

The Interacting Quark Diquark Model

qD model

▪ Diquark

- Two strongly correlated quarks (S wave)

- Baryon in 1_c color representation \rightarrow diquark in $\bar{3}_c$

- Diquark WF: $\square \otimes \square = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} \oplus \begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \Psi_D$ (spin-flavor) antisymmetric
 $6 \otimes 6 = 15 \oplus 21 \rightarrow 15 (A)$ repr. not present

• ~~SU(6) representations for baryons~~

~~$$\begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} \otimes \square = \begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \end{array} \oplus \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array}$$

$$15 \otimes 6 = 20(A) \oplus 70(MA)$$

$$\begin{array}{|c|c|} \hline \square & \square \\ \hline \end{array} \otimes \square = \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} \oplus \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \end{array}$$

$$21 \otimes 6 = 70(MS) \oplus 56(S)$$~~

- Problem of missing resonances

Scalar & axial-vector diquarks

- **21 $SU(6)_{sf}$ representation**

- Decomposed in $SU(2)_s \times SU(3)_f$

- $[\bar{3},0]$ & $[6,1]$ representations. Notation: [flavor,spin]

- **“Good” & “bad” diquarks**

- According to OGE-calculations, $[\bar{3},0]$ is energetically favored
[Wilczek, Jaffe]

- $[\bar{3},0]$: good (scalar) diquark

- $[6,1]$: bad (axial-vector) diquark

Problem of missing resonances

▪ 3-quark QMs

- Excessive number of th. states (much more than experimental ones)
- Several experiments (CLAS, CB-ELSA, TAPS, GRAAL, SAPHIR, etc.) provided no evidence for these states.
- Possible explanation: resonances weakly coupled to the single pion, may decay in 2 (or more) pions/other mesons.

• qD models

- The number of missing resonances decreases notably
- 15 representation for diquark is neglected

Evidences of diquark correlations

- **Regge behavior of hadrons**

Baryons arranged in rotational Regge trajectories ($J=\alpha+\alpha'M^2$) with the same slope of the mesonic ones.

- **$\Delta = 1/2$ rule in weak nonleptonic decays**

Neubert and Stech, Phys. Lett. B **231** (1989) 477; Phys. Rev. D **44** (1991) 775

- **Regularities in parton distribution functions and in spin-dependent structure functions**

Close and Thomas, Phys. Lett. B **212** (1988) 227

- **Regularities in $\Lambda(1116)$ and $\Lambda(1520)$ fragmentation functions**

Jaffe, Phys. Rept. **409** (2005) 1 [Nucl. Phys. Proc. Suppl. **142** (2005) 343]

Wilczek, hep-ph/0409168

- **Any interaction that binds π and ρ mesons in the rainbow-ladder approximation of the DSE will produce diquarks**

Cahill, Roberts and Praschifka, Phys. Rev. D **36** (1987) 2804

- **Indications of diquark confinement**

Bender, Roberts and Von Smekal, Phys. Lett. B **380** (1996) 7

the Interacting qD model

E. Santopinto, PRC72, 022201 (2005)

▪ Hamiltonian

$$H = \frac{p^2}{2m} - \frac{\tau}{r} + \beta r + [B\delta_{S_{12},1} + C\delta_0] +$$
$$+ (-1)^{l+1} 2Ae^{-\alpha r} [(\vec{s}_{12} \cdot \vec{s}_3) + (\vec{t}_{12} \cdot \vec{t}_3) + (\vec{s}_{12} \cdot \vec{s}_3)(\vec{t}_{12} \cdot \vec{t}_3)]$$

- Non-rel. Kinetic energy + Coulomb + linear confining terms
- Splitting between scalar & axial-vector diquarks
- Exchange potential
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Rel. Interacting qD model

J. Ferretti, E. Santopinto & A. Vassallo, PRC83, 065204 (2011)

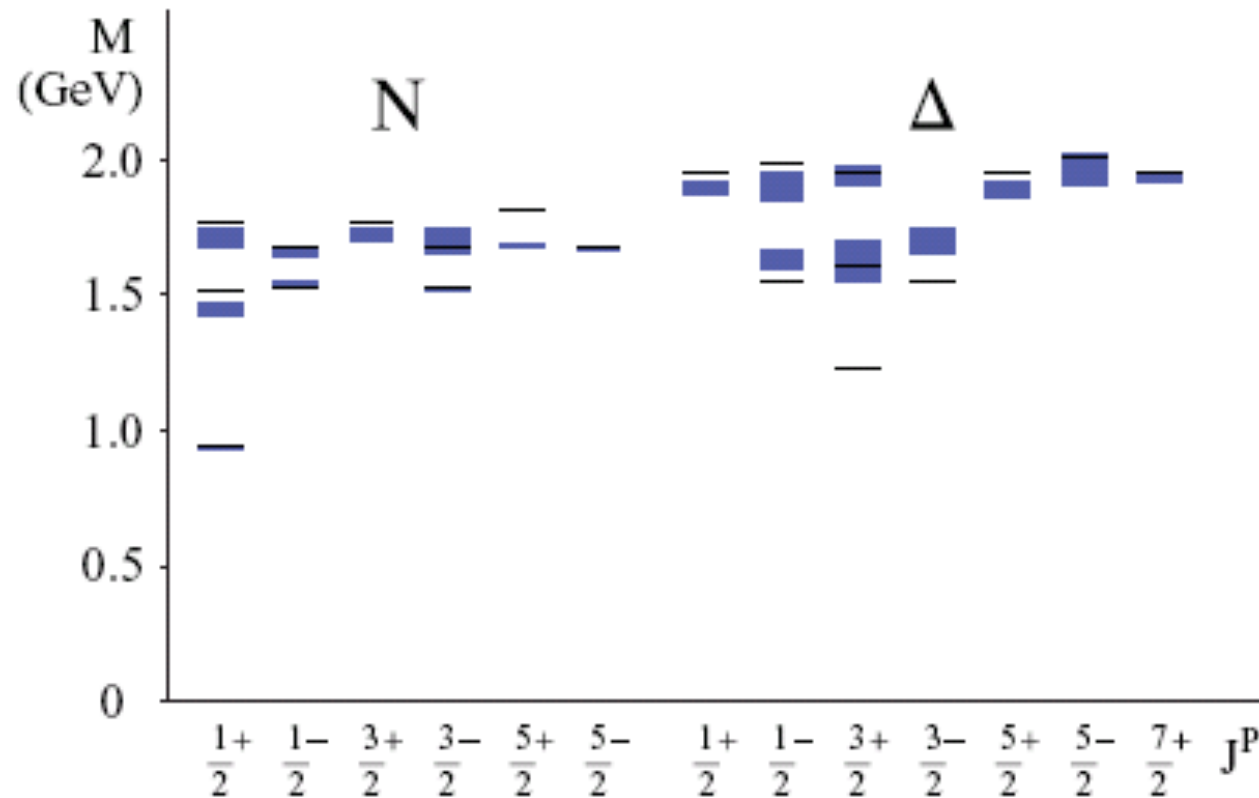
- Relativistic extension of the previous model (point-form formalism).

$$M = E_0 + \sqrt{q^2 + m_1^2} + \sqrt{q^2 + m_2^2} + M_{\text{dir}}(r) + M_{\text{cont}}(r) + M_{\text{ex}}(r),$$
$$M_{\text{dir}}(r) = -\frac{\tau}{r}(1 - e^{-\mu r}) + \beta r.$$
$$M_{\text{ex}}(r) = (-1)^{l+1} e^{-\sigma r} [A_S(\vec{s}_1 \cdot \vec{s}_2) + A_I(\vec{t}_1 \cdot \vec{t}_2) + A_{SI}(\vec{s}_1 \cdot \vec{s}_2)(\vec{t}_1 \cdot \vec{t}_2)],$$
$$M_{\text{cont}} = \left(\frac{m_1 m_2}{E_1 E_2}\right)^{1/2+\epsilon} \frac{\eta^3 D}{\pi^{3/2}} e^{-\eta^2 r^2} \delta_{L,0} \delta_{s_1,1} \left(\frac{m_1 m_2}{E_1 E_2}\right)^{1/2+\epsilon}$$

- Numerical solution with variational program
- Parameters fitted to nonstrange baryon spectrum

Rel. Interacting qD model

J. Ferretti, E. Santopinto & A. Vassallo, PRC83, 065204 (2011)



Resonance	Status	M^{expt} (MeV)	J^P	L^P	S	s_1	n_r	M^{calc} (MeV)
$N(939) P_{11}$	****	939	$\frac{1}{2}^+$	0^+	$\frac{1}{2}$	0	0	939
$N(1440) P_{11}$	****	1420–1470	$\frac{1}{2}^+$	0^+	$\frac{1}{2}$	0	1	1513
$N(1520) D_{13}$	****	1515–1525	$\frac{3}{2}^-$	1^-	$\frac{1}{2}$	0	0	1527
$N(1535) S_{11}$	****	1525–1545	$\frac{1}{2}^-$	1^-	$\frac{1}{2}$	0	0	1527
$N(1650) S_{11}$	****	1645–1670	$\frac{1}{2}^-$	1^-	$\frac{1}{2}, \frac{3}{2}$	1	0	1671
$N(1675) D_{15}$	****	1670–1680	$\frac{5}{2}^-$	1^-	$\frac{3}{2}$	1	0	1671
$N(1680) F_{15}$	****	1680–1690	$\frac{5}{2}^+$	2^+	$\frac{1}{2}$	0	0	1808
$N(1700) D_{13}$	***	1650–1750	$\frac{3}{2}^-$	1^-	$\frac{1}{2}, \frac{3}{2}$	1	0	1671
$N(1710) P_{11}$	***	1680–1740	$\frac{1}{2}^+$	0^+	$\frac{1}{2}$	1	0	1768
$N(1720) P_{13}$	****	1700–1750	$\frac{3}{2}^+$	0^+	$\frac{1}{2}$	1	0	1768
$\Delta(1232) P_{33}$	****	1231–1233	$\frac{3}{2}^+$	0^+	$\frac{1}{2}$	1	0	1233
$\Delta(1600) P_{33}$	***	1550–1700	$\frac{3}{2}^+$	0^+	$\frac{1}{2}$	1	1	1602
$\Delta(1620) S_{31}$	****	1600–1660	$\frac{1}{2}^-$	1^-	$\frac{1}{2}$	1	0	1554
$\Delta(1700) D_{33}$	****	1670–1750	$\frac{3}{2}^-$	1^-	$\frac{1}{2}$	1	0	1554
$\Delta(1900) S_{31}$	**	1850–1950	$\frac{1}{2}^-$	1^-	$\frac{1}{2}$	1	1	1986
$\Delta(1905) F_{35}$	****	1865–1915	$\frac{7}{2}^+$	2^+	$\frac{1}{2}$	1	0	1952
$\Delta(1910) P_{31}$	****	1870–1920	$\frac{3}{2}^+$	2^+	$\frac{1}{2}$	1	0	1952
$\Delta(1920) P_{33}$	***	1900–1970	$\frac{3}{2}^+$	2^+	$\frac{1}{2}$	1	0	1952
$\Delta(1930) D_{35}$	***	1900–2020	$\frac{5}{2}^-$	1^-	$\frac{1}{2}$	1	0	2005
$\Delta(1950) F_{37}$	****	1915–1950	$\frac{7}{2}^+$	2^+	$\frac{1}{2}$	1	0	1952
$N(2100) P_{11}$	*	1855–1915	$\frac{1}{2}^+$	0^+	$\frac{1}{2}$	0	2	1893
$N(2090) S_{11}$	*	1869–1987	$\frac{1}{2}^-$	1^-	$\frac{1}{2}$	0	1	1882
$N(1900) P_{13}$	**	1820–1974	$\frac{3}{2}^+$	2^+	$\frac{1}{2}$	0	0	1808
$N(2080) D_{13}$	**	1740–1940	$\frac{3}{2}^-$	1^-	$\frac{1}{2}$	0	1	1882
$\Delta(1750) P_{31}$	*	1708–1780	$\frac{1}{2}^+$	0^+	$\frac{1}{2}$	1	0	1858
$\Delta(1940) D_{33}$	*	1947–2167	$\frac{3}{2}^-$	1^-	$\frac{1}{2}$	1	1	1986

**0 missing resonances
below 2 GeV**

Model Parameters

$m_q = 200 \text{ MeV}$	$m_S = 600 \text{ MeV}$	$m_{AV} = 950 \text{ MeV}$
$\tau = 1.25$	$\mu = 75.0 \text{ fm}^{-1}$	$\beta = 2.15 \text{ fm}^{-2}$
$A_S = 375 \text{ MeV}$	$A_I = 260 \text{ MeV}$	$A_{SI} = 375 \text{ MeV}$
$\sigma = 1.71 \text{ fm}^{-1}$	$E_0 = 154 \text{ MeV}$	$D = 4.66 \text{ fm}^2$
$\eta = 10.0 \text{ fm}^{-1}$	$\epsilon = 0.200$	

Rel. Interacting qD model – strange B.

E. Santopinto & J. Ferretti, arXiv: 1412.7571

▪ Model

- Model extended to strange sector
- Hamiltonian:

$$M = E_0 + \sqrt{\vec{q}^2 + m_1^2} + \sqrt{\vec{q}^2 + m_2^2} + M_{\text{dir}}(r) + M_{\text{ex}}(r)$$

$$M_{\text{ex}}(r) = (-1)^{L+1} e^{-\sigma r} [A_S \vec{s}_1 \cdot \vec{s}_2 + A_F \vec{\lambda}_1^f \cdot \vec{\lambda}_2^f + A_I \vec{t}_1 \cdot \vec{t}_2]$$

$$M_{\text{dir}}(r) = -\frac{\tau}{r}(1 - e^{-\mu r}) + \beta r.$$

- Gursev-Radicati inspired exchange interaction
- Parameters fitted to strange baryon spectrum

Rel. Interacting qD model – strange B.

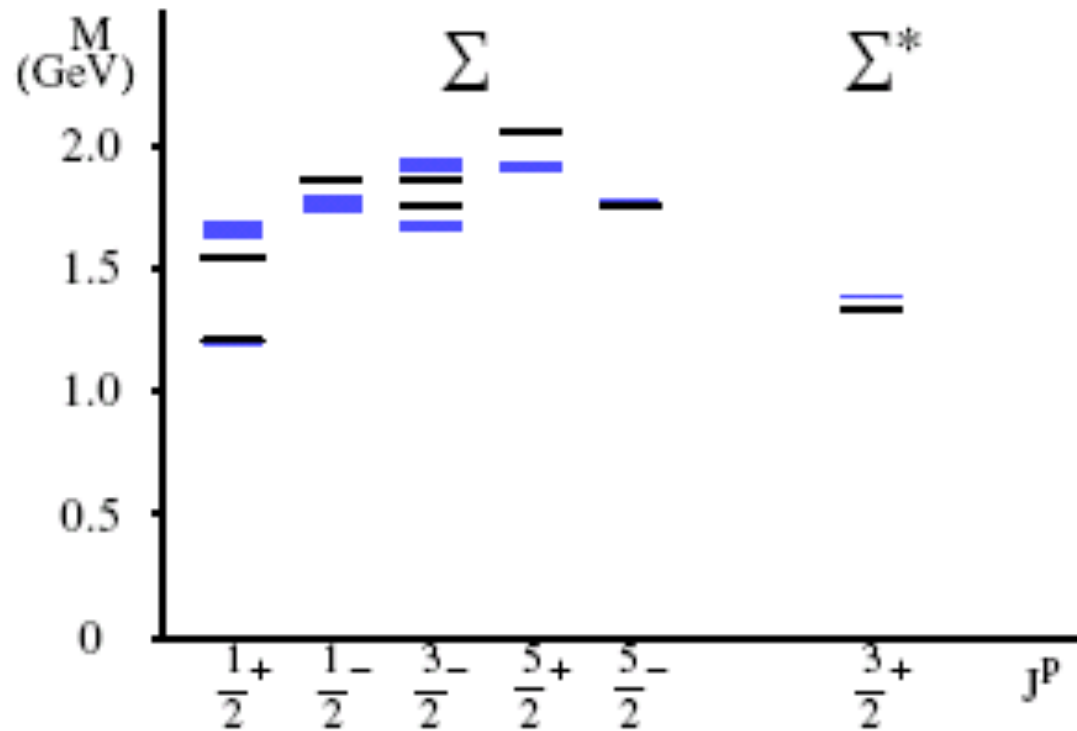
E. Santopinto & J. Ferretti, arXiv: 1412.7571

Parameters

Parameter	Value (Fit 1)	Value (Fit 2)	Parameter	Value (Fit 1)	Value (Fit 2)
m_n	200 MeV	159 MeV	m_s	550 MeV	213 MeV
$m_{[n,n]}$	600 MeV	607 MeV	$m_{[n,s]}$	900 MeV	856 MeV
$m_{\{n,n\}}$	950 MeV	963 MeV	$m_{\{n,s\}}$	1200 MeV	1216 MeV
$m_{\{s,s\}}$	1580 MeV	1352 MeV	τ	1.20	1.02
μ	75.0 fm ⁻¹	28.4 fm ⁻¹	β	2.15 fm ⁻²	2.36 fm ⁻²
A_S	350 MeV	-436 MeV	A_F	100 MeV	193 MeV
A_I	250 MeV	791 MeV	σ	2.30 fm ⁻¹	2.25 fm ⁻¹
E_0	141 MeV	150 MeV	ϵ	0.37	–
D	6.13 fm ²	–	η	11.0 fm ⁻¹	–

Rel. Interacting qD model – strange B.

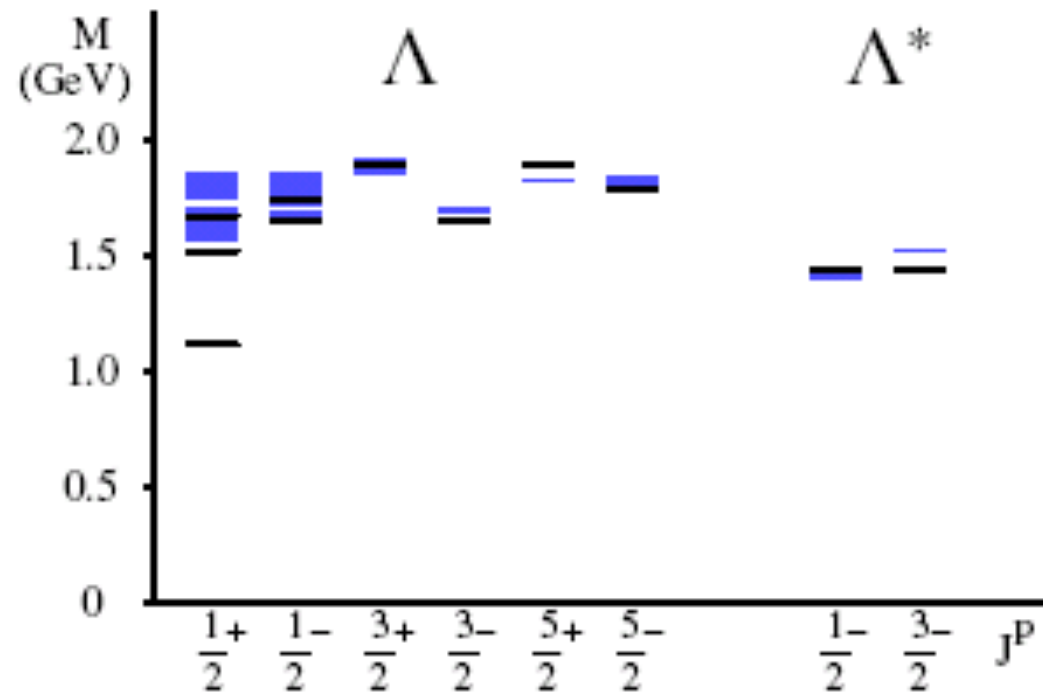
E. Santopinto & J. Ferretti, arXiv: 1412.7571

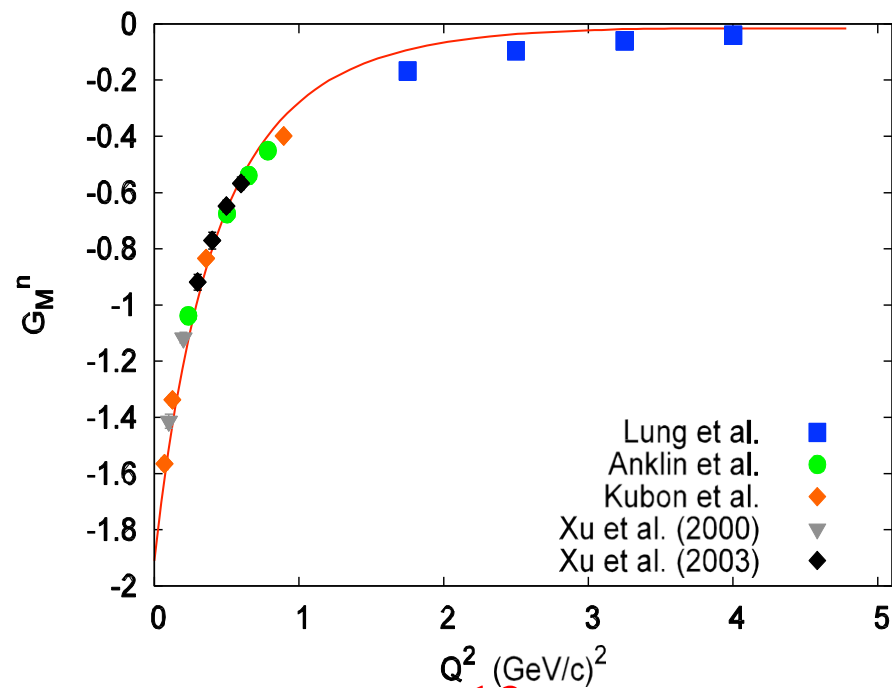
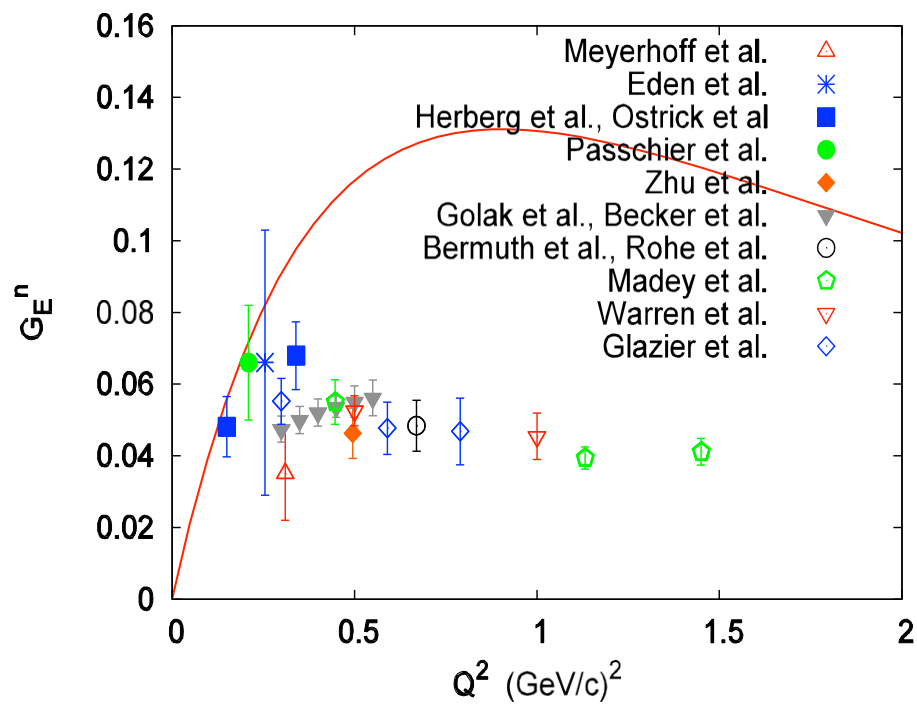
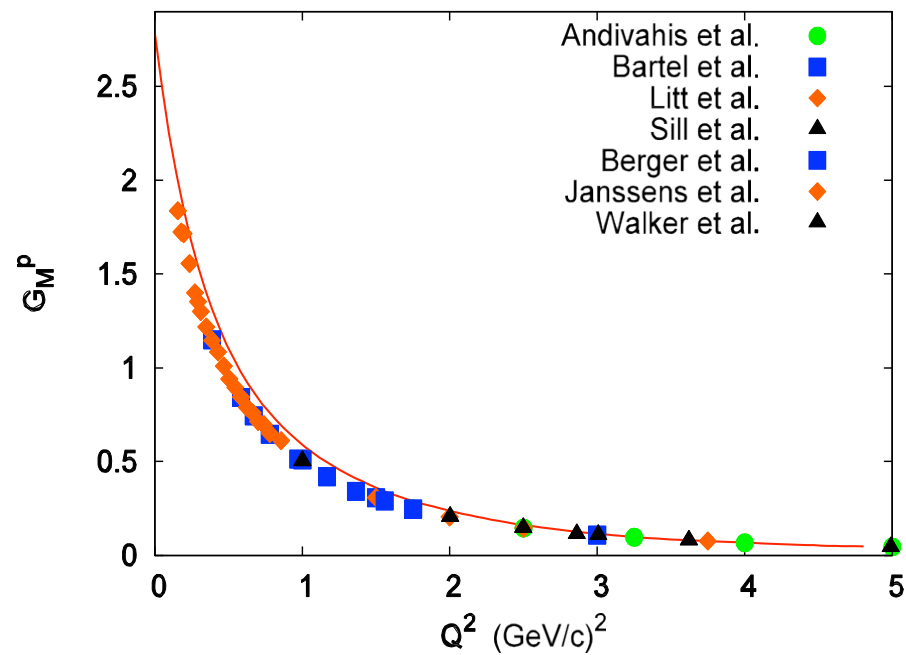
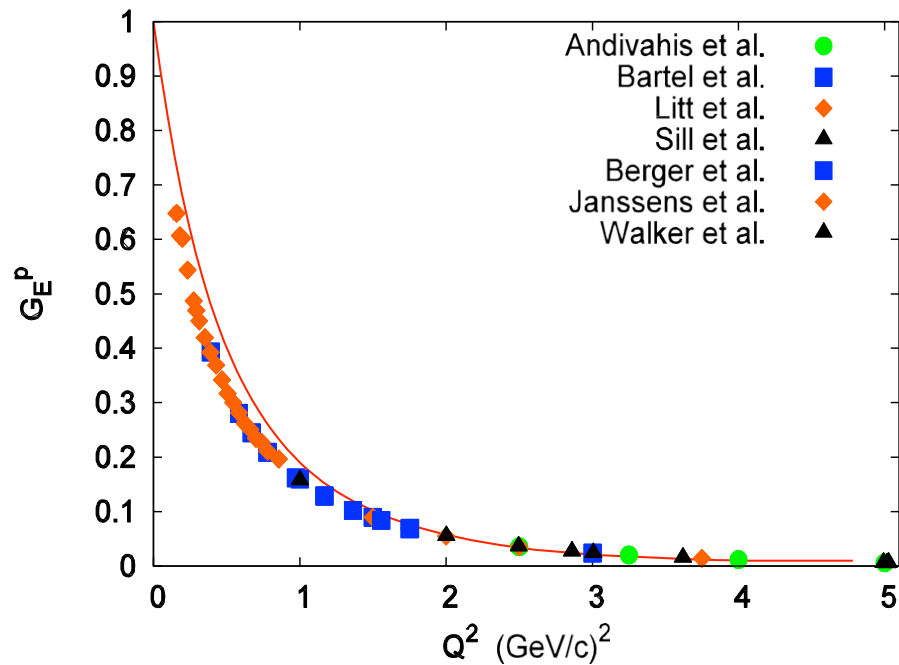


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E. Santopinto & J. Ferretti, arXiv: 1412.7571

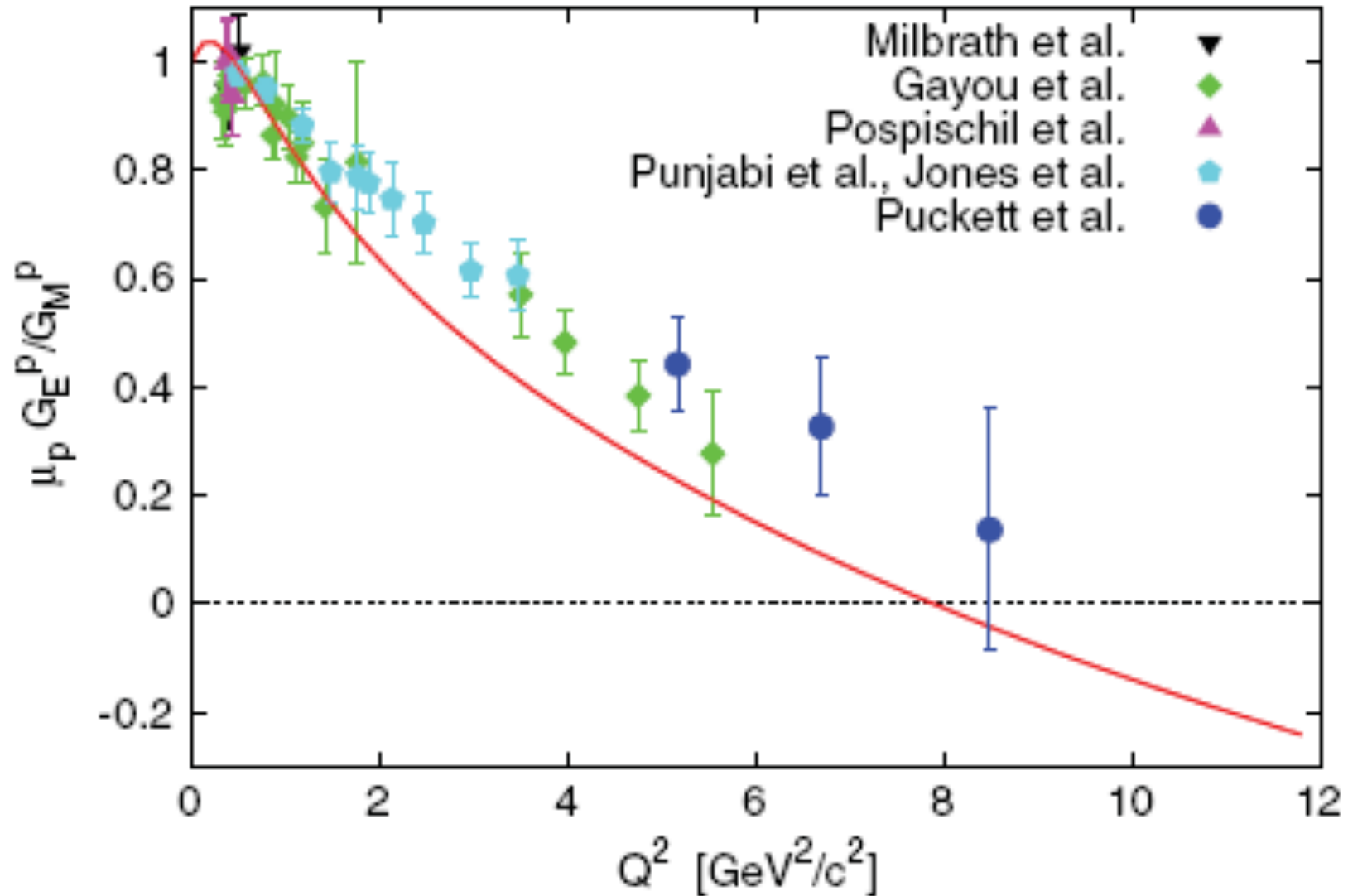
■ Lambda & Lambda* states





Ratio $\mu_p G_E^p/G_M^p$

De Sanctis, Ferretti, Santopinto, Vassallo, Phys. Rev. C 84, 055201 (2011)



Interacting Quark Diquark model , *E. Santopinto, Phys. Rev. C 72, 022201(R) (2005)*

QCD shows no q-diquark

evidence of the type of the Lichtenberg Model.

q-diqu. m. does not mix different diquark configurations.

**D. B. Lichtenberg and L. J. Tassie, Phys. Rev. 155, 1601
(1967).**

$$M_{\text{tr}}(r) = V_0 e^{-\frac{1}{2}\nu^2 r^2} (\vec{s}_2 \cdot \vec{S})(\vec{t}_2 \cdot \vec{T}) ,$$

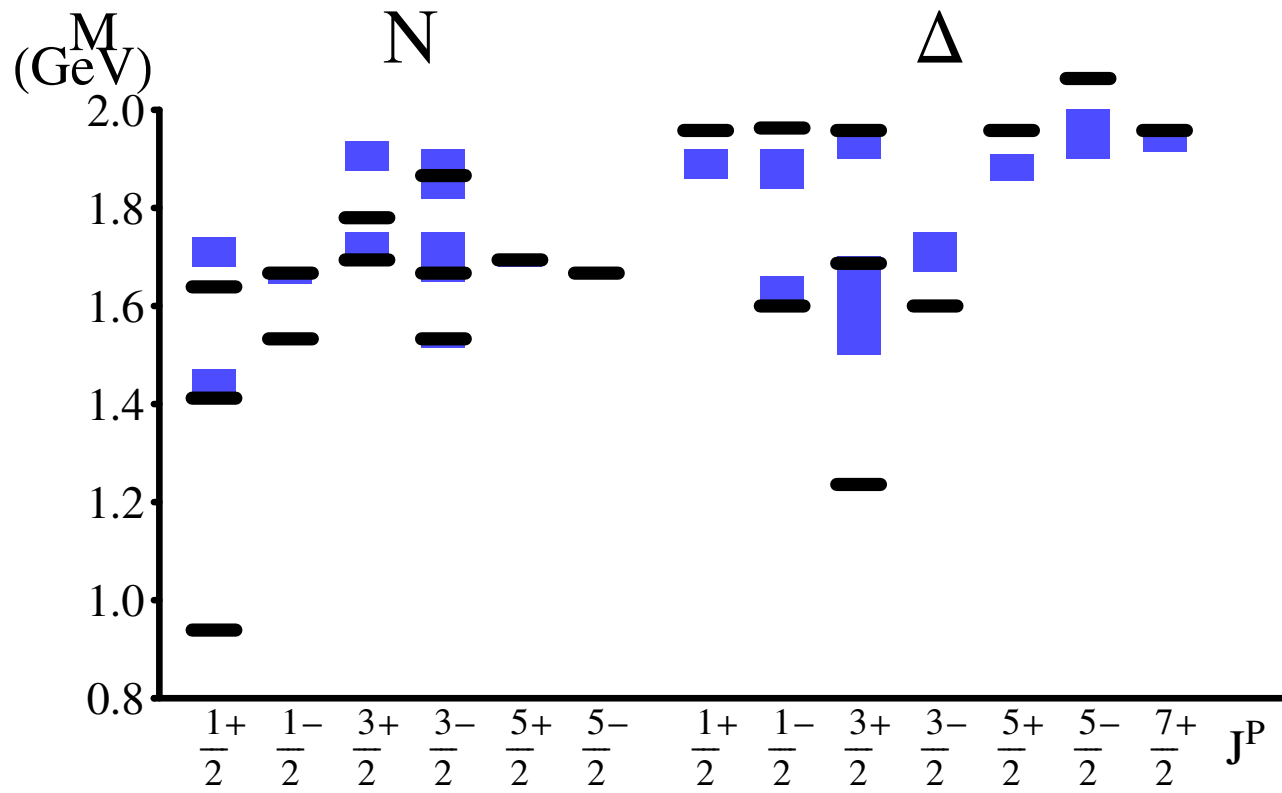
$$\langle s'_1, m'_{s_1} | S_\mu^{[1]} | s_1, m_{s_1} \rangle \neq 0 \text{ for } s'_1 \neq s_1 ,$$

$$\langle 1 || S_1 || 0 \rangle = 1 ,$$

$$\langle 0 || S_1 || 1 \rangle = -1$$

$$\begin{aligned} \langle \Phi' | M_{\text{tr}} | \Phi \rangle &= \frac{1}{4} V_0 \delta_{s'_1, s_1 \pm 1} \delta_{S \frac{1}{2}} \delta_{t'_1, t_1 \pm 1} \delta_{T \frac{1}{2}} \\ &\times \langle \Phi'(\vec{r}) | e^{-\frac{1}{2}\nu^2 r^2} | \Phi(\vec{r}) \rangle , \end{aligned}$$

$$|N\rangle = a_S |qD_S, L = 0\rangle + a_{AV} |qD_{AV}, L = 0\rangle$$



Thank you for your attention