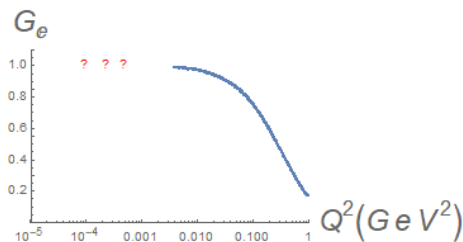


A new approach to limit the extrapolation errors of the proton size problem

Proton charge radius of e-scattering has been estimated by extrapolation of Sachs electric form factor $G_E(Q^2) = F_1(Q^2) - (Q^2/4M^2)F_2(Q^2)$ to $Q^2 = 0$.



Current estimates **CONTROVERSIAL**: fits allow $r_p \in [0.84, 89] \text{ fm}$ (Horbatsch 2016).

The plan: limit extrapolation errors by making use of *analytic structure* of G_E .

Two relevant theorems: theorem #1

Theorem #1, on the convergence of Chebyshev polynomials

Chebyshev polynomial fits to a function f on $[-1, 1]$ converge geometrically as $n \rightarrow \infty$ iff f is analytic [Bernstein 1911, 1912].

Cheby polynomials: set of orthogonal polynomial functions ...

Geometric convergence: convergence at a rate $O(C^{-n})$ for $C > 1$

Errors look like straight lines on a semilog scale.

Intro to Cheby polynomials: **Chebfun (L.N. Trefethen)**

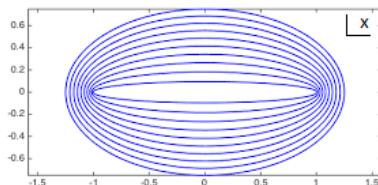
Two relevant theorems: theorem #2

Theorem #2, on Chebyshev coefficients of analytic functions

Let a function f analytic in $[-1, 1]$ be analytically continuable to the open Bernstein ellipse E_ρ , where it satisfies $|f(x)| < M$ for some M .

Then its Chebyshev coefficients satisfy $|c_0| < M$ and $|c_k| < 2M\rho^{-k}$, $k > 1$.

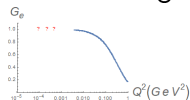
Bernstein Ellipses E_ρ



How the theorems fit together: the big picture

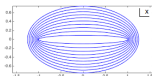
Thm #1:

We do Chebyshev fit to G_E data. Having G_E data is NOT same as having an analytic function for G_E . Thm #1 doesn't automatically apply; exponential convergence as $n \rightarrow \infty$ is not guaranteed.



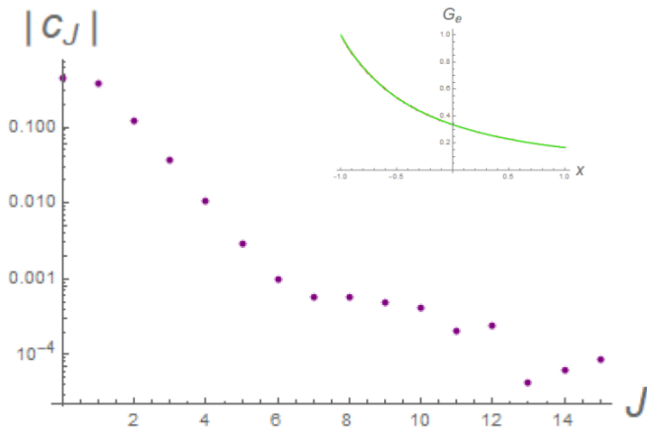
Thm #2:

FF ansatz (dipole + branch cut) + thm #2 provide a test of the convergence of the Cheby fit above— i.e. they constrain $\{|c_k|\}$.

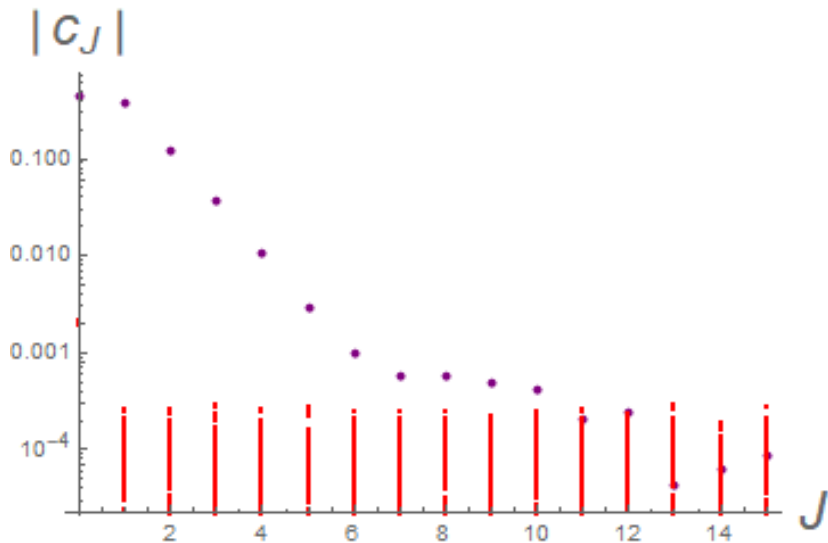


$$|c_k| < 2M\rho^{-k}, \text{ where physical singularities fix } \rho$$

Chebyshev fit to G_E data



Chebyshev fit to G_E data

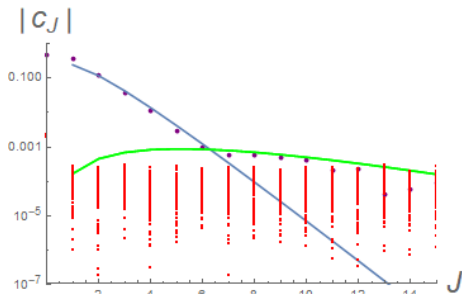


red = 2σ Gaussian noise, approximating effects of $\square \delta G_E$

Adding in the analyticity bounds

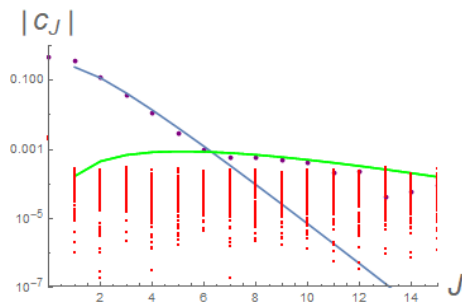
$$F_{\text{ansatz}}(Q^2) = \frac{a_1}{(Q^2 + m_{di}^2)^2} + a_2 \cdot \left(1 - \frac{\sqrt{Q^2 + 4m_\pi^2 + 0.08}}{\sqrt{Q^2 + 4m_\pi^2 - 0.08}}\right)^4 \xrightarrow{\text{thm 2}} \text{bounds}$$

Blue is constraint from dipole term; green from branch cut term.



Upshot: Cheb fit obeys analyticity bounds !

The (preliminary) final analysis

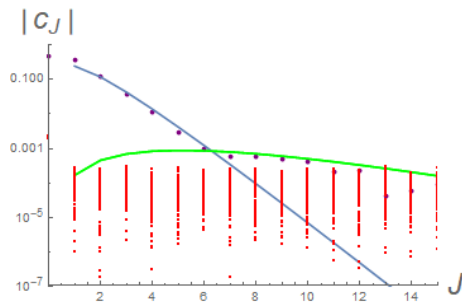


! PRELIMINARY !

We find

- 1) a radius value consistent with the 'small' radius values $\sim 0.84 \text{ fm}$ and 2) a factor of 2 reduction in the relative uncertainty of r_p .

The (preliminary) final analysis



! PRELIMINARY !

We find

- 1) a radius value consistent with the 'small' radius values $\sim 0.84 \text{ fm}$ and 2) a factor of 2 reduction in the relative uncertainty of r_p .

Thanks!

Model form factor is

$$F(Q^2) = \frac{a_1}{(Q^2 + m_{di}^2)^2} + a_2 \cdot \left(1 - \frac{\sqrt{Q^2 + 4m_\pi^2} + 0.08}{\sqrt{Q^2 + 4m_\pi^2} - 0.08}\right)^4$$

Dipole term has a singularity at $q^2 = -m_{di}^2$, that is mathematically allowed, if a bit unphysical. It does not include the width of the resonance, and its imaginary part is quite singular.

The other term has been concocted to include a branch cut from the 2-pion continuum, beginning at $Q^2 < -4m_\pi^2$ and falling fast enough as $Q^2 \rightarrow \infty$. The important thing is that singularities of the model occur at physically realistic locations in the time-like region