

Day 6

① Real analytic function (Hermitian analytic)

(z Domain $D := \{s, \text{Im } s > 0\}$)

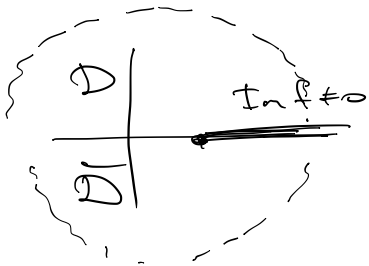


$f(z)$ is analytic, $\text{Im } f(z) \neq 0$ at the border of the domain $z \in \text{Real}$

\rightarrow Schwarz reflection principle $f(z)$ is analytic in $D \cup \bar{D}$

$f(z^*) = f^*(z)$, $\rightarrow f$ is Hermitian analytic

if $f(z)$ has region $\text{Im } f(x) \neq 0$, $x \in [a, b)$ in border

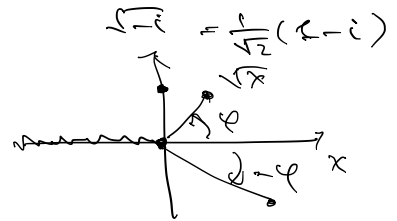
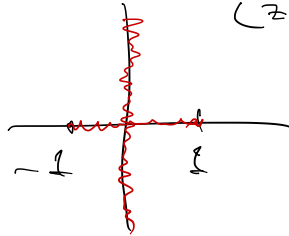


Example: $f(z) = \sqrt{z^2 - 1}$

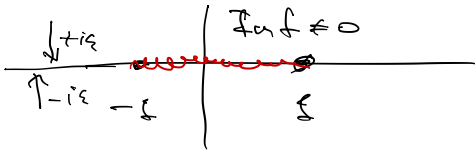
(z)

$$\sqrt{1} = 1, \sqrt{-1} = \frac{1}{\sqrt{2}}(1+i)$$

$$\sqrt{-i} = \frac{1}{\sqrt{2}}(1-i)$$



$$b) f(z) = \sqrt{z-1} \sqrt{z+1}$$



② Dispersion relation

$T(s, t)$ is scattering amplitude
for t to be negative, in s -channel

$$s_{th} = 4m^2, \quad t \rightarrow -! \quad s + t + u = \sum p_i^2 = 4m^2$$

$$u_{th} = 4m^2 \Rightarrow s_{th} = -t$$

$T(s, t) = f(s)$ real analytic

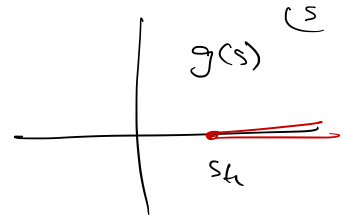
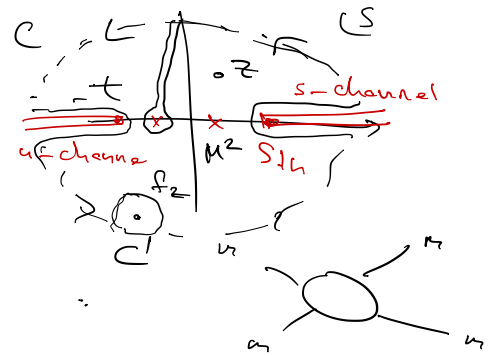
Cauchy theorem
$$f(z) = \frac{1}{2\pi i} \int_C \frac{f(z')}{z' - z} dz'$$

$$F(z', z) = \frac{f(z')}{z' - z}$$

$$\frac{1}{2\pi i} \int_C F(z', z) dz' = \sum_{z_i, z_j, z'_j} \text{res } F(z', z) z_p = M^2, \text{ bound state in } s\text{-channel}$$

$$f(z) + \frac{g}{M^2 - z} + \frac{g'}{4m^2 - t - M^2 - z} =$$

$$= \frac{1}{2\pi i} \left[\int_{\gamma_1} + \int_{\gamma_2} + \int_{C_\infty} \right] \frac{f(z')}{z' - z} dz'$$

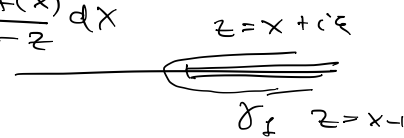


$$z'_p = 4m^2 - t - M^2 = s'_p$$



if $|f(z)| \rightarrow 0$ when $z \rightarrow \infty$ then $\int_{\gamma_\infty} \frac{f(z')}{z'-z} dz' = 0$

$$\Rightarrow f(z) = \frac{g}{M^2 - z} + \frac{g'}{4M^2 - M^2 - z} + \frac{1}{2\pi i} \int_{\gamma_1} + \int_{\gamma_2} \frac{f(z')}{z'-z} dz'$$

$$\int_{\gamma_1} \frac{f(z)}{z'-z} dz = \int_{th}^{+\infty} dx \frac{f(x+i\epsilon) - f(x-i\epsilon)}{x-z} = \int_{th} \frac{\text{Disc } f(x)}{x-z} dx$$


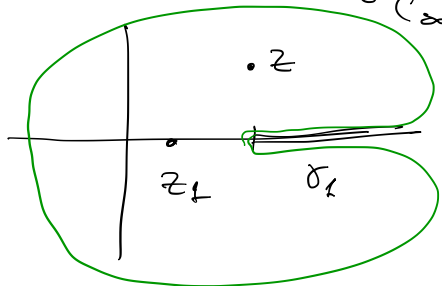
$$\text{Disc } f(x) = f(x+i\epsilon) - f(x-i\epsilon)$$

$$f(z) \text{ is real analytic} \Rightarrow \text{Disc } f(x) = 2i \text{Im } f(x)$$

$$f(z) = \frac{g}{M^2 - z} + \frac{g'}{4M^2 - M^2 - z} + \frac{1}{\pi} \int_{th}^{+\infty} \frac{\text{Im } f(x)}{x-z} dx + \frac{1}{\pi} \int_{-\infty}^{th'} \frac{\text{Im } f(x)}{x-z} dx$$

③ Subtraction

What if $f(z) \neq 0$, $g(z) \neq 0$
 $z \rightarrow \infty$



$$F(z', z, z_1) = \frac{f(z')}{(z' - z)(z' - z_1)}$$

$$G(z', z, z_1) = \frac{g(z')}{(z' - z)(z' - z_1)}$$

$$\frac{1}{2\pi i} \int_C G(z', z, z_1) dz' = \sum \text{res}$$

$$\Rightarrow \text{Now } \left| \frac{g(z')}{z' - z_1} \right| \xrightarrow{z' \rightarrow \infty} 0$$

$$\frac{1}{\pi} \int_{\text{th}} \frac{\text{Im} g(x)}{(x - z)(x - z_1)} dx = \frac{g(z)}{z - z_1} + \frac{g(z_1)}{z_1 - z} \Rightarrow$$

$$g(z) = g(z_1) + \frac{(z - z_1)}{\pi} \int_{\text{th}} \frac{\text{Im} g(x)}{(x - z)(x - z_1)} dx$$

Why subtraction?

$$g(z) = \frac{1}{2\pi i} \int_C \frac{g(z')}{z' - z} dz'$$

$$g(z_1) = \frac{1}{2\pi i} \int_C \frac{g(z')}{z' - z_1} dz'$$



$$g(z) - g(z_1) = \frac{z - z_1}{2\pi i} \int_C \frac{g(z')}{(z' - z)(z' - z_1)} dz'$$

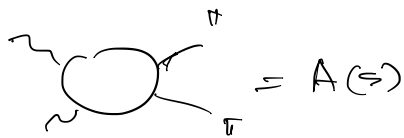
Multiple subtraction

$$F(z', z, z_1, \dots, z_n) \sim \frac{f(z')}{(z'-z)(z'-z_1)\dots}$$

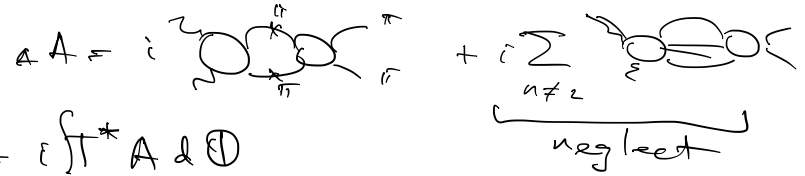
$$\left| \frac{f(z')}{(z'-z_1)\dots(z'-z_n)} \right| \begin{array}{l} |z| \rightarrow \infty \\ \rightarrow 0 \end{array}$$

$$g(z) = g(0) + z g'(0) + \frac{z^2}{2} \int \frac{g(z')}{(z'-z) z'^2} dz'$$

④. Omnes functions



weakly coupled channel



$$\Delta A = i \int T^* A d\Omega$$

partial waves $\Delta f_\ell = 2i f_\ell^* g f_\ell$, $g = \frac{1}{2} \Phi = \frac{1}{16\pi} \sqrt{1 - \frac{4m^2}{s}}$

$$T(s, \theta) = \sum (2\ell + 1) f_\ell(s) P_\ell(\cos \theta)$$

elastic PW amplitude

$$\begin{cases} \Delta b_\ell = 2i f_\ell^* g f_\ell \\ \Delta f = 2i f_\ell^* g f_\ell \end{cases}$$

$$\rightarrow f_\ell = a_\ell + b_\ell \leftarrow \text{only left hand cut}$$

has only right hand cut

$$\Delta a_\ell = 2i f_\ell^* g (a_\ell + b_\ell) \leftarrow \text{Dunne's-like equation}$$

$a_\ell = P_\ell \Omega_\ell$ → phase function, defined by t
 \uparrow defined by b_ℓ

$$\Delta a_e = z i f^* g(a_e + b_e) \quad \rightarrow \quad a^+ - a^- = z i f^* g(a^+ + b)$$

$$a_e(s+i\varepsilon) = a_e^+, \quad a_e(s-i\varepsilon) = a_e^-$$

$$\Delta \Omega = (\Omega_+ - \Omega_-) = z i g^* \Omega(t)$$

$$\Rightarrow a^+(1 - z i g^* f^*) = a^- + z i g^* f^* b \quad \leftarrow \text{inhomogeneous}$$

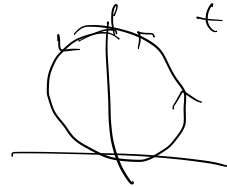
Onnes equations

Homogeneous equation

$$a^+(1 - z i g^* f^*) = a^- \quad , \quad t = \frac{1}{z g i} [e^{2i\delta} - 1]$$

$$a^+ e^{-2i\delta} = a^- \quad \leftarrow$$

$$\ln a^+ - 2i\delta = \ln a^- \quad \Rightarrow \quad \Delta \ln a^+ = 2i\delta$$

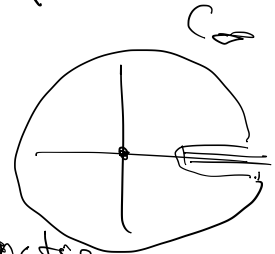


$$\Rightarrow \ln a = + \frac{\beta}{\pi} \int_{t_0}^+ \frac{\delta(s')}{(s'-s)s'} ds' + \ln a(0)$$

subtraction
const

$$a = a(0) \exp\left(+ \frac{\beta}{\pi} \int_{t_0}^+ \frac{\delta(s')}{s'-s} ds'\right)$$

$$\Omega = \Omega(0) \exp\left(+ \frac{\beta}{\pi} \int_{t_0}^+ \frac{\delta(s')}{s'-s} ds'\right)$$



Onnes function

• same phase as t
• no Lft

Solution of inhomogeneous equations

$$a^+(1 - 2\epsilon g f^+) = a^- + 2\epsilon g b f^+ \quad , \quad a = P \Omega$$

↑ function

$$P^+ \Omega^+ (1 - 2\epsilon g f^+) = P^- \Omega^- + 2\epsilon g b f^+$$

$$\Rightarrow \Omega^- P^+ = P^+ \Omega^- + 2\epsilon g b f^+ \Rightarrow$$

$$\Rightarrow A P = 2\epsilon g b f^+ / \Omega^- \Rightarrow P = \frac{f}{\pi} \int \frac{b f^+}{\Omega(s'-s)} ds +$$

Solution $a(s) = \Omega(s) \left[\begin{matrix} \text{regular} \\ \text{function} \end{matrix} + \frac{f}{\pi} \int \frac{b f^+}{\Omega(s'-s)} ds' \right] + \text{regular function}$

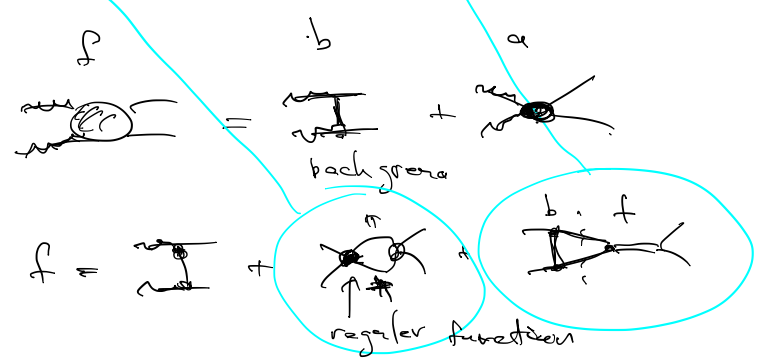
if b is not present $\Rightarrow a(s) = \Omega(s) \cdot \text{regular function}$

$$f(s) = a(s) + b(s)$$

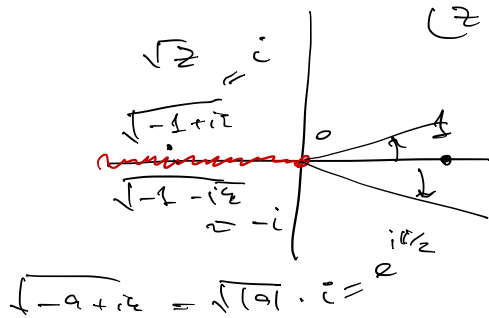
↑ $\pm H.C.$

↑ only R.H.E

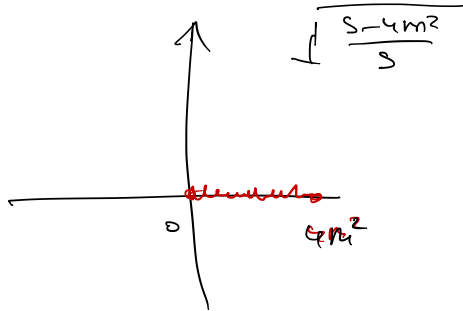
Linearization \longrightarrow



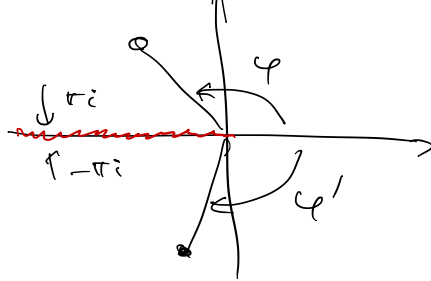
Solution 6.5



$$\sqrt{-1+i} = \sqrt{|z|} \cdot i = e$$



$$\ln z = \ln |z| + i \text{Arg } z$$



$$\sqrt{x} = 0 + \frac{x}{2\pi i} \int_{-\infty}^{\infty} \frac{\Delta \sqrt{x'}}{(x'-x)x'} dx'$$

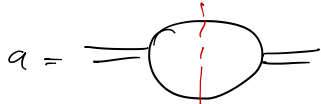
$$\Delta \sqrt{x} = 2\sqrt{x}$$

$$\sqrt{x} = 0 + \frac{x}{\pi} \int_{-\infty}^{\infty} \frac{\sqrt{-x'}}{(x'-x)x'} dx'$$

$$\Delta \ln x = 2\pi i$$

$$\ln x = \ln 1 + \frac{(x-1)}{2i\pi} \int_{-\infty}^{\infty} \frac{2\pi i dx'}{(x'-x)(x'-1)}$$

Scalar loop



$$a = \int \frac{d^4k}{(2\pi)^4} \frac{1}{(k^2 - m^2)(k^2 - p^2)}$$

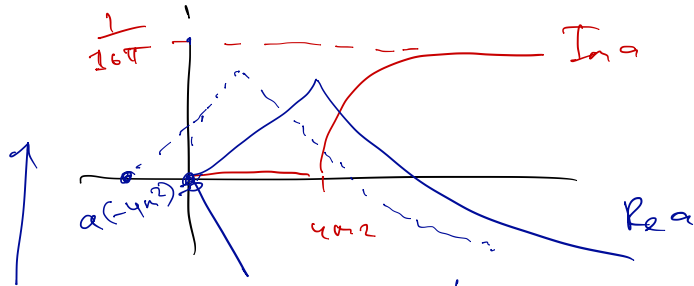
integral diverges

$$\text{Im } a = \frac{1}{2} \circledast = \frac{1}{16\pi} \sqrt{1 - \frac{4m^2}{s}}, \quad s > 4m^2$$

$$= 0, \quad \text{when } s \leq 4m^2$$

$$a = a(0) + \frac{1}{\pi} \int \frac{\rho(s')}{(s' - s)s'} ds' = \Sigma(s)$$

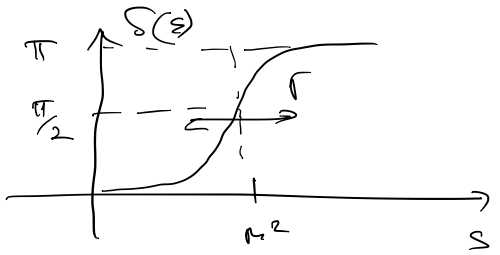
renormalization



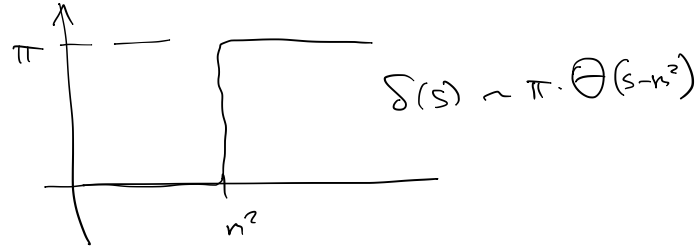
shift is determined by subtraction constant

6.2. Solution

$$F_{BW} = \frac{g^2}{m^2 - s - i\epsilon \Gamma(s)}, \quad \arg F_{BW} = \arctan \frac{\text{Im} F_{BW}}{\text{Re} F_{BW}} = \arctan \frac{m \Gamma}{m^2 - s}$$



$\Gamma \rightarrow 0$



$$\Omega = \Omega(s) \exp\left(-\frac{\Gamma s}{\pi} \int \frac{\Theta(s' - m^2)}{s'(s' - s)} ds'\right) =$$

$$= \Omega(s) \exp\left(-s \int_{m^2}^{+\infty} \frac{ds'}{(s' - s)s'}\right) = \frac{1}{m^2 - s}$$

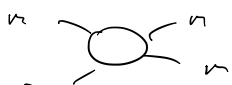
→ We have got
Breit-Wigner

back

$$\approx \log(m^2 - s)$$

→ omnes is just t without LHC.

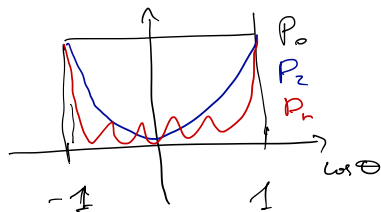
Lecture: High energy for partial waves


 $A(s, t) = \sum (2\ell + 1) f_\ell P_\ell(z_s), \quad z = \cos \Theta$

take s to be large, t is small and fixed

f) P_ℓ for high ℓ -?

$$\cos \Theta = 1 + \frac{2t}{s - 4a^2} \rightarrow \text{I when } s \rightarrow \infty$$



$P_n(\cos \Theta)$ has peaks forward and backward

$$\rightarrow \Theta \sim \frac{1}{k_{\max}}$$

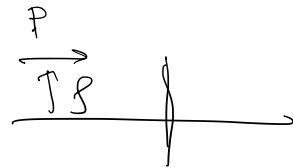
Impact parameter representation

$$P_\ell(\cos \Theta) \Rightarrow J_0(\ell \Theta) \text{ for small } \Theta$$

$$A(s, t) = \sum (2\ell + 1) f_\ell P_\ell(\cos \Theta) \rightarrow \int \rho d\ell f_\ell(\Theta) J_0(\ell \Theta)$$

$$f_\ell(s) = f(\ell, s) \text{ w/ amplitude}$$

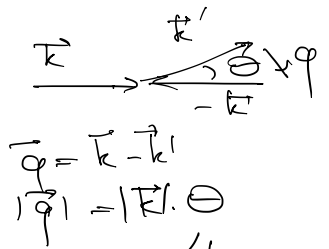
$$b = \frac{\sqrt{s}}{2} = \ell \Rightarrow s = \frac{2\ell}{k}, \quad k = \frac{\sqrt{s}}{2}$$



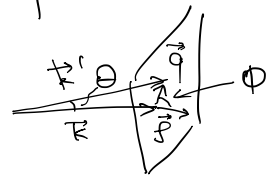
$f(\xi, \eta)$ is impact parameter representation of PW

$$A(s, t) = 2 \int_{-\infty}^{\infty} d\ell \cdot \ell f(\ell, s) J_0(k\ell\theta) \Rightarrow$$

$$= 2k^2 \int d\xi \xi f(\xi, s) J_0(q\xi)$$



$$\int_a^b f(x) dx = \sum f(x_i) \Delta x$$



$$J_0(q\xi) = \int_0^{2\pi} e^{+i\vec{q}\cdot\vec{\xi}} \frac{d\phi}{2\pi}$$

$$t = -q^2$$

$$A(s, t) = k^2 \int d\xi d\eta f(\xi, \eta) \int \frac{d\phi}{2\pi} e^{+i\vec{\xi}\cdot\vec{q}}$$

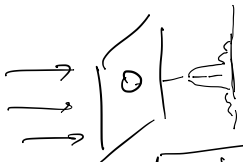
$$\Rightarrow \frac{k^2}{\pi} \int d\xi_x d\xi_y f(\xi, \eta) e^{+i\vec{\xi}\cdot\vec{q}} = A(s, -q^2)$$

$$|\vec{k}| = \frac{\sqrt{s}}{2}$$

Very similar to optics, Fraunhofer scattering

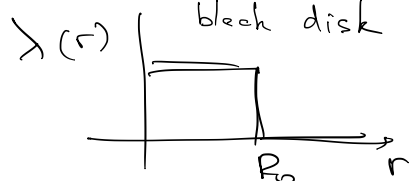


$$A = \int_S dx dy e^{i\varphi(r)} = \int \underbrace{\lambda(x, y)}_{\text{density of sources}} e^{+i\vec{q}\cdot\vec{r}} dx dy$$



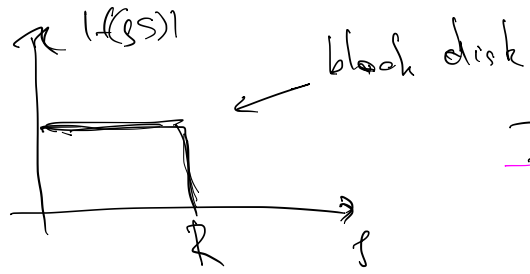
$$I(q) = I_0(q)$$

Hold disk



$$A(\vec{q}) + A_0(\vec{q}) = 0 \Rightarrow A(\vec{q}) = -A_0(\vec{q})$$

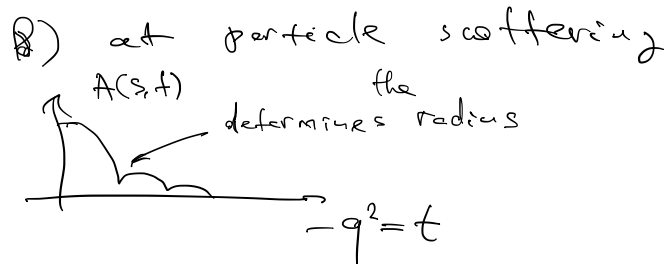
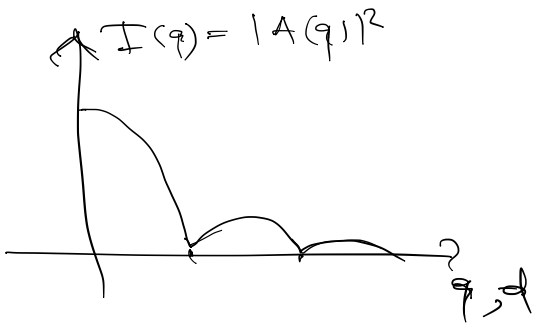
$f(\vec{q}, s)$ can be viewed as density just behind the target.



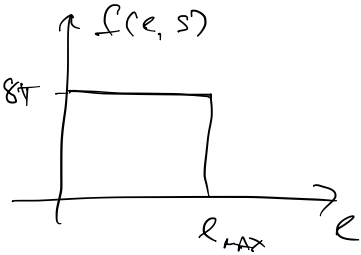
Interference picture.

a) optics $\int_0^R dx dy e^{iq} = \frac{R I_0(Rq)}{q}$

$(q) = |K| \cdot \Theta \Rightarrow \Theta = \frac{|\vec{q}|}{|\vec{k}|}, d = L \cdot \Theta$



Cross-section



$$A(r, t) = \sum_0^{r_{max}} (2r+1) f_r P_r(\theta)$$

$$\sigma = \frac{1}{s} \int_0^s A(r, 0) \quad , \quad A(r, 0) = \sum f_r (2r+1)$$

$t=0, \cos \theta = 1$

$$\sigma = \frac{1}{s} \cdot 8\pi \cdot \sum_{r=0}^{r_{max}} (2r+1) =$$

$r_{max} = R \cdot \frac{\sqrt{s}}{2}$

$$= \frac{1}{s} 8\pi (R_{max} + 1)^2$$

$$= \frac{1}{s} \cdot 8\pi \frac{s}{4} R^2 = \boxed{2\pi R^2}$$

← Twice area of the disc

$$I_{in} A = \frac{1}{2} \sum_{r=0}^{r_{max}} \text{diagram} = \frac{1}{2} \sum_{r=0}^{r_{max}} \text{diagram} + \frac{1}{2} \sum_{r=0}^{r_{max}} \text{diagram}$$

elastic

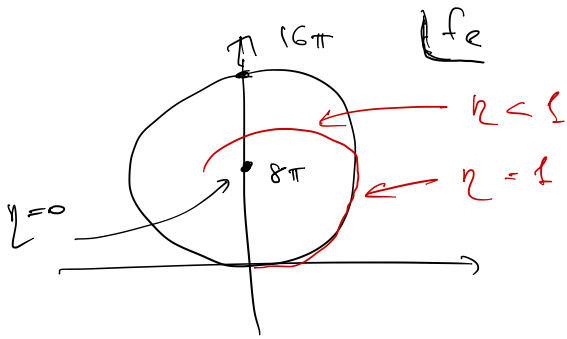
Pur expansion

$$I_{el} r^2 = I_{in} r^2 + (R_{el} r_{el})^2$$

$$I_{in} f_r = g f_r^* f_r + \Delta r \quad , \quad \frac{1}{2} \sum_{r=0}^{r_{max}} \text{diagram} = \sum \Delta r P_r(2r+1)$$

$\Delta r > 0$

$$\rightarrow f_r = \frac{\eta R^{2\delta} - f}{2ig} \quad , \quad \eta \text{ is elasticity } \eta^2 = f - \epsilon g \Delta$$



black disk $\rightarrow \eta = 0$

$$\text{Im } f_e = \sum f_e^2 + \Delta e$$

\swarrow inelastic
 \nwarrow elastic

$$\sigma = \frac{1}{s} \text{Im } A = \frac{1}{s} \sum (2l+1) \text{Im } f_e$$

$$\sigma = \sigma_d + \sigma_{inel} ; \quad \sigma_{elastic} = \frac{1}{s} \sum (2l+1) \text{Im } f_e =$$

$$= \frac{1}{s} \sum_{l=0}^{l_{max}} (2l+1) \int |f_e|^2 =$$

$$\sigma_{inelastic} = \frac{1}{s} \sum_{l=0}^{l_{max}} (2l+1) \Delta e = \frac{1}{s} \sum_{l=0}^{l_{max}} (2l+1) [1 - \eta^2] \frac{1}{4\beta}$$

$$\sigma_{elastic} = \frac{1}{s} \sum (2l+1) [1 + \eta^2 - 2 \cos \delta \eta] \frac{1}{4\beta}$$

$$\eta = 0 \rightarrow \sigma_{el} = \sigma_{in} = \pi R^2 \Rightarrow \sigma_{TOT} = 2\pi R^2$$

$$\rightarrow \mathcal{N}^+ - \mathcal{N}^- = \mathcal{L} \left(\frac{e^{\epsilon} \mathcal{N}^+}{s} \right) \quad \mathcal{L}(s) \quad \omega \mathcal{L}(s)$$

$$e^{\epsilon} = \mathcal{L}(s + i\epsilon)$$

$$\mathcal{N}^+(s + i\epsilon) = \mathcal{N}^+(s_+) = \mathcal{N}(s_-)$$

$$\mathcal{N}^{\pm} = \mathcal{N}(s \pm i\epsilon)$$

$$\mathcal{N} = \mathcal{P} \quad \mathcal{N}(s)$$



$$e^{\epsilon} = 1 + \epsilon e^{i\mathcal{P}(s)}$$

$$\mathcal{N}(s \pm i\epsilon) = \mathcal{N}(s) e^{\pm i\mathcal{P}(s)}$$

$$\frac{\mathcal{N}^+ - \mathcal{N}^-}{2i} = \lim_{\epsilon \rightarrow 0} \mathcal{N}(s) = \text{real part}$$

$$e^{-i\mathcal{P}(s)} = \int_{\mathcal{C}_\epsilon} \frac{ds' \mathcal{P}(s')}{s' - s}$$

$$e^{-i\mathcal{P}(s)} = 1 \quad (\text{approx})$$

$$\frac{1}{s \pm i\epsilon} = \text{P.V.} \frac{1}{s} \mp i\pi \delta(x)$$

$$\mathcal{N}(s) = \left(R + P \left(\frac{1}{T_1} s \right) \frac{\mathcal{D}(s)}{s-s} \right)$$

$$\mathcal{N}(s) \rightarrow \text{flurbivans}(s) \rightarrow \underline{E_{dr}}$$

$$\underline{\mathcal{D}(s)} = \frac{e^{Ts} - e^{-Ts}}{L}$$

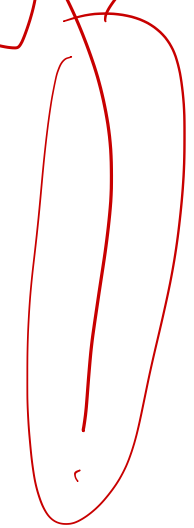




S
 $2A$

$N(s) \mathcal{O}(s)$
 (s)

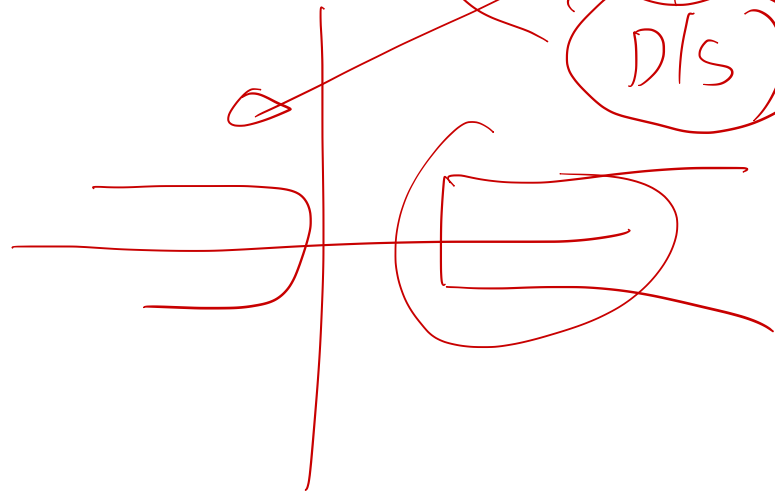
N_2



$$\underbrace{P(s)}_{\text{numerator}} \underbrace{Q(s)}_{\text{denominator}}$$

$$d(s) \Rightarrow \frac{1}{Q(s)} = \frac{1}{D(s)}$$

$$\frac{N(s)}{D(s)}$$



$$\ln \quad s^+ s = 1$$

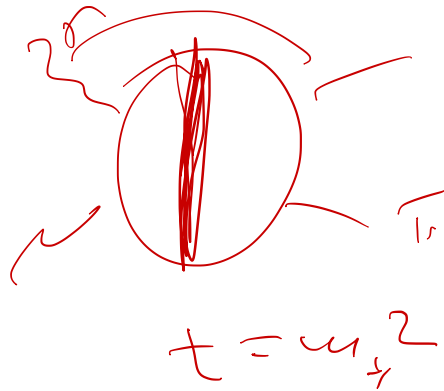
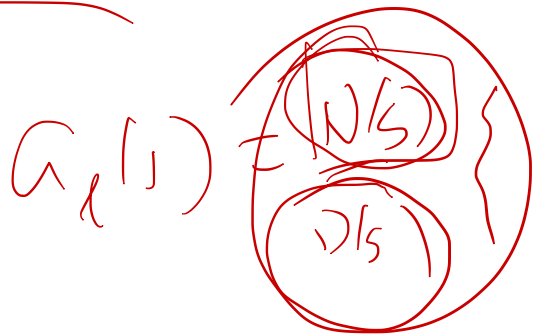
$$\ln T = \ln T^+$$

$$T = A(s, \omega)$$

$$\rightarrow \underbrace{a_p(s)}_{\text{pole}} P(\omega)$$

$$\ln \underline{a_p(s)} = \ln (a_p(s) \underbrace{g/s})$$

$$a_k(s) = \frac{1}{\pi} \int ds' \frac{\text{Im } a_k(s')}{s' - s} |a|^2$$



$N(s) = \text{same function}$

that has a branch point

at $s = s_0$ from OPE

$$S = S_0$$

$$\frac{\sqrt{s-s_0} - \sqrt{s}}{\sqrt{s-s_0} + \sqrt{s}}$$

