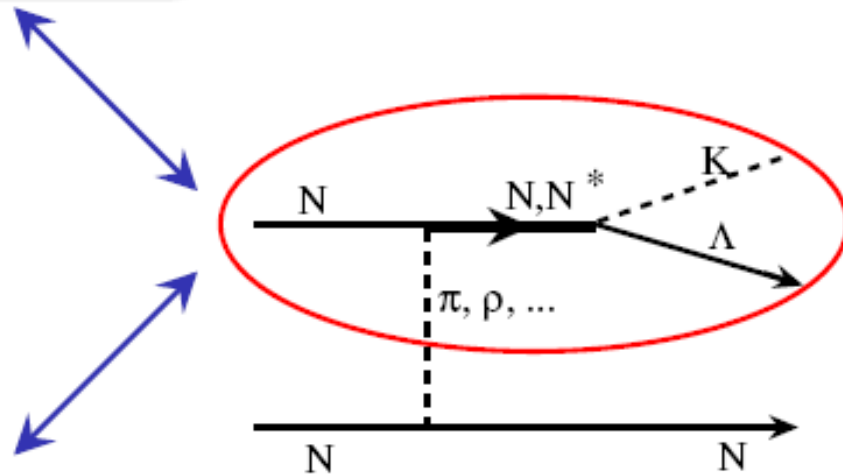
(LEPS data: R. G. T. Zegers *et al.*, Phys. Rev. Lett. 91, 092001 (2003))

Differences between SAPHIR and CLAS cross section data have minor influence on the spin observables! (PRC 72:015210 (2005))

Λ polarization for $\gamma+p \rightarrow K + \Lambda$.

Data: SAPHIR98, SAPHIR04, CLAS, and CORNELL.

- reaction $pp \rightarrow pK^+\Lambda$
- hypernuclei: $A(p, K^+)_{\Lambda}B$



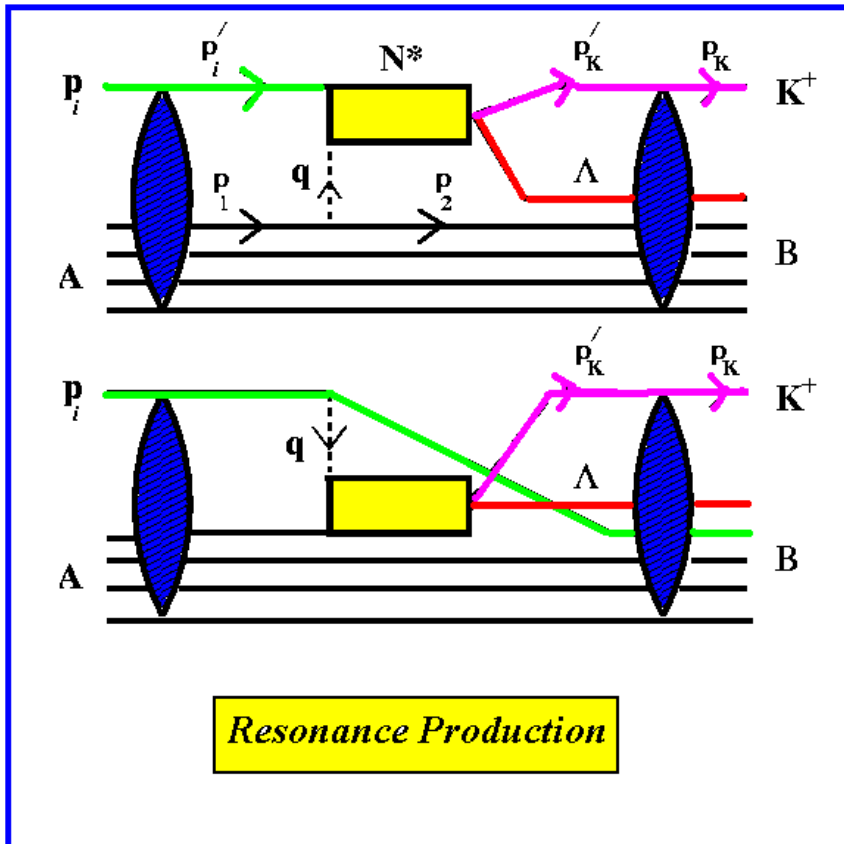
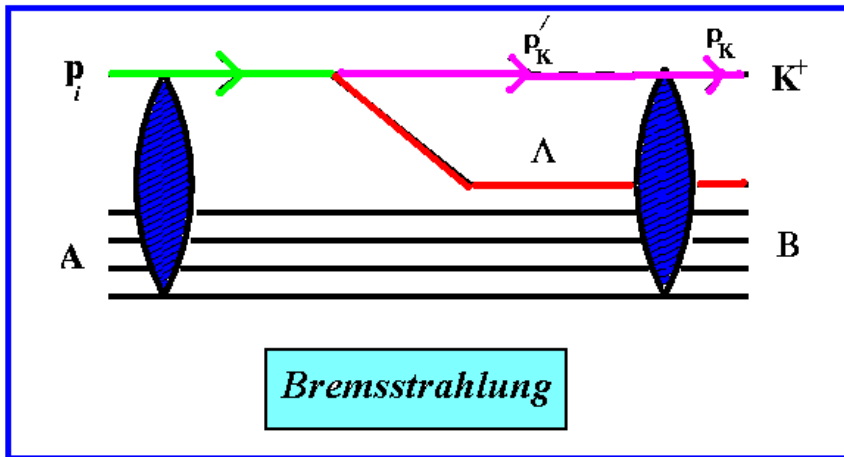
To constrain $pp \rightarrow pK\Lambda$ reaction:
combined analysis of

- $\pi N \rightarrow K\Lambda$
- $\gamma N \rightarrow K\Lambda$

is necessary

Relation to Hadronic $K\Lambda$ Production on the Nucleon and on Nuclei

R. Shyam, H. Lenske,
 U. Mosel:
 nucl-th/0505043 &
 Nucl.Phys. A (2005);
 nucl-th/0308085 &
 Phys.Rev. C69 (2004)
 065205



Associated (p, A) Strangeness Production

- exploratory for hadronic production at COSY
- elementary process for HI production
- extension to electro-production at ELSA, MAMI and JLAB in preparation

$$N^*(1650)[\frac{1}{2}^-], N^*(1710)[\frac{1}{2}^+], N^*(1720)[\frac{3}{2}^+]$$

R. Shyam, H.L., PRC 69: 065205 (2004),
nucl-th/0505043 & NPA (2005) in print

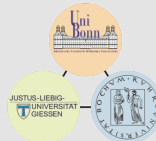
Summary and Outlook

- The Giessen Model: Lagrangian Approach to Coupled Channels K-Matrix
- Coherence of resonance and background contributions
- Simultaneous description of photo- and hadro-production
- Photoproduction and reaction dynamics of ωN and $K\Lambda$ states
- Spectral structures: interference effects or resonances?
- in progress: dynamical correlations in 2-pion channels
- hadro- and electro-production of $K\Lambda$ on nucleons and nuclei

In collaboration with: **V. Shklyar**, U. Mosel



EuroGK



SFB/TR 16



MCTS